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**Mechanism Test Bed
Flexible Body Model
Report**

Prepared for:

**National Aeronautics and Space Administration
George C. Marshall Space Flight Center
Marshall Space Flight Center, AL 35812**

**Under Contract No.
NAS8-38771**

by

Jimmy Compton

July 1991

LOGICON

Control Dynamics Company

600 BOULEVARD SOUTH SUITE 304 ■ HUNTSVILLE, AL 35802 ■ TELEPHONE: 205 882-2650

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LIST OF ACRONYMS AND SYMBOLS

ASCII	-	American standard code for information interchange
B1	-	Body number 1 (target) center of mass/undeformed coordinate frame
B2	-	Body number 2 (chase) center of mass/undeformed coordinate frame
CDy	-	Control Dynamics Company
CM	-	Center of mass
DOF	-	Degree(s) of freedom
D1	-	Body number 1 (target) docking node/deformed coordinate frame
D2	-	Body number 2 (chase) docking node/deformed coordinate frame
EOM	-	Equations of motion
EHOMIT	-	Eliminate higher order modal integral terms
ERNLT	-	Eliminate remaining right hand side nonlinear terms
FEM	-	Finite element model
F/M	-	Force and moment
HWIL	-	Hardware in the loop
Hz	-	Hertz
K	-	Kilo
kg	-	Kilograms
LU	-	Lower-upper; refers to matrices in lower-upper triangular form
m	-	Meters
MTB	-	Mechanism Test Bed
NB	-	Body number

Nt	-	Newtons
PC-386	-	386 class personal computer
SMIS	-	Structures and matrix interpretive system
S1	-	Force and moment sensor location/coordinate frame on body number 1 (target)
wrt	-	With respect to

1.0 INTRODUCTION

The Space Station Mechanism Test Bed is a six degree of freedom (DOF) motion simulation facility used to evaluate docking and berthing hardware mechanisms. The major components of the Mechanism Test Bed (MTB) are shown in Figure 1.0-1. The chase vehicle docking mechanism is mounted on the hydraulically driven, computer controlled, six DOF motion system. The target vehicle docking mechanism is mounted in conjunction with a force and moment sensor to the facility ceiling. Mechanism contact forces and moments are measured and supplied to the host computer (i.e., currently, an Alliant computer) for use in the dynamics model.

Under contract NAS8-36570, Control Dynamics (CDy) developed a generalized rigid body math model to replace the "old" model which was based on several restrictive assumptions (e.g., one body was assumed to have much greater mass than the second and therefore was unaffected by contact forces and moments). The "new" model allowed the computation of vehicle relative motion in six DOF due to forces and moments from mechanism contact, attitude control systems, and gravity. No vehicle size limitations were imposed in the model. The equations of motions were based on Hill's equations for translational motion with respect to a nominal circular earth orbit and Newton-Euler equations for rotational motion. Over the past several years, CDy has worked with NASA in refining this rigid body model and the supporting software.

This report documents the development of a generalized flexible body math model to further enhance the MTB simulation capabilities. Although the original contract plan was to modify the current rigid body model, early investigations showed that a "fresh start" approach to the flex body model would yield a more efficient simulation. The development and major components of the flex body math model parallel those of the rigid body model.

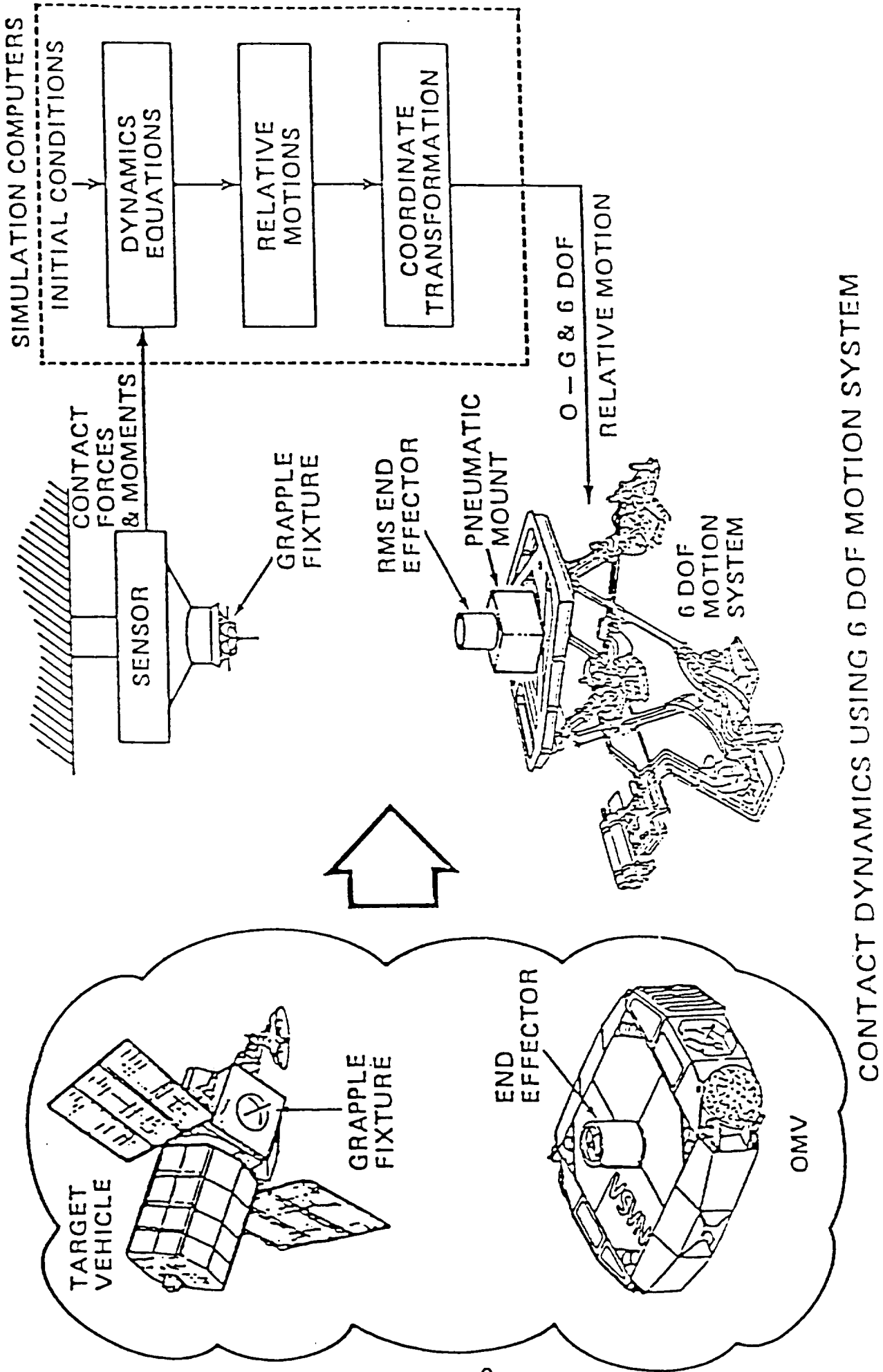


Figure 1.0-1

Section 2.0 of this report summarizes the rather mathematically intense derivation of the equations of motion for a single generic flexible body. The derivation is based on Lagrange's quasi-coordinate equations. Section 3.0 describes the method used to transform contact forces and moments from the sensor location to each body docking port. Section 4.0 discusses the computation of the relative body motion data: (1) relative orientation, (2) relative position, (3) relative translational velocity, and (4) relative angular velocity. This data is required in the interface between the dynamic math model and the main simulation. Section 5.0 describes the major components of the math model FORTRAN simulation including the required input data and a pre-processing algorithm for flexible body data. The model was coded with user selectable options regarding the method of integration and the complexity of the equations of motion. These options are discussed in Section 5.0 along with the governing input data. Section 6.0 discusses the many tests used to verify the flex body math model in progressive steps. Time domain comparisons to the multi-flex body code called TREETOPS are also presented. Section 7.0 contains a brief summary with concluding remarks. Appendices A and B contain a listing of the flexible body math model code and definitions for the math model global simulation variables, respectively. Finally, Appendix C contains a listing of the flexible body data pre-processing algorithm.

2.0 FLEXIBLE BODY MODEL

Table 2.0-1 lists the most popular methods, along with advantages and disadvantages, for deriving equations of motion (EOM) of dynamical systems. The D'Alembert / Newton-Euler equations are best suited for rigid body applications; for complex systems, the form of the equations depend to a large extent on the insight of the analyst. The other methods (i.e., Lagrange's Eqns. and Kane's Eqns.) offer a more systematic approach.

Lagrange's generalized coordinate equations are represented by the vector equation

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\mathbf{q}}} \right) - \frac{\partial L}{\partial \mathbf{q}} = \mathbf{Q} \quad (2.0-1)$$

where L = kinetic energy - potential energy (Lagrangian)

\mathbf{q} = vector of generalized coordinates

\mathbf{Q} = vector of generalized forces.

This "standard" form of Lagrange's equations is well known and produces EOM with the same basic structure for all applications. A limitation of Lagrange's "standard" equations is that the generalized coordinate vector derivative ($\dot{\mathbf{q}}$) must be integrable to a set of coordinates (\mathbf{q}) which completely describe the system configuration. Specifically, this precludes the use of body angular rates ($\boldsymbol{\omega}$) in the \mathbf{q} vector (i.e., the integral of $\{\boldsymbol{\omega} dt\}$ does not adequately describe body attitude). This limitation can be overcome by employing a modified form of Lagrange's equations known as Lagrange's quasi-coordinate equations or the Boltzmann-Hamel equations. Quasi-coordinates (\mathbf{w}) are quantities whose differentials may be written as linear combinations of generalized coordinate derivatives, e.g.,

Table 2.0-1 Candidate Methods for Equations of Motion Derivation.

METHOD	ADVANTAGES	DISADVANTAGES
1. D'Alembert / Newton - Euler Eqs.	Physically visualizable quantities Vector / dyadic representation	Max number of DOFs No standard equation form
2. Lagrange's Generalized Coordinate Eqs.	Systematic approach Invariant D.E. structure Equations facilitate integration	Differentiation of "complex" scalar energy functions. Must use "integrable" DOFs
3. Lagrange's Quasi-Coordinate Eqs. (Boltzmann-Hamel)	Subsumes method (2) Relieves constraint of using generalized coordinate derivatives: $T(\mathbf{q}, \dot{\mathbf{q}}, t) \rightarrow \bar{T}(\mathbf{w}, \mathbf{q}, t)$	Differentiation of energy functions
4. Kane's Eqs. (TREETOPS)	Systematic approach Reduces system to minimal representation	Lack of physical insight Not as well known

$$\frac{d}{dt}(w_1) = a_1 \frac{d}{dt}(q_1) + a_2 \frac{d}{dt}(q_2) + \dots \quad (2.0-2)$$

Thus, the quasi-coordinate form of Lagrange's equations allows the direct use of body angular rates in the EOM coordinate vector which replaces the q coordinate vector (i.e., $(q, \dot{q}, t) \rightarrow (w, \dot{q}, t)$). Lagrange's quasi-coordinate method was used to derive the EOM for the MTB generic flexible body model.

Kane's equations combines certain advantages from D'Alembert's principle and Lagrange's method. The multi-flex body code TREETOPS is based on Kane's equations. However, Kane's method is not as well known as the Lagrangian formulation.

2.1 Lagrange's Quasi-Coordinate Equations

The quasi-coordinate vector definition used in the flex body model is

$$w = \begin{Bmatrix} \dot{p}^B \\ \omega^B \\ \dot{\eta} \end{Bmatrix} \quad (2.1-1)$$

where \dot{p}^B = translational velocity of undeformed body CM with respect to (wrt) inertial space expressed in the body frame

ω^B = angular velocity of undeformed body frame wrt inertial space expressed in the body frame

$\dot{\eta}$ = time derivative of the generalized modal coordinate vector.

The associated generalized coordinate vector was defined using the four Euler parameters (i.e., a quaternion) for attitude

$$\mathbf{q} = \begin{Bmatrix} \mathbf{p}^I \\ \boldsymbol{\epsilon}_{BI} \\ \boldsymbol{\eta} \end{Bmatrix} \quad (2.1-2)$$

where \mathbf{p}^I = position of undeformed body CM wrt inertial space
expressed in the inertial frame

$\boldsymbol{\epsilon}_{BI}$ = inertial to body quaternion

$\boldsymbol{\eta}$ = generalized modal coordinate vector.

Notice that the generalized coordinates are not independent due to the unity magnitude quaternion constraint

$$\left[\sum_{i=1}^4 (\epsilon_{BI_i})^2 \right]^{1/2} = 1 \quad (2.1-3)$$

The transformation from the quasi-coordinate vector to the derivative of the generalized coordinate vector can be expressed as

$$\dot{\mathbf{q}} = \begin{bmatrix} [\mathbf{B} \mathbf{I}]^T & 0 & 0 \\ 0 & \frac{1}{2} \mathbf{E}^T & 0 \\ 0 & 0 & \mathbf{I} \end{bmatrix} \mathbf{w} \quad (2.1-4)$$

where $[\mathbf{B} \mathbf{I}]^T$ = 3 x 3 body to inertial transformation matrix

$(1/2) \mathbf{E}^T$ = 4 x 3 transformation from the body angular velocity
vector to the quaternion time derivative

$$\mathbf{E}^T = \begin{bmatrix} \epsilon_4 & -\epsilon_3 & \epsilon_2 \\ \epsilon_3 & \epsilon_4 & -\epsilon_1 \\ -\epsilon_2 & \epsilon_1 & \epsilon_4 \\ -\epsilon_1 & -\epsilon_2 & -\epsilon_3 \end{bmatrix}$$

$I = 3 \times 3$ identity matrix.

Starting with Lagrange's "standard" equations (2.0-1), Lagrange's quasi-coordinate equations can be derived for the coordinate vector definitions given in eqns. (2.1-1) through (2.1-4). The EOM separate nicely into vector equations for translation, rotation, and flexibility. The resulting equations can be expressed in general as

Translational EOM:

$$\frac{d}{dt} \left(\frac{\partial \bar{T}}{\partial \dot{\mathbf{P}}^B} \right) + \boldsymbol{\omega}^B \times \frac{\partial \bar{T}}{\partial \dot{\mathbf{P}}^B} - [BI] \frac{\partial \bar{T}}{\partial \mathbf{P}^I} + [BI] \frac{\partial V^*}{\partial \mathbf{P}^I} = [BI] \mathbf{Q}_P \quad (2.1-5)$$

Rotational EOM:

$$\frac{d}{dt} \left(\frac{\partial \bar{T}}{\partial \boldsymbol{\omega}^B} \right) + \boldsymbol{\omega}^B \times \frac{\partial \bar{T}}{\partial \boldsymbol{\omega}^B} + \dot{\mathbf{P}}^B \times \frac{\partial \bar{T}}{\partial \dot{\mathbf{P}}^B} - \frac{1}{2} \mathbf{E} \frac{\partial \bar{T}}{\partial \boldsymbol{\varepsilon}} + \frac{1}{2} \mathbf{E} \frac{\partial V^*}{\partial \boldsymbol{\varepsilon}} = \frac{1}{2} \mathbf{E} \mathbf{Q}_\varepsilon \quad (2.1-6)$$

Flexibility EOM:

$$\frac{d}{dt} \left(\frac{\partial \bar{T}}{\partial \dot{\boldsymbol{\eta}}} \right) - \frac{\partial \bar{T}}{\partial \boldsymbol{\eta}} + \frac{\partial V^*}{\partial \boldsymbol{\eta}} = \mathbf{Q}_\eta \quad (2.1-7)$$

where \bar{T} = kinetic energy function in terms of quasi-coordinates

V^* = potential energy function augmented with constraint equations.

Notice that application of Lagrange's equations requires a derivation of both kinetic and potential energy functions (scalars) for the system under consideration.

Constraints are included with the potential energy function using the Lagrange multiplier technique to form an augmented potential energy function

$$V^* = V(q) - \lambda^T \phi(q) \quad (2.1-8)$$

where λ is a vector of Lagrange multipliers and the constraints are of the form

$$\phi(q) = 0. \quad (2.1-9)$$

For an unconstrained flexible body (as considered here), the only constraint equations will be those resulting from the use of dependent coordinates (e.g., quaternion for attitude).

2.2 Flexible Body Kinetic Energy Function

The kinetic energy function for a flexible body can be derived using a mass particle approach. Using the configuration definitions given in Fig. 2.2-1, the inertial position of the i^{th} mass particle of a generic flexible body may be written as

$$\mathbf{r}_i = \mathbf{R}_o + \mathbf{p}_i + \mathbf{\delta}_i. \quad (2.2-1)$$

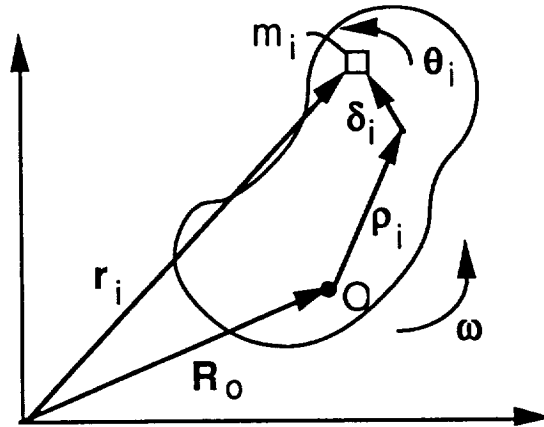
Taking the time derivative of the mass particle position vector with respect to the inertial frame leads to

$$\dot{\mathbf{r}}_i = \dot{\mathbf{R}}_o + \boldsymbol{\omega} \times (\mathbf{p}_i + \mathbf{\delta}_i) + (\dot{\mathbf{\delta}}_i)_B \quad (2.2-2)$$

where $(\dot{\mathbf{\delta}}_i)_B$ is a derivative wrt the body fixed frame.

The total kinetic energy (i.e., translational plus rotational contributions) for a body may be written as

$$T = \frac{1}{2} \sum_i \left[m_i (\dot{\mathbf{r}}_i \cdot \dot{\mathbf{r}}_i) + (\boldsymbol{\omega} + \dot{\boldsymbol{\theta}}_i)^T \mathbf{I}_i (\boldsymbol{\omega} + \dot{\boldsymbol{\theta}}_i) \right] \quad (2.2-3)$$



- m_i = i^{th} mass particle
- O = reference point on flexible body
- r_i = mass particle inertial position vector
- R_o = body reference point inertial position vector
- p_i = mass particle nominal position vector wrt O
- δ_i = translational deformation
- θ_i = rotational deformation
- ω = "rigid body" angular velocity vector

Figure 2.2-1 Generic Flexible Body Configuration.

The somewhat lengthy derivation proceeds by substituting eqn. (2.2-2) into eqn. (2.2-3) and manipulating the kinetic energy function into a useful form. The translational and rotational elastic deformation vectors for the i^{th} mass particle are computed from modal expansions as given in eqn. (2.2-4).

$$\delta_i = \sum_j^{\text{\#modes}} \Phi_{ij} \eta_j \quad \theta_i = \sum_j^{\text{\#modes}} \Psi_{ij} \eta_j \quad (2.2-4)$$

The translational and rotational mode shapes (Φ and Ψ , respectively) are commonly obtained from a finite element model (FEM) of the body. The resulting kinetic energy function is written in "shorthand" notation as

$$\begin{aligned}
T = & \frac{1}{2} \mathbf{M} \dot{\mathbf{R}}_o^B + \frac{1}{2} \dot{\boldsymbol{\eta}}^T \mathbf{M} \dot{\boldsymbol{\eta}} + \dot{\mathbf{R}}_o^B \cdot \left[\boldsymbol{\omega} \times \left(\overline{\mathbf{M}} \boldsymbol{\rho} + \sum_k^{nm} \Gamma_{ok} \boldsymbol{\eta}_k \right) \right] + \dot{\mathbf{R}}_o^B \cdot \sum_k^{nm} (\Gamma_{ok} \dot{\boldsymbol{\eta}}_k) \\
& + \boldsymbol{\omega} \cdot \sum_j^{nm} \sum_k^{nm} (\Gamma_{1k} \dot{\boldsymbol{\eta}}_k + \Gamma_{2jk} \boldsymbol{\eta}_j \dot{\boldsymbol{\eta}}_k) + \frac{1}{2} \boldsymbol{\omega}^T \sum_j^{nm} \sum_k^{nm} (\mathbf{I}_{2jk} \boldsymbol{\eta}_j \boldsymbol{\eta}_k + 2 \mathbf{I}_{1k} \boldsymbol{\eta}_k + \mathbf{I}_o) \boldsymbol{\omega}
\end{aligned}
\tag{2.2-5}$$

where Γ_{ok} and Γ_{1k} are the "standard" modal integrals and Γ_{2jk} , \mathbf{I}_{1k} , and \mathbf{I}_{2jk} are called the higher order modal integrals. The modal integral terms are functions of the mode shapes, FEM element masses, and FEM element inertias. The modal integrals are constants and appear in the final form of the equations of motion; they are typically computed from the FEM output data for use as inputs to a flex body dynamic simulation.

2.3 Flexible Body Potential Energy Function

The potential energy function for a flexible body may be assumed to consist only of the strain energy due to linear elastic deformation. In this case, all other conservative forces which are derivable from potential functions (e.g., gravity effects) must be applied as forcing functions (i.e., external or generalized forces) on the right hand side of the EOM.

The strain energy in a linear elastic material may be expressed in the form

$$V = \frac{1}{2} \bar{\boldsymbol{\delta}}^T \bar{\mathbf{K}} \bar{\boldsymbol{\delta}} \tag{2.3-1}$$

where $\bar{\boldsymbol{\delta}}$ is the deformation due to flexibility and $\bar{\mathbf{K}}$ is the FEM stiffness matrix. Here $\bar{\boldsymbol{\delta}}$ represents both the translational and rotational deformation. Replacing $\bar{\boldsymbol{\delta}}$ in eqn. (2.3-1) with a modal expansion yields

$$V = \frac{1}{2} \boldsymbol{\eta}^T \bar{\boldsymbol{\Phi}}^T \bar{\mathbf{K}} \bar{\boldsymbol{\Phi}} \boldsymbol{\eta} \tag{2.3-2}$$

where the modal matrix $\overline{\Phi}$ includes both the translational and rotational modal gains. If the mode shapes, $\overline{\Phi}$, are the system orthonormal eigenvectors, then a new diagonalized stiffness matrix can be defined

$$\mathbf{K} = \overline{\Phi}^T \overline{\mathbf{K}} \overline{\Phi} = \begin{bmatrix} \omega_1^2 & 0 & 0 \\ 0 & \omega_2^2 & 0 \\ 0 & 0 & \omega_n^2 \end{bmatrix} \quad (2.3-3)$$

where the diagonal elements are the squared natural frequencies of the system modes. The assumption of orthonormal modes follows from NASA's request in the statement of work for the capability to input second order bending modes of the form

$$\ddot{\eta}_i + 2 \zeta_i \omega_i \dot{\eta}_i + \omega_i^2 \eta_i = F_i. \quad (2.3-4)$$

The potential energy or strain energy may be written in terms of the diagonalized stiffness matrix

$$V = \frac{1}{2} \boldsymbol{\eta}^T \mathbf{K} \boldsymbol{\eta} \quad (2.3-5)$$

Notice that the potential energy function is a function of the generalized modal coordinates $\boldsymbol{\eta}$.

The augmented potential energy function used in Lagrange's quasi-coordinates eqns. (2.1-5) through (2.1-7) included constraints with Lagrange multipliers. The only constraint equation associated with the unconstrained single flex body model is the unity magnitude quaternion constraint. The augmented potential energy function can now be written as

$$V^* = V(\mathbf{q}) - \boldsymbol{\lambda}^T \boldsymbol{\phi}(\mathbf{q})$$

$$V^* = \frac{1}{2} \eta^T K \eta^T - \lambda_1 (\epsilon_1^2 + \epsilon_2^2 + \epsilon_3^2 + \epsilon_4^2 - 1) \quad (2.3-6)$$

where λ_1 is the Lagrange multiplier.

2.4 Equations of Motion

The kinetic energy function (2.2-5) and the augmented potential energy function (2.3-6) may now be incorporated into Lagrange's quasi-coordinate equations (2.1-5) through (2.1-7). The body reference point (i.e., point O in Fig. 2.2-1) is assumed to be located at the body center of mass in order to simplify the equations somewhat. After lengthy algebraic manipulations, the "full-up" nonlinear EOM can be expressed in the following "shorthand" notation

$$\begin{bmatrix} \mathbf{M} & \widetilde{\Gamma_0 \eta} & \Gamma_0 \\ \widetilde{\Gamma_0 \eta} & \overline{\mathbf{I} \mathbf{I}} & \Gamma_1 + \Gamma_2 \\ \Gamma_0^T & \Gamma_1^T + \Gamma_2^T & M \end{bmatrix} \begin{Bmatrix} \ddot{\mathbf{P}}_{CM}^B \\ \dot{\boldsymbol{\omega}}^B \\ \ddot{\boldsymbol{\eta}} \end{Bmatrix} = \mathbf{F}_L + \mathbf{F}_{NL} + \mathbf{F}_{APPL}$$

$$= \left(\begin{aligned} & -\boldsymbol{\omega}^B \times (M \dot{\mathbf{P}}_{CM}^B + 2 \Gamma_0 \eta) - \boldsymbol{\omega}^B \times (\boldsymbol{\omega}^B \times \Gamma_0 \eta) + \sum_m^{APPL} \mathbf{F}_m^B \\ & - \dot{\mathbf{I} \mathbf{I}} \boldsymbol{\omega}^B - \boldsymbol{\omega}^B \times \mathbf{I} \boldsymbol{\omega}^B - \Gamma_0 \eta \times (\boldsymbol{\omega}^B \times \dot{\mathbf{P}}_{CM}^B) - \boldsymbol{\omega}^B \times (\Gamma_1 + \Gamma_2) \ddot{\boldsymbol{\eta}} + \sum_n^{APPL} \mathbf{T}_n^B \\ & - \boldsymbol{\omega}_k^2 \eta - 2 \zeta_k \boldsymbol{\omega}_k \eta + \dot{\mathbf{P}}_{CM}^B \cdot (\boldsymbol{\omega}^B \times \Gamma_0) + \sum_m^{APPL} \Phi_m^T \mathbf{F}_m^B + \sum_n^{APPL} \Psi_n^T \mathbf{T}_n^B \\ & - 2 \boldsymbol{\omega}^B \cdot \dot{\Gamma}_2 + \frac{1}{2} \boldsymbol{\omega}^{BT} \left[2 \mathbf{I}_{1r} + \sum_j^{nm} (\mathbf{I}_{2jr} + \mathbf{I}_{2rj}) \eta_j \right] \boldsymbol{\omega}^B \end{aligned} \right) \quad (2.4-1)$$

where parts = number of mass particles
 appl = number of externally applied forces or torques.

$$\begin{aligned}
 \mathbf{M} &= \begin{bmatrix} \mathbf{M} & 0 & 0 \\ 0 & \mathbf{M} & 0 \\ 0 & 0 & \mathbf{M} \end{bmatrix} & M &= \sum_i^{\text{parts}} \begin{bmatrix} \Phi_i \\ \Psi_i \end{bmatrix} \begin{bmatrix} \mathbf{M}_i & 0 \\ 0 & \mathbf{I}_i \end{bmatrix} \begin{bmatrix} \Phi_i \\ \Psi_i \end{bmatrix} \\
 \Gamma_0 &= [\Gamma_{01} \ \Gamma_{02} \ \dots \ \Gamma_{0n}] & \Gamma_1 &= [\Gamma_{11} \ \Gamma_{12} \ \dots \ \Gamma_{1n}] \\
 \Gamma_2 &= \sum_j^{nm} \eta_j [\Gamma_{2j1} \ \Gamma_{2j2} \ \dots \ \Gamma_{2jn}] & \dot{\Gamma}_2 &= \sum_j^{nm} \dot{\eta}_j [\Gamma_{2j1} \ \Gamma_{2j2} \ \dots \ \Gamma_{2jn}] \\
 \Phi_i &= [\Phi_{i1} \ \Phi_{i2} \ \dots \ \Phi_{in}] & \Psi_i &= [\Psi_{i1} \ \Psi_{i2} \ \dots \ \Psi_{in}] \\
 \Pi &= \mathbf{I}_{CM} + (\mathbf{I}_{1k} + \mathbf{I}_{1k}^T) \eta_k + \frac{1}{2} (\mathbf{I}_{2jk} + \mathbf{I}_{2kj}) \eta_j \eta_k \\
 \dot{\Pi} &= (\mathbf{I}_{1k} + \mathbf{I}_{1k}^T) \dot{\eta}_k + \frac{1}{2} (\mathbf{I}_{2jk} + \mathbf{I}_{2kj}) \dot{\eta}_j \eta_k + \frac{1}{2} (\mathbf{I}_{2jk} + \mathbf{I}_{2kj}) \eta_j \dot{\eta}_k \quad (2.4-2)
 \end{aligned}$$

"Standard" Modal Integrals:

$$\begin{aligned}
 \Gamma_{0k} &= \sum_i^{\text{parts}} m_i \Phi_{ik} \\
 \Gamma_{1k} &= \sum_i^{\text{parts}} (m_i \rho_i \times \Phi_{ik} + \mathbf{I}_i \Psi_{ik}) \quad (2.4-3)
 \end{aligned}$$

Higher Order Modal Integrals:

$$\begin{aligned}
 \Gamma_{2jk} &= \sum_i^{\text{parts}} m_i \Phi_{ij} \times \Phi_{ik} \\
 \mathbf{I}_{1k} &= \sum_i^{\text{parts}} m_i [(\rho_i \cdot \Phi_{ik}) \mathbf{1} - \Phi_{ik} \rho_i^T] \\
 \mathbf{I}_{2jk} &= \sum_i^{\text{parts}} m_i [(\Phi_{ij}^T \Phi_{ik}) \mathbf{1} - \Phi_{ij} \Phi_{ik}^T] \quad (2.4-4)
 \end{aligned}$$

where $\text{parts} = \text{number of mass particles}$.

Notice that for the "full-up" nonlinear case the mass matrix has time varying elements; this implies that periodic updates of the time varying elements are required in the flexible body math model.

Examination of the "full-up" nonlinear EOM reveals that the equations can be written in levels of increasing simplicity by neglecting certain higher order effects. First, neglecting the higher order modal integral terms allows the EOM to be written as

$$\begin{bmatrix} \mathbf{M} & -\widetilde{\Gamma_0 \eta} & \Gamma_0 \\ \widetilde{\Gamma_0 \eta} & \mathbf{I}_{CM} & \Gamma_1 \\ \Gamma_0^T & \Gamma_1^T & M \end{bmatrix} \begin{bmatrix} \ddot{\mathbf{P}}_{CM}^B \\ \dot{\boldsymbol{\omega}}^B \\ \ddot{\boldsymbol{\eta}} \end{bmatrix} = \begin{bmatrix} -\boldsymbol{\omega}^B \times (M \dot{\mathbf{P}}_{CM}^B + 2 \Gamma_0 \dot{\boldsymbol{\eta}}) - \boldsymbol{\omega}^B \times (\boldsymbol{\omega}^B \times \Gamma_0 \boldsymbol{\eta}) + \sum_m^{\text{APPL}} \mathbf{F}_m^B \\ -\boldsymbol{\omega}^B \times \mathbf{I}_{CM} \boldsymbol{\omega}^B - \Gamma_0 \dot{\boldsymbol{\eta}} \times (\boldsymbol{\omega}^B \times \dot{\mathbf{P}}_{CM}^B) - \boldsymbol{\omega}^B \times (\Gamma_1 \dot{\boldsymbol{\eta}}) + \sum_n^{\text{APPL}} \mathbf{T}_n^B \\ -\omega_k^2 \boldsymbol{\eta} - 2 \zeta_k \omega_k \dot{\boldsymbol{\eta}} + \dot{\mathbf{P}}_{CM}^B \cdot (\boldsymbol{\omega}^B \times \Gamma_0) + \sum_m^{\text{APPL}} \boldsymbol{\Phi}_m^T \mathbf{F}_m^B + \sum_n^{\text{APPL}} \boldsymbol{\Psi}_n^T \mathbf{T}_n^B \end{bmatrix} \quad (2.4-5)$$

Eliminating the higher order modal integrals reduces the computations per simulation cycle. Notice in eqn. (2.4-5) that if the "standard" modal integrals were zero (i.e., free-free modes assumption discussed below), then neglecting the higher order modal integrals would eliminate the time varying terms from the mass matrix; this would allow a "one-time" computation of the mass matrix in the initialization and a significant decrease in simulation cycle time.

Secondly, elimination of the remaining nonlinear terms (i.e., includes elimination of the higher order modal integrals) on the right hand side of the EOM leads to another level of simplification

$$\begin{bmatrix} \mathbf{M} & -\widetilde{\Gamma_0} \boldsymbol{\eta} & \Gamma_0 \\ \widetilde{\Gamma_0} \boldsymbol{\eta} & \mathbf{I}_{CM} & \Gamma_1 \\ \Gamma_0^T & \Gamma_1^T & M \end{bmatrix} \begin{Bmatrix} \ddot{\mathbf{P}}_{CM}^B \\ \dot{\boldsymbol{\omega}}^B \\ \ddot{\boldsymbol{\eta}} \end{Bmatrix} = \begin{pmatrix} \sum_m^{APPL} \mathbf{F}_m^B \\ \sum_n^{APPL} \mathbf{T}_n^B \\ -\omega_k^2 \boldsymbol{\eta} - 2 \zeta_k \omega_k \dot{\boldsymbol{\eta}} + \sum_m^{APPL} \boldsymbol{\Phi}_m^T \mathbf{F}_m^B + \sum_n^{APPL} \boldsymbol{\Psi}_n^T \mathbf{T}_n^B \end{pmatrix} \quad (2.4-6)$$

Eliminating the nonlinear terms on the right hand side greatly reduces the computational requirements per simulation cycle and again decreases cycle time.

Another simplification level results from the use of free-free modes for the unconstrained flexible body. The "standard" modal integral terms are identically zero when free-free modes are used. The result is a significant simplification of the mass matrix (and the right hand side if the previous simplification was not invoked)

$$\begin{bmatrix} \mathbf{M} & 0 & 0 \\ 0 & \mathbf{I}_{CM} & 0 \\ 0 & 0 & M \end{bmatrix} \begin{Bmatrix} \ddot{\mathbf{P}}_{CM}^B \\ \dot{\boldsymbol{\omega}}^B \\ \ddot{\boldsymbol{\eta}} \end{Bmatrix} = \begin{pmatrix} \sum_m^{APPL} \mathbf{F}_m^B \\ \sum_n^{APPL} \mathbf{T}_n^B \\ -\omega_k^2 \boldsymbol{\eta} - 2 \zeta_k \omega_k \dot{\boldsymbol{\eta}} + \sum_m^{APPL} \boldsymbol{\Phi}_m^T \mathbf{F}_m^B + \sum_n^{APPL} \boldsymbol{\Psi}_n^T \mathbf{T}_n^B \end{pmatrix} \quad (2.4-7)$$

The first and second simplifications resulting in eqns (2.4-5) and (2.4-6) are not prerequisites for applying the free-free mode simplification. However, using the free-free mode simplification in conjunction with eliminating the higher order modal integral terms produces a constant mass matrix which can be computed once in the initialization. Notice that the translational, rotational, and flexibility equations are completely decoupled with respect to one another when the following simplifications are applied: (1) neglect higher order mass integral terms, (2) neglect remaining right hand side nonlinear terms, and (3) utilize free-free modes.

NOTE: It is important to invoke the free-free mode option explicitly in the math model configuration data when free-free modes are being used to characterize body flexibility. Invoking the free-free option eliminates unnecessary computations that will otherwise be performed with standard modal integrals which are zero (or very small due to numerical rounding).

A final simplification of the EOM results when no flex modes are included (i.e., rigid body). In this case, the EOM reduce to the well known linear form of the Newton-Euler equations for translation and rotation

$$\begin{bmatrix} \mathbf{M} & 0 \\ 0 & \mathbf{I}_{CM} \end{bmatrix} \begin{Bmatrix} \ddot{\mathbf{P}}_{CM}^B \\ \dot{\boldsymbol{\omega}}^B \end{Bmatrix} = \begin{Bmatrix} \sum_m^{APPL} \mathbf{F}_m^B \\ \sum_n^{APPL} \mathbf{T}_n^B \end{Bmatrix} \quad (2.4-8)$$

All of the above mentioned EOM complexity levels are made available to the user of the two-flex body math model. The user will configure the math model to the desired complexity level through the use of flags and parameters in ASCII input data files.

2.5 Equations of Motion Solution

Regardless of the level of complexity, the equations of motion can be expressed in the general form

$$\mathbf{A} \cdot \mathbf{x} = \mathbf{b} \quad (2.5-1)$$

where \mathbf{A} = system mass matrix
 \mathbf{x} = quasi-coordinate vector derivative
 \mathbf{b} = right hand side vector.

The mass matrix (\mathbf{A}) and the right hand side vector (\mathbf{b}) are computed (or known) at each point in time. The objective is to solve the linear equation set for the quasi-coordinate vector derivative (\mathbf{x}) such that numerical integration can be performed to propagate the system states.

Recall that depending on the EOM complexity options selected by the user, the system mass matrix may or may not vary with time. If the mass matrix is not time varying (i.e., constant elements) then a "one-time" solution approach to the linear equation set of (2.5-1) is desired. This would allow a portion of the computations associated with solving (2.5-1) to be performed once in an initialization routine and, thus, reduce the math model cycle time.

A numerical manipulation method called LU (lower-upper) decomposition was selected for solving eqn. (2.5-1). The LU decomposition decomposes the mass matrix into a product of a lower triangular matrix (\mathbf{L}) and an upper triangular matrix (\mathbf{U}).

$$\mathbf{A} \cdot \mathbf{x} = (\mathbf{L} \cdot \mathbf{U}) \mathbf{x} = \mathbf{L} \cdot (\mathbf{U} \cdot \mathbf{x}) = \mathbf{b}$$
$$\mathbf{L} = \begin{bmatrix} \alpha_{11} & 0 & 0 \\ \alpha_{21} & \alpha_{22} & 0 \\ \alpha_{n1} & \alpha_{n2} & \alpha_{nn} \end{bmatrix} \quad \mathbf{U} = \begin{bmatrix} \beta_{11} & \beta_{12} & \beta_{1n} \\ 0 & \beta_{22} & \beta_{2n} \\ 0 & 0 & \beta_{nn} \end{bmatrix} \quad (2.5-2)$$

The decomposition is based on a numerical manipulation technique called Crout's algorithm. Once the mass matrix has been decomposed, the problem of solving for the desired vector \mathbf{x} is broken into two smaller linear equation problems which can be solved using backwards and forwards substitution

$$\mathbf{L} \cdot \mathbf{y} = \mathbf{b} \quad (\text{forward}) \quad (2.5-3)$$

$$\mathbf{U} \cdot \mathbf{x} = \mathbf{y} \quad (\text{backward}) \quad (2.5-4)$$

Eqn. (2.5-3) can be solved using forward substitution since \mathbf{L} is lower triangular; eqn. (2.5-4) can be solved using backward substitution since \mathbf{U} is upper triangular.

Notice that the LU decomposition of the mass matrix can be used to solve for \mathbf{x} with any right hand side \mathbf{b} . Thus, for a constant mass matrix the LU decomposition needs only to be computed once during initialization; given a new right hand side vector \mathbf{b} during each simulation cycle, only the forward/backward substitution operations associated with eqns. (2.5-3) and (2.5-4) are required to solve for \mathbf{x} . Another possible solution method is to compute the mass matrix inverse \mathbf{A}^{-1} . For the constant mass matrix case only a matrix-vector multiplication (given the computed \mathbf{b} vector) would be required during each cycle to solve for \mathbf{x}

$$\mathbf{x} = \mathbf{A}^{-1} \mathbf{b}. \quad (2.5-4)$$

However, for the time varying mass matrix case the LU decomposition (performed each cycle) is numerically more efficient than computing the matrix inverse.

3.0 CONTACT FORCE / MOMENT TRANSFORMATION

During hardware-in-the-loop (HWIL) simulations, the math model requires input from the force/moment (F/M) sensor located on the ceiling of the MTB six DOF facility. The sensor location is assumed to correspond to some arbitrary position on the target body. Transformations are required in order to determine the appropriate force/moment applied to the docking node of each body. The body docking node is defined as the single point on the body through which all contact forces and moments are applied to the body.

The F/M sensor measures the total forces and torques applied to the target body due to hardware contact between the target and chase bodies. It is assumed that a rigid link exists in the target body between the F/M sensor and the docking mechanism; therefore, all contact forces and moments, from all points of contact, are detected by the F/M sensor.

It is also assumed that the docking mechanism mass (for each body) is much smaller than the body masses. This assumption implies that the inertial forces acting on the moving parts of the docking mechanisms may be neglected. In this case, the contact forces and moments between the two bodies are equal and opposite for collocated docking nodes.

Figure 3.0-1 defines the position vectors from the F/M sensor to the target docking node (D1) and to the chase docking node (D2). Due to the rigid link assumption, \mathbf{P}_{D1S1} remains constant in a frame attached to either the sensor location or the target docking node. However, \mathbf{P}_{D2S1} includes the effects of deformation of both the sensor location (or D1 since they are rigidly connected) and the chase docking node. \mathbf{F}_{S1} and \mathbf{M}_{S1} are defined as the sensed force and moment. \mathbf{F}_{D1} , \mathbf{M}_{D1} , \mathbf{F}_{D2} and \mathbf{M}_{D2} are defined as the contact force and moment applied at the target docking node and chase docking node, respectively.

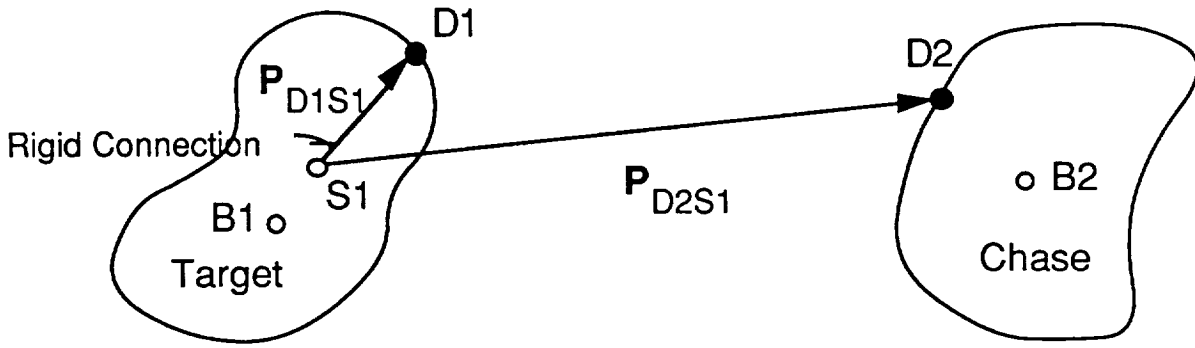


Figure 3.0-1 F/M Sensor to Docking Node Position Vector Definitions.

Due to the rigid link assumption, the forces can be transferred directly from the sensor location to the docking nodes as

$$\mathbf{F}_{D1} = -\mathbf{F}_{D2} = \mathbf{F}_{S1} \quad (3.0-1)$$

The moment transformations include the effect of non collocated docking nodes and are expressed as

$$\mathbf{M}_{D1} = \mathbf{M}_{S1} - \mathbf{P}_{D1S1} \times \mathbf{F}_{D1} \quad (3.0-2)$$

$$\mathbf{M}_{D2} = -\mathbf{M}_{S1} - \mathbf{P}_{D2S1} \times \mathbf{F}_{D2} \quad (3.0-3)$$

Notice that if the docking nodes are collocated (i.e., $\mathbf{P}_{D2S1} = \mathbf{P}_{D1S1}$) the contact moments are also equal and opposite.

4.0 RELATIVE BODY MOTION

The role of the two-flex body math model is to simulate the dynamic response of the two bodies when acted upon by contact, control, and any other modeled external forces and torques. The most important output from the math model is the relative motion between the target and chase docking nodes. This data is used to generate the commands which drive the six DOF motion system.

The transformation matrix describing the orientation of the chase docking frame with respect to the target docking frame must include the effects of flexibility. This is most easily expressed by defining undeformed node frames (i.e., rigid body node frames) for use as intermediate transform orientations. Following this approach, the 3 x 3 transformation matrix from the D2 to the D1 node frame may be written as

$$\begin{aligned} [D1 \ D2] &= [I \ D1]^T [I \ D2] \\ &= [D1_0 \ D1]^T [B1 \ D1_0]^T [I \ B1]^T [I \ B2] [B2 \ D2_0] [D2_0 \ D2] \end{aligned} \quad (4.0-1)$$

where $D1_0$ and $D2_0$ are undeformed docking node coordinate frames and I represents the inertial reference frame. The transformation from the deformed node frame to the undeformed node frame can be expressed in terms of the infinitesimal rotations about each axis

$$[D1_0 \ D1] = [R_{x,d\theta_x} \ R_{y,d\theta_y} \ R_{z,d\theta_z}] = [1 + \widetilde{d\theta}_{D1}] \quad (4.0-2)$$

Figure 4.0-1 defines the vectors which are used in the equations describing the relative motion. The vectors representing deformation are not shown explicitly in Fig. 4.0-1. The position of the chase docking node with respect to the target docking node may be expressed as

$$\mathbf{P}_{D2D1} = \mathbf{P}_{B2I} + \mathbf{P}_{D2B2} + \boldsymbol{\delta}_{D2D2_0} - \mathbf{P}_{B1I} - \mathbf{P}_{D1_0B1} - \boldsymbol{\delta}_{D1D1_0} \quad (4.0-3)$$

where the $\boldsymbol{\delta}$ vectors represent the effects of flexibility. The result is a vector and may be expressed in any coordinate frame.

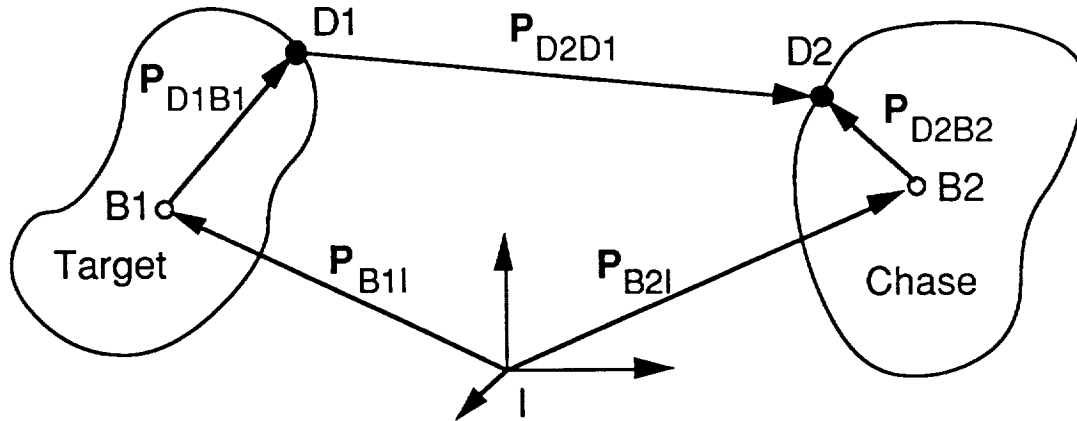


Figure 4.0-1 Relative Body Position Vector Definitions.

The relative velocity between docking nodes is defined as the velocity of the chase docking node as seen by an observer moving with the target docking node. Therefore, the derivative is taken wrt the target docking node frame. Equation (4.0-4) is the well known relationship for the time derivative of a vector viewed from two frames rotating at different angular rates.

$$(\dot{\mathbf{R}})_A = (\dot{\mathbf{R}})_B + \boldsymbol{\omega}_{BA} \times \mathbf{R} \quad (4.0-4)$$

where $(\dot{\mathbf{R}})_A$ = time derivative of an arbitrary vector \mathbf{R} wrt A frame

$(\dot{\mathbf{R}})_B$ = time derivative of the vector \mathbf{R} wrt B frame

$\boldsymbol{\omega}_{BA}$ = angular velocity of B frame wrt A frame.

Using the vector definitions in Fig. 4.0-1 and applying eqn. (4.0-4) results in the expression for the time derivative of the position of D2 wrt D1 taken wrt the D1 frame.

$$\begin{aligned}
 (\dot{\mathbf{P}}_{D2D1})_{D1} = & \dot{\mathbf{P}}_{B2I} - \dot{\mathbf{P}}_{B1I} + \boldsymbol{\omega}_{B2I} \times \left(\mathbf{P}_{D2_0B2} + \sum_j^{nm} \boldsymbol{\Phi}_{D2j} \eta_j \right) + \sum_j^{nm} \boldsymbol{\Phi}_{D2j} \dot{\eta}_j \\
 & - \boldsymbol{\omega}_{B1I} \times \left(\mathbf{P}_{D1_0B1} + \sum_j^{nm} \boldsymbol{\Phi}_{D1j} \eta_j \right) - \sum_j^{nm} \boldsymbol{\Phi}_{D1j} \dot{\eta}_j \\
 & - \left(\boldsymbol{\omega}_{B1I} + \sum_j^{nm} \boldsymbol{\Psi}_{D1j} \dot{\eta}_j + \boldsymbol{\omega}_{B1I} \times \sum_j^{nm} \boldsymbol{\Psi}_{D1j} \eta_j \right) \times \mathbf{P}_{D2D1}
 \end{aligned} \tag{4.0-5}$$

The deformation vectors are represented by their respective modal expansions in this expression.

The relative angular velocity of the chase docking frame wrt the target docking frame is given in eqn. (4.0-6).

$$\begin{aligned}
 \boldsymbol{\omega}_{D2D1} = & \boldsymbol{\omega}_{B2I} + \sum_j^{nm} \boldsymbol{\Psi}_{D2j} \dot{\eta}_j + \boldsymbol{\omega}_{B2I} \times \left(\sum_j^{nm} \boldsymbol{\Psi}_{D2j} \eta_j \right) - \boldsymbol{\omega}_{B1I} \\
 & - \sum_j^{nm} \boldsymbol{\Psi}_{D1j} \dot{\eta}_j - \boldsymbol{\omega}_{B1I} \times \left(\sum_j^{nm} \boldsymbol{\Psi}_{D1j} \eta_j \right)
 \end{aligned} \tag{4.0-6}$$

Again, the rotational deformation vectors are expressed as a modal expansion.

5.0 SIMULATION DESCRIPTION

The computer code which implements the MTB flexible body math model is coded in FORTRAN 77 and all floating point numbers are maintained with double precision accuracy. A listing of the computer code, along with input data files, is contained in Appendix A. The math model is coded in a modular form to facilitate debugging and modifications. The math model is controlled by the driver subroutine "DYNAMIC" which is called by the host simulation. This driver subroutine and the common declarations in the include file "INTFAC.INC" serve as the interface to the host simulation.

The following nomenclature is used for specific data extracted from coordinate vectors. Position vectors begin with the letter P. Translational velocity vectors begin with the letters PD. The vector $P_A_B_C$ is defined as the position vector from point B to point A expressed in the C frame coordinates; $PD_A_B_C$ is the time derivative of $P_A_B_C$. Angular velocity vectors begin with the letter W and derivatives follow the pattern described above. Transformation matrices begin with the letter T. The transformation T_A_B is a 3 x 3 matrix which transforms vectors from the B frame to the A frame. The labels used to define bodies, docking nodes, etc., in Figures 3.0-1 and 4.0-1 are also used in the simulation code.

5.1 Major Subroutines

Figure 5.1-1 is a flow diagram of the MTB two flex body math model. A general description of each of the major subroutines and include files is given in Table 5.1-1 and Table 5.1-2, respectively. The first pass through the model initializes the simulation variables. Data from the force/moment sensor is supplied to the subroutine "FMTRANS" through common declarations in the include file "INTFAC.INC". The relative motion data is made available to the host simulation through common declarations in the include file "INTFAC.INC".

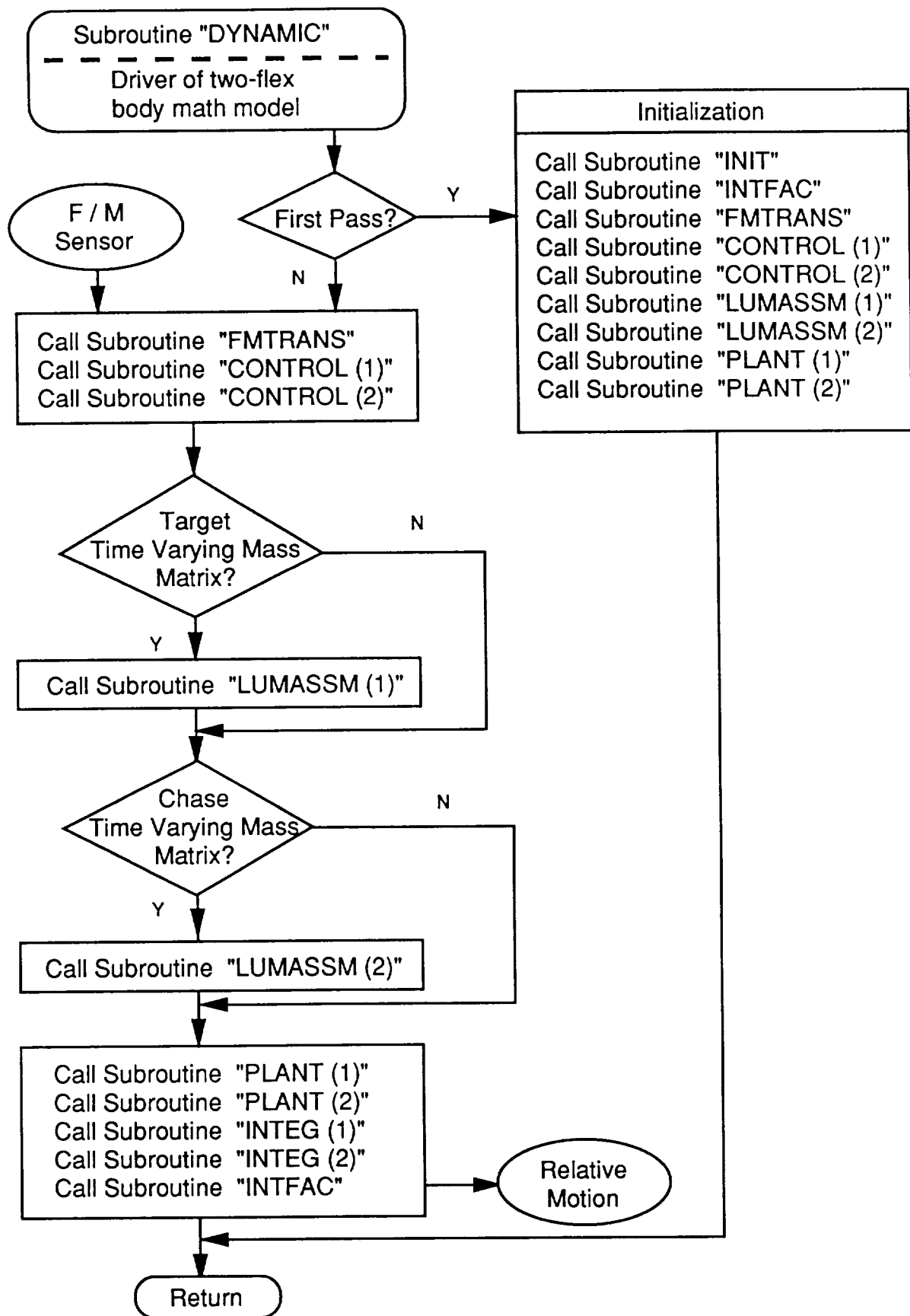


Figure 5.1-1 Two Flexible Body Math Model Flow Diagram.

Table 5.1-1 Two Flex Body Math Model Subroutines.

<u>Subroutine</u>	<u>Description</u>
DYNAMIC	Driver of two-flex body dynamics model; called by main routine.
INIT	Initialization of mass properties, geometry, state vectors, flexibility data, etc.
FMTRANS	Computes contact forces/torques at docking port nodes based on force/torque sensor data.
CONTROL (NB)	Computes control forces/torques applied to NB th body.
LUMASSM (NB)	Computes mass matrix of NB th body in lower-upper decomposed form.
PLANT (NB)	Computes right hand side terms (e.g., damping, stiffness, forcing functions, etc.) of equations of motion and solves $\mathbf{A} \mathbf{x} = \mathbf{b}$ for NB th body.
INTEG (NB)	Integrates derivatives of quasi and generalized coordinate vectors and updates generalized coordinate vector derivative for NB th body.
INTFAC	Updates relative transformations, position vectors and rate vectors with flexibility included.

Table 5.1-2 Two Flex Body Simulation Include Files.

<u>INCLUDE FILE</u>	<u>DESCRIPTION</u>
BODY.INC	Variable/parameter declarations and common specifications for mass properties, coordinate vectors, flexibility, etc. for both bodies; also simulation control variables.
INTFAC.INC	Variable declarations and common specifications for interfacing to main routine; includes all necessary relative motion data.

5.2 Input Data

The MTB flex body simulation has three ASCII input data files. Table 5.2-1 gives a general description of the contents of these input data files. A computer listing of example input data files is contained in Appendix A.

The ASCII input data may be modified using any available editor (e.g., the VI editor on the UNIX based ALLIANT computer system). The existing input files contain comments (see example files in Appendix A) which help to define the input data. Most of the input data is well described by the comments. The EOM complexity options are defined with the following acronyms:

<u>EOM Option</u>	<u>Description</u>
EHOMIT	Eliminate higher order modal integral terms
ERNLT	Eliminate remaining right hand side nonlinear terms Note: Enabling ENRLT encompasses EHOMIT regardless of the EHOMIT option status.
FREE-FREE	Free-Free modes

The EOM complexity options are declared as logical variables. Therefore, ".T." is TRUE and enables the option; ".F." is FALSE and disables the option.

The pre-processed flex data (see Table 5.2-1) forms the bulk of the required flexibility data and is generated by the flexible body data pre-processing algorithm discussed in Section 5.3. The remaining body data can be generated using a file editor and then concatenated with the output of the pre-processing algorithm to produce the whole input data file.

Currently, all the input data specified in Table 5.2-1 is read regardless of the EOM options. For example, enabling the EHOM-IT option does not eliminate the need for higher order modal integral data in the respective body input data file even though the data is not used in the simulation. This follows from the fact that the

Table 5.2-1 Two Flex Body ASCII Input Data Files.

<u>INPUT FILE</u>	<u>DESCRIPTION</u>
MTBSIM.INP	Simulation control data: <ul style="list-style-type: none"> Integration step size Screen output step size Data file output step size Stop time Integration method flag
BODY1.INP	Target body data: <ul style="list-style-type: none"> Mass Inertia matrix Initial P_B_I_I Initial PD_B_I_I Initial w_B_I_B [I B] Euler rotation sequence [I B] Euler rotation angles Number of nodes; dock node num Node geometry *P_D1_S1_S1 [B1 D1o] Transformation *[D1 S1] Transformation EOM complexity options Flexibility Data: <ul style="list-style-type: none"> Number of flexible modes Mode frequencies Modal damping Initial generalized modal coords. Initial gen modal coord derivatives <hr/> Pre-Processed Flex Data: <ul style="list-style-type: none"> Generalized mass matrix Translational modal gains Rotational modal gains Standard modal integrals Higher order modal integrals
BODY2.INP	Chase body data: <ul style="list-style-type: none"> Same as target body data omitting *

flexible body data pre-processing routine (Section 5.3) computes and outputs all modal integral terms. An advantage of this approach is that the input data files (for a given body) have a "standard" structure (i.e., same contents).

5.3 Flexible Body Data Pre-Processing Algorithm

Flexible body simulations require a great deal more input data than pure rigid body simulations. Typically, the simulation input data (e.g., modal integrals and generalized mass matrix) is computed from the output of a finite element model (FEM). To assist users of the MTB flexible body simulation, Control Dynamics coded a flexible body data pre-processing algorithm. This algorithm acts as a post-processor for FEM data and a pre-processor for the MTB flexible body simulation. Figure 5.3-1 is a block diagram which illustrates the role of the algorithm.

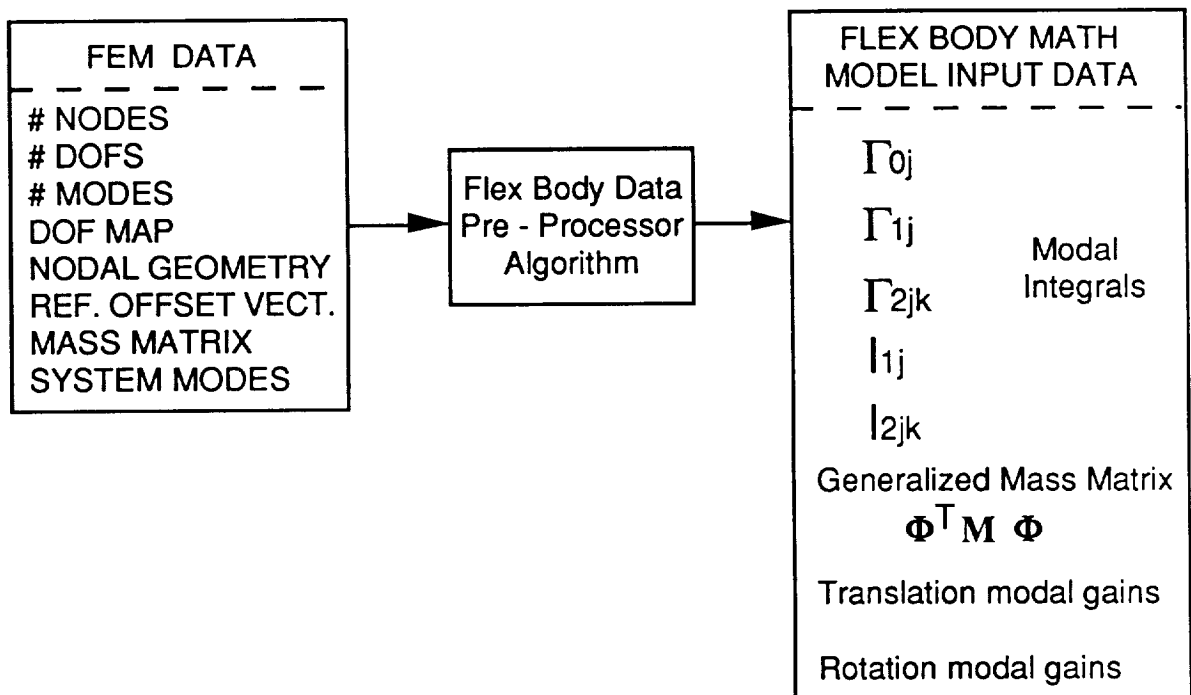


Figure 5.3-1 Flexible Body Data Pre-Processing Algorithm.

Table 5.3-1 describes the flexible body data required in the ASCII input file for the pre-processing algorithm. This data is read using free format (see pre-processor code listing in Appendix C). The output (see flexibility data in Table 5.2-1) is written to the user specified file in a format consistent with the inputs to the MTB math model simulation. The output file from the pre-processor contains the data denoted as pre-processed flex data in Table 5.2-1.

Table 5.3-1 Flexible Body Data Pre-Processor Input.

<u>Input Description</u>	<u>Explicit Input Data</u>
Degree of Freedom Map	# nodes, # DOF, # modes Blank line node num (i), DOF type (i) Blank line
Node Geometry Table	*Offset vector Node position vector (i) Blank line
Mass Matrix	Mass_mat (# DOF, # DOF) Blank line
Modal Matrix	Phi (# DOF, # modes)
*A displacement vector subtracted from each node position vector which allows the body reference point to be relocated.	

6.0 SIMULATION VERIFICATION

The math model simulation was verified in a series of progressive steps on a PC-386 computer at Control Dynamics Company prior to installation on the host computer . This section of the report documents the various tests used to verify the math model simulation. Time domain results for all of the various simulation component tests are not presented; only the results from an "all-encompassing" two flexible body comparison to the multi-flex body code TREETOPS is shown. The "full-up" flexible body verification run presented in section 6.4 was repeated after the model was installed on the host computer to ensure proper operation of the model in the MTB facility. This "full-up" run exercises all aspects of the simulation and, therefore, eliminates the need to reverify the simulation in a component manner on the host computer.

6.1 Equations of Motion Solution Verification

The LU decomposition technique used to solve for the derivative of the quasi-coordinate vector is described in Section 2.5 of this report. Since the LU decomposition and forward/backward substitution algorithms were coded to solve a general linear set of equations $\mathbf{A} \mathbf{x} = \mathbf{b}$ (i.e., the symmetry property of the mass matrix was not considered), the algorithm may be verified using any arbitrary deterministic system with a square matrix \mathbf{A} . The LU decomposition algorithms were used to solve several linear systems of various size. The same systems were then solved using the well established computer analysis tool called MATLAB. The solution vectors (\mathbf{x}) were compared to ensure proper operation of the EOM solution technique.

6.2 Rigid Body Verification

The rigid body portion of the math model was tested by two means. First, each degree of freedom (i.e., 3 translations and 3 rotations per body) was tested independently through the application of a constant force/moment. The results were verified by hand calculations.

Secondly, the simple two rigid body system shown in Figure 6.2-1 was used to test portions of the relative vehicle motion calculations and the contact force/moment transformation. For the system in Fig. 6.2-1, an omnidirectional spring of stiffness K and undeformed length r was assumed at the target docking node in order to simulate contact forces. Various initial conditions on the chase vehicle were used to induce contact between the docking ports in each translational axis. The resulting rigid body motion was verified through hand calculations. Also, forces and torques representing contact were input as "sensed" values to simulate data from the F/M sensor in the MTB facility. The math model operated on these "sensed" values to produce contact forces and torques at the body docking nodes. Again, the results were validated using hand calculations.

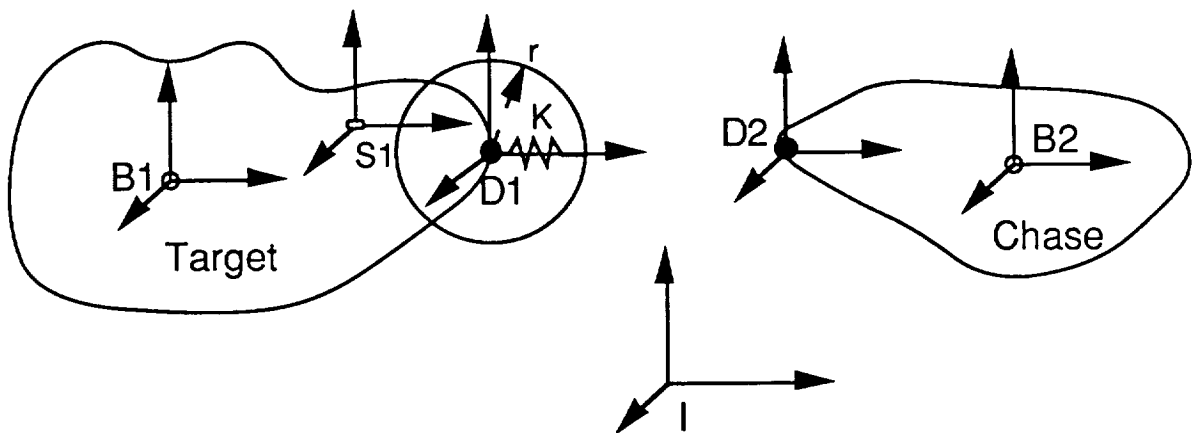
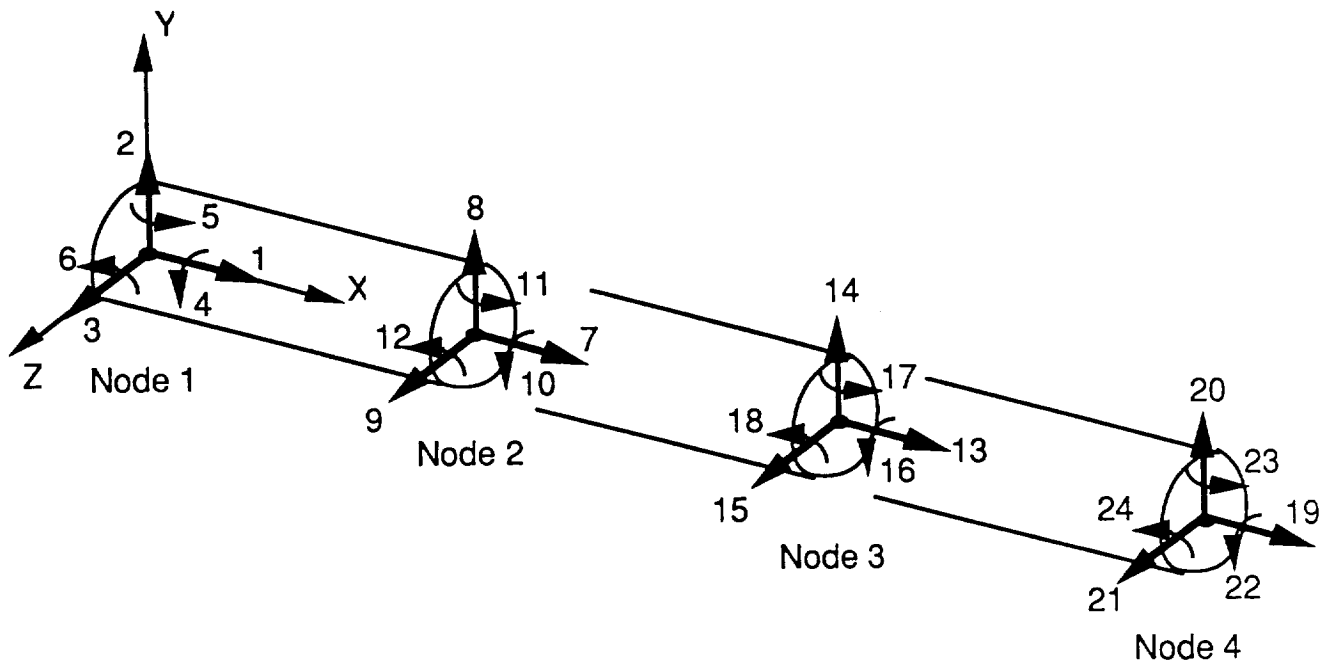


Figure 6.2-1 Two Rigid Body Contact Test Case.

6.3 Flexible Beam Model

The flexible body math model was tested using a free-free flexible beam model created in the structural computer code SMIS (Structures and Matrix Interpretive System). The beam geometry and properties are shown in Figure 6.3-1.



Length = 24.78 m

Radius = 0.1 m

Density = 8410 kg/m³

Elasticity = 100 x 10⁹ Nt/m²

Rigidity = 30 x 10⁹ Nt/m²

$$ICM = \begin{bmatrix} 32.74 & 0 & 0 \\ 0 & 335K & 0 \\ 0 & 0 & 335K \end{bmatrix} \text{ kg} \cdot \text{m}^2$$

Figure 6.3-1 Properties of Free-Free Flexible Test Beam.

The beam was sized to yield a first bending mode near 1 Hz. Table 6.3-1 lists the first four flexible mode frequencies from the SMIS model of the beam described in Figure 6.3-1.

Table 6.3-1 First Four Bending Modes of Free-Free Flexible Test Beam.

<u>Mode Number</u>	<u>Frequency (Hz)</u>	<u>Description</u>
1	1.003	Z Translation/Y Rotation
2	1.003	Y Translation/Z Rotation
3	2.770	Z Translation/Y Rotation
4	2.770	Y Translation/Z Rotation

The SMIS flexible beam model (i.e., mass matrix, mode shapes, DOF map, node geometry, etc.) was processed by the flexible body data processing algorithm discussed in Section 5.3. The processor produced the generalized mass matrix, modal gains, and modal integrals required for the MTB flexible body input data files (See Section 5.2).

6.4 Flexible Body Verification

The flexible beam model was implemented in two other flexible body simulations for comparisons to the MTB math model. Table 6.4-1 summarizes the important differences among the simulations. First, the model was implemented in an in-house (CDy) flex body simulation based on Lagrange's quasi-coordinate EOM (see Section 2.0), but coded independently. This simulation comparison facilitated rapid debugging since the terms in the equations of motion had a "one-to-one" correspondence. Secondly, the flexible beam model was implemented in the multi-flex body simulation called TREETOPS. TREETOPS is based on Kane's equations and gives a totally independent verification of the MTB math model. The majority of the data in TREETOPS is maintained as single precision. Notice also (from Table 6.4-1) that both test simulations use a different integration method than the MTB simulation. Only a comparison between the MTB math model and

TREETOPS is presented in this report.

Table 6.4-1 Simulation Comparison.

<u>Simulation</u>	<u>Precision</u>	<u>Integration Method(s)</u>
MTB Math Model	Double	HWIL Schemes
CDy In-House	Double	4-Pass Runge-Kutta
TREETOPS	Single	4-Pass Runge-Kutta

Several flexible body test cases were investigated using different "contact" forcing functions and body configurations. Only one "full-up" two flexible body comparison to TREETOPS is presented here; this run exercises "all" features of the MTB flexible body math model.

Figure 6.4-1 illustrates the system configuration and contact force/moment application points for the flexible body verification run. Equations (6.4-1) and (6.4-2) define the "contact" forcing function applied to the target docking node (D1). Notice

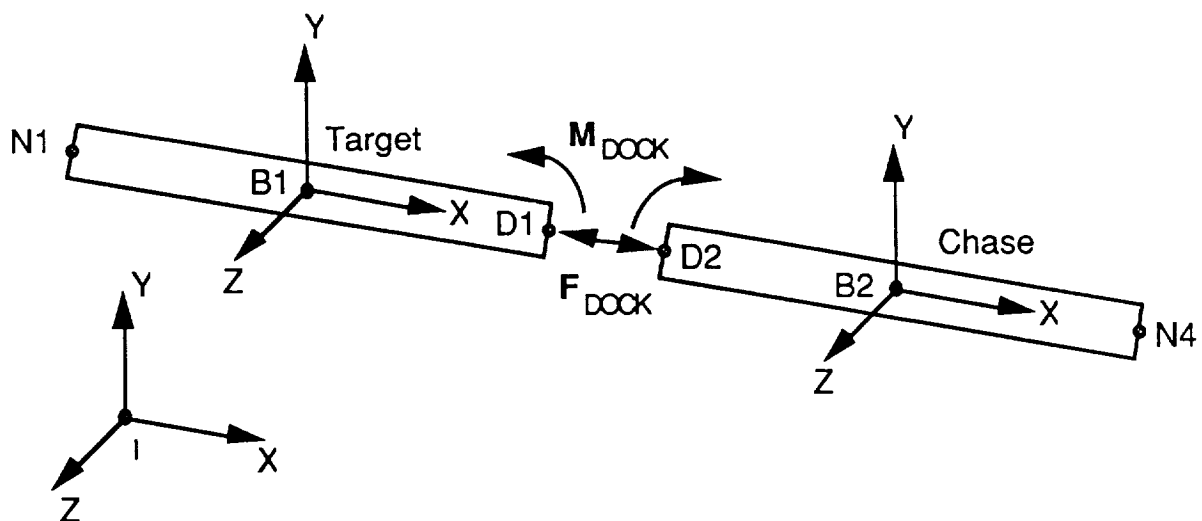


Figure 6.4-1 Two Flexible Beam Test Case.

$$\begin{Bmatrix} \mathbf{F}_{\text{DOCK}} \\ \mathbf{M}_{\text{DOCK}} \end{Bmatrix} = \begin{Bmatrix} \begin{pmatrix} -3K \\ 2K \\ 0 \end{pmatrix} \text{ Nt.} \\ \begin{pmatrix} 0 \\ 15K \\ 0 \end{pmatrix} \text{ Nt-m.} \end{Bmatrix} \quad 0 \leq t < 0.02 \text{ sec.} \quad (6.4-1)$$

$$\begin{Bmatrix} \mathbf{F}_{\text{DOCK}} \\ \mathbf{M}_{\text{DOCK}} \end{Bmatrix} = \begin{Bmatrix} \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix} \text{ Nt.} \\ \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix} \text{ Nt-m.} \end{Bmatrix} \quad t \geq 0.02 \text{ sec.} \quad (6.4-2)$$

that the "contact" forcing function is applied over a short time period (impulsive loading) which is characteristic of actual contact.

Figure 6.4-2 is a TREETOPS block representation of the two flexible beam test case shown in Figure 6.4-1 with the forcing functions defined in eqns. (6.4-1) and (6.4-2). The function generators are step functions which produce the forcing functions; position and negative step functions are used to "turn-on" and "turn-off" the forces and moments. Actuators 1, 2, 3, and 4 are reaction jets which apply a force at a given node; these actuators implement the equal and opposite "contact" forces at the body docking nodes. Actuators 5 and 6 are moments actuators which implement the equal and opposite "contact" moments. In TREETOPS, the axis of application of these actuators changes in the body fixed frame (i.e., undeformed body frame) due to body flexing effects. Body 1 and 2 are both three mode flexible models of the beam shown in Figure 6.3-1. Hinges 1 and 2 are configured with zero translational and rotational stiffness to allow the bodies to respond as "free-floaters" (i.e., unconstrained bodies). Sensor 1 is an inertial position vector sensor which outputs the position of the chase docking node (D2) with respect to the target docking node (D1). Sensor 2 is an inertial velocity sensor which outputs the velocity of D2 with respect to D1. The TREETOPS simulation used Runge-Kutta

integration with a step size of 0.01 seconds; data was also written out every 0.01 seconds.

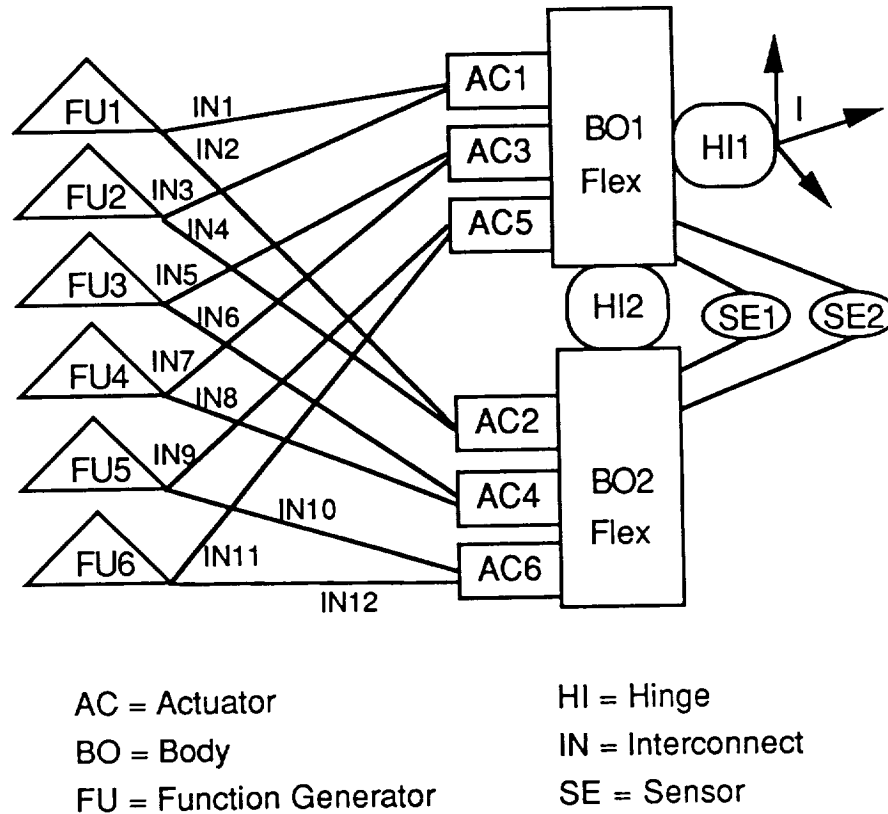


Figure 6.4-2 TREETOPS Block Representation of Two Flexible Beam Test Case.

The "contact" forcing function defined in eqns (6.4-1) and (6.4-2) was implemented in the MTB math model by computing an equivalent "sensed" force and moment (i.e., at some arbitrary F/M sensor location on the target body). This "sensed" force and moment was then used as if the data came from a physical sensing device; the "sensed" force and moment was transformed to the appropriate body docking nodes and applied as an external forcing function. The MTB was initialized such that the docking nodes were collocated at time zero. Referring to Section 3.0 of this report, the "contact" forces will be equal and opposite as in the

TREETOPS simulation. However, the "contact" moments will be equal and opposite only for collocated docking nodes; the TREETOPS simulation of this two flexible beam problem will apply equal and opposite "contact" moments over the entire duration of the forcing function. The results should still be comparable since the "contact" forcing function is of short duration; thus, the relative displacement of the docking nodes should be small during the forcing period.

In the MTB simulation, both the target and chase bodies are three mode flexible models of the beam shown in Figure 6.3-1. The modal damping was zero and the higher order modal integral terms were neglected; TREETOPS has no direct way of including non-zero modal damping and requires the use of a COSMIC NASTRAN FEM model to include higher order modal integrals. The MTB simulation was configured to use Adams (3-1)/2 integrations with a step size of 0.002 seconds. The smaller step size (factor of 5 lower than TREETOPS step size) should give results comparable to the multi-pass Runge-Kutta scheme.

The results of the MTB and TREETOPS simulation runs will now be presented. Since both bodies are based on the same generic flexible body model, the presented data focuses on the target body and the relative motion data.

Figure 6.4-3 shows the time response of the "contact" forcing functions applied to the target docking node. These correspond to equations (6.4-1) and (6.4-2). Figure 6.4-4 contains both the MTB and TREETOPS results for the target body center of mass (CM) velocity vector expressed in the inertial frame. The corresponding curves are plotted with the same line style and the plot scale does not allow the observer to differentiate between comparable curves. Figure 6.4-5 shows the vector difference between the TREETOPS and the MTB velocity vectors for the target body. The remaining data will be presented in the format of Figures 6.4-4 and 6.4-5; that is, the comparable MTB and TREETOPS data (usually a 3 element vector) will be plotted on the same graph at the top of the page and the

vector difference (TREETOPS - MTB) will be shown in the next figure at the bottom of the page.

Figures 6.4-6 and 6.4-7 contain results for the target body CM position expressed in the inertial frame. Figures 6.4-8 and 6.4-9 show the target body angular velocity vector data. There is significant difference between the MTB and the TREETOPS results for the angular velocity about the body X-axis. The moment arm associated with the torque about the X-axis which produces this angular velocity arises solely from body flexing. The "contact" moment about the body y-axis displaces the target docking node in the negative Z direction. The "contact" force along the positive Y direction, in conjunction with the Z displacement, produces the torque about the X-axis. Investigations into the different results showed that the X-axis angular velocity is sensitive to modal damping; this follows since the moment arises solely from flexibility. In turn, the X-axis angular velocity is also sensitive to the method of integration; different numerical integration methods induce different levels of numerical damping. For example, repeating the MTB simulation run using rectangular integration produces an X-axis angular velocity value much closer to the TREETOPS result. The conclusion is that the difference in integration methods, combined with the effects of single versus double precision and the sensitivity of the X-axis angular velocity in this particular case, is the source of the differing results.

Figures 6.4-10 and 6.4-11 contain the generalized modal coordinate derivative data for the target body. Figures 6.4-12 and 6.4-13 show the generalized modal coordinate data for the target body. Figures 6.4-14 and 6.4-15 show the translational deformation data associated with the target docking node. This data is expressed in the target body frame and represents deformation with respect to the rigid body. Figures 6.4-16 and 6.4-17 contain data for the rotational deformation of the target docking node with respect to the undeformed position.

Figures 6.4-18 and 6.4-19 show the results for the relative translational velocity of the chase docking node with respect to the target docking node expressed in the inertial frame. Figures 6.4-20 and 6.4-21 contain data for the relative docking node position vector (D2 with respect to D1), expressed in the inertial frame. Notice in general that the difference between the MTB and TREETOPS results appears to be growing with increasing time. This can be attributed to the difference in numerical damping associated with the respective integration techniques and the effects of single versus double precision simulation variables.

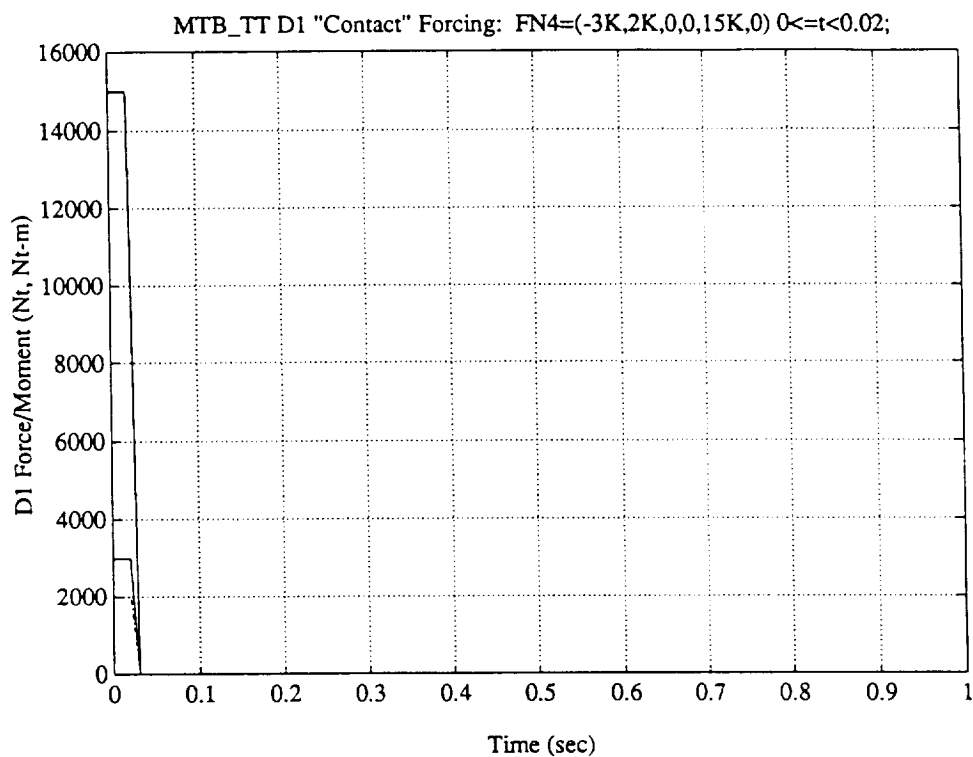


Figure 6.4-3 Target Docking Node "Contact" Forcing Functions.

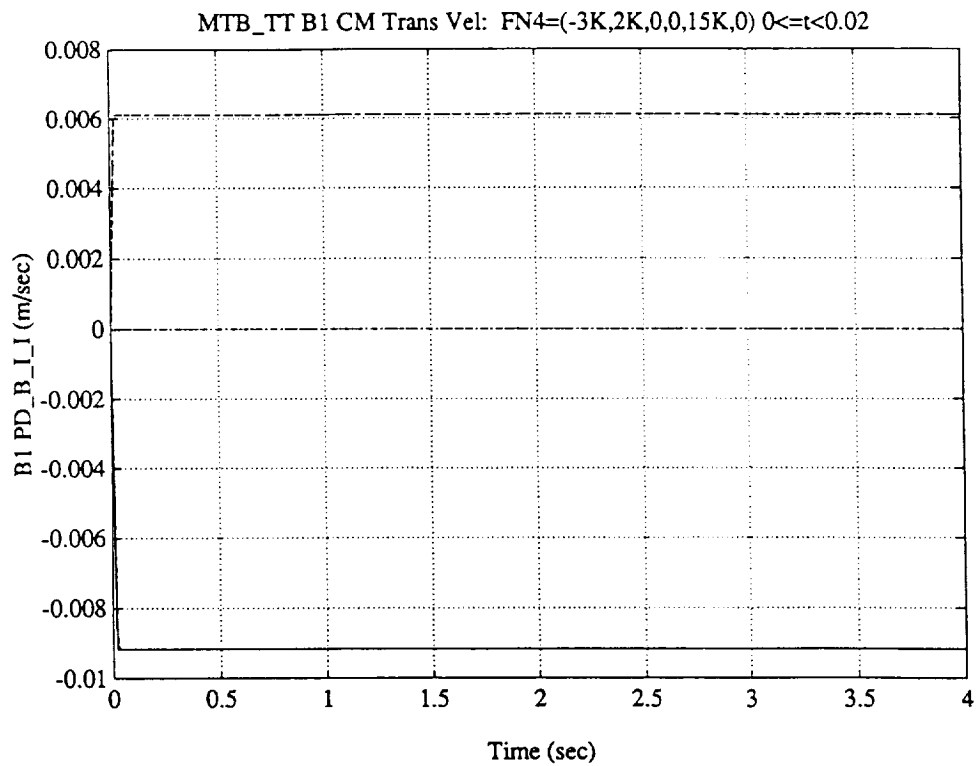


Figure 6.4-4 Target Body CM Translational Velocity.

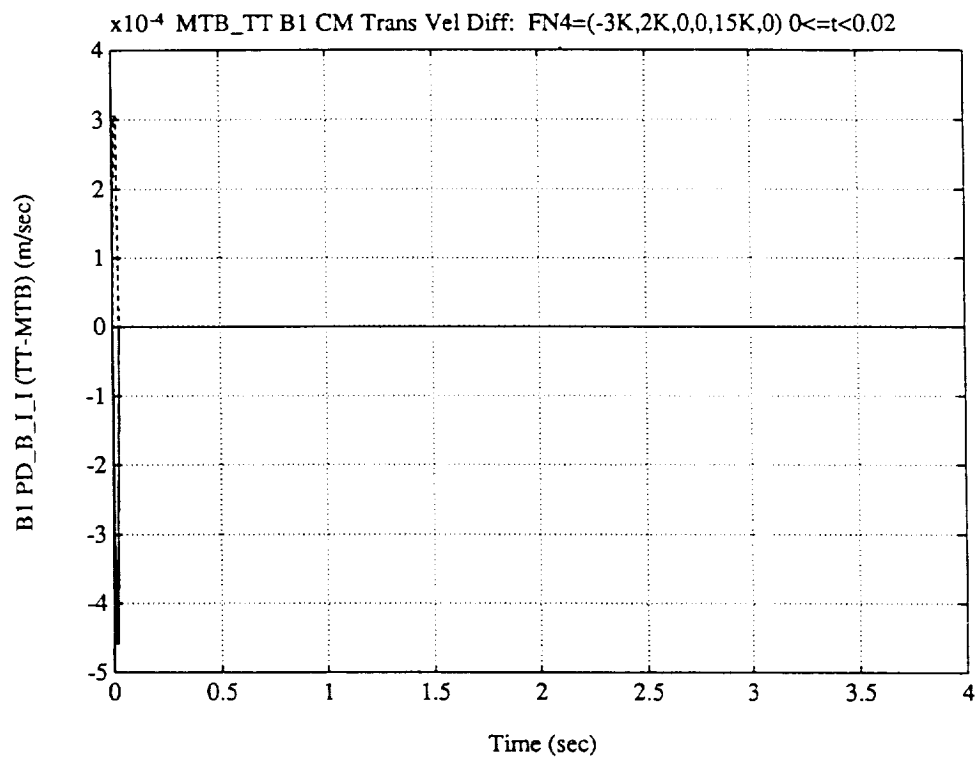


Figure 6.4-5 TREETOPS - MTB Target Body CM Translational Velocity.

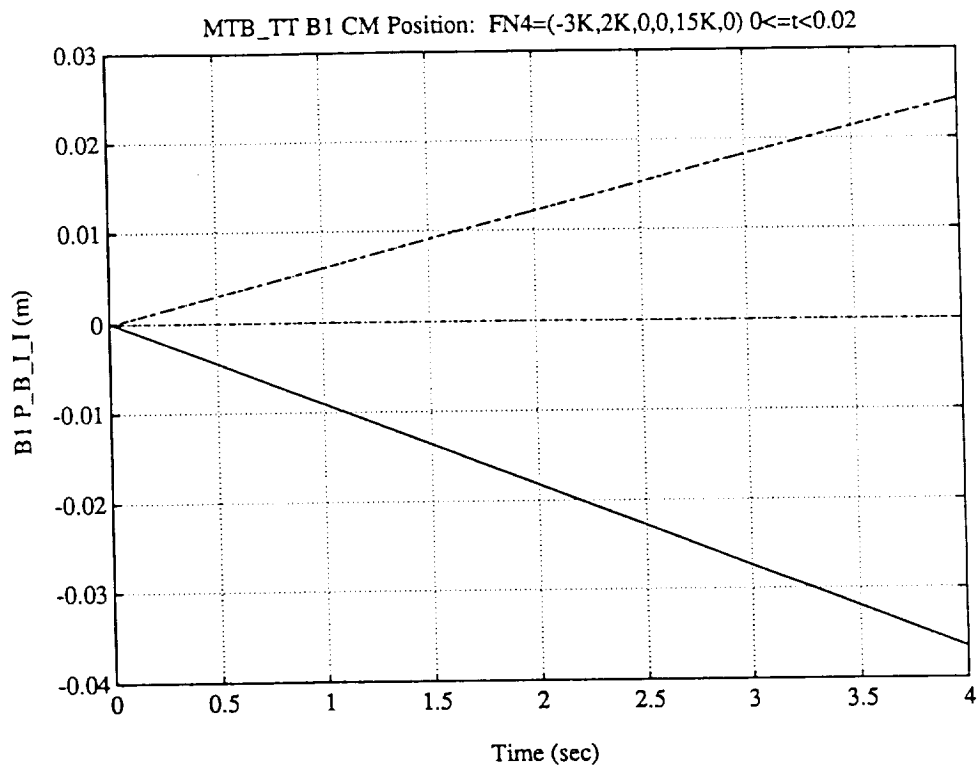


Figure 6.4-6 Target Body CM Position.

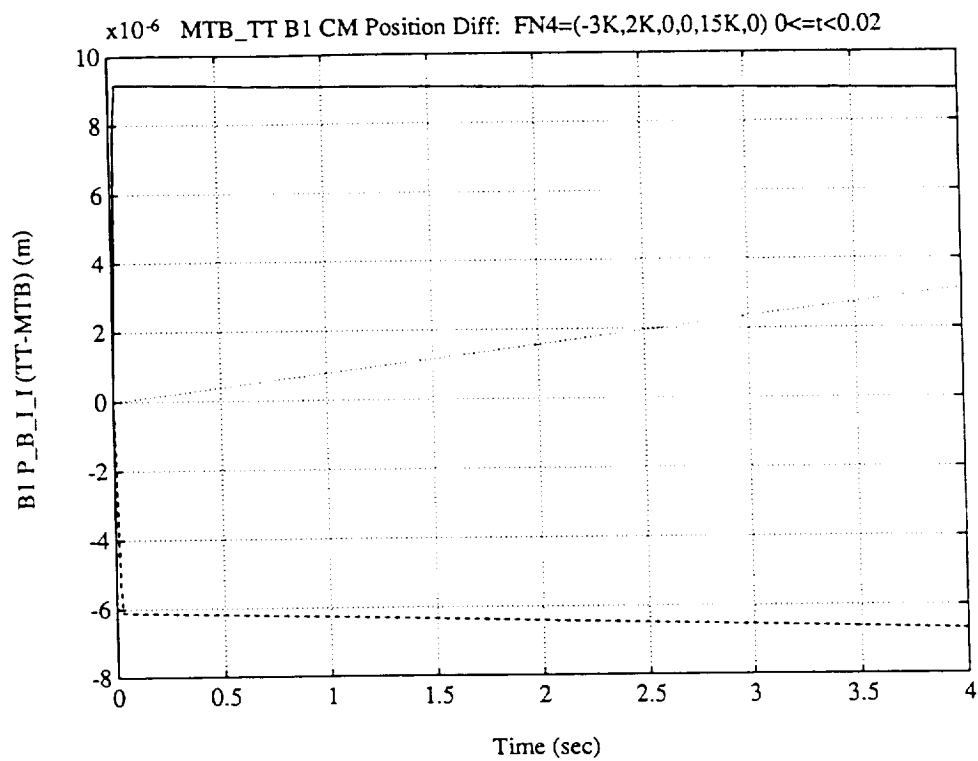


Figure 6.4-7 TREETOPS - MTB Target Body CM Position.

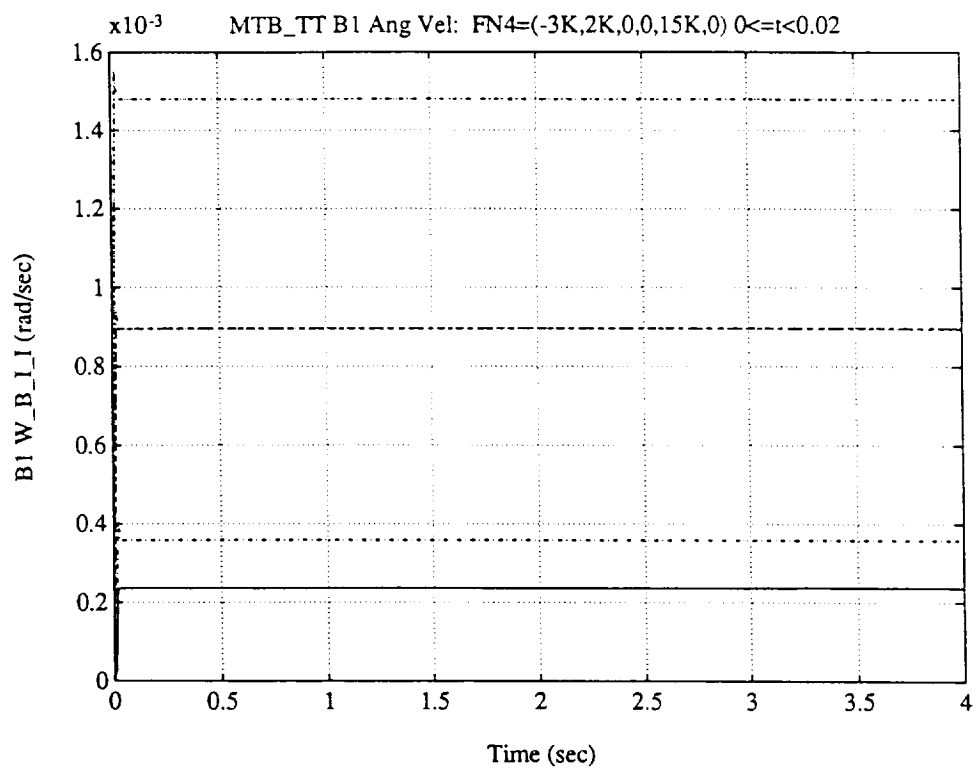


Figure 6.4-8 Target Body Angular Velocity.

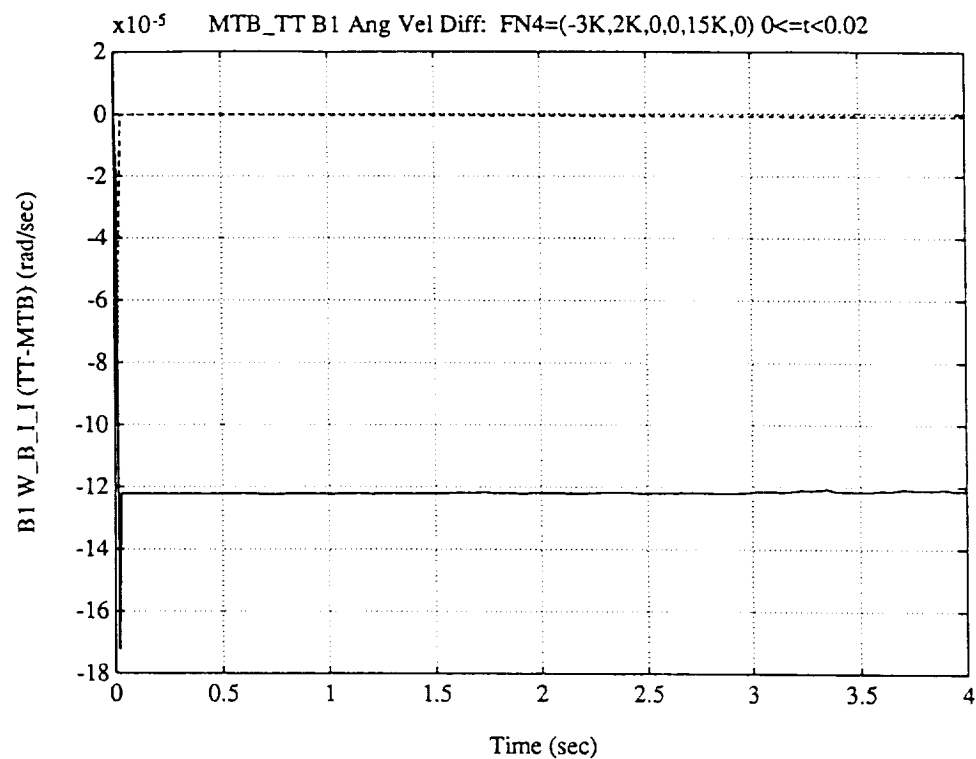


Figure 6.4-9 TREETOPS - MTB Target Body Angular Velocity.

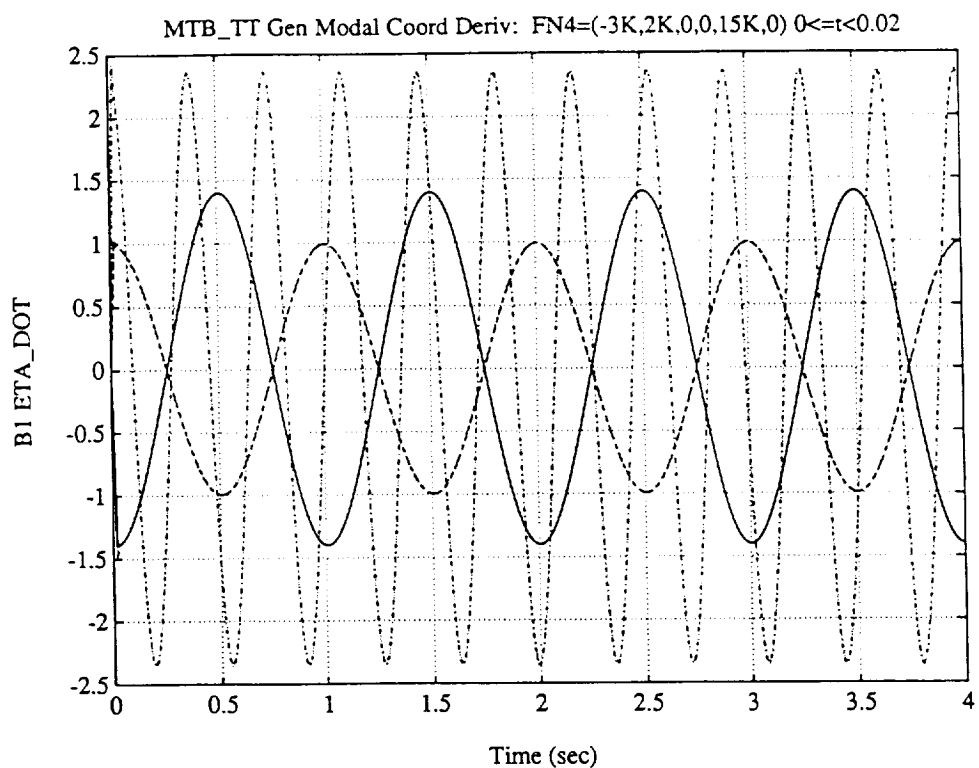


Figure 6.4-10 Target Body Generalized Modal Coordinate Derivatives.

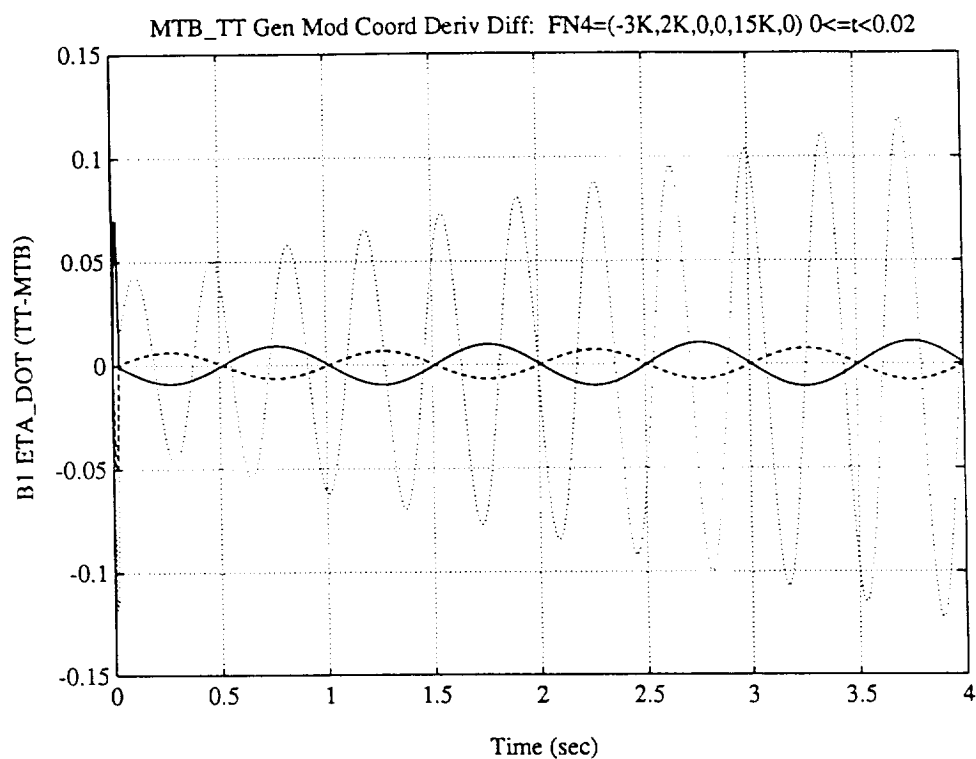


Figure 6.4-11 TREETOPS - MTB Target Body Generalized Modal Coordinate Derivatives.

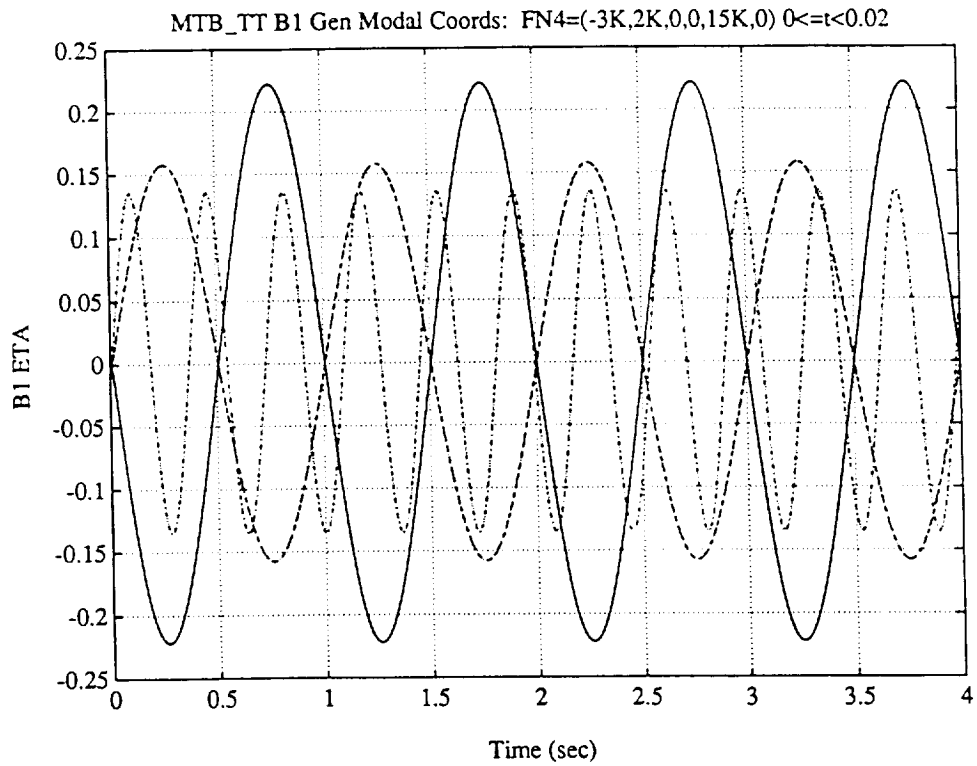


Figure 6.4-12 Target Body Generalized Modal Coordinates.

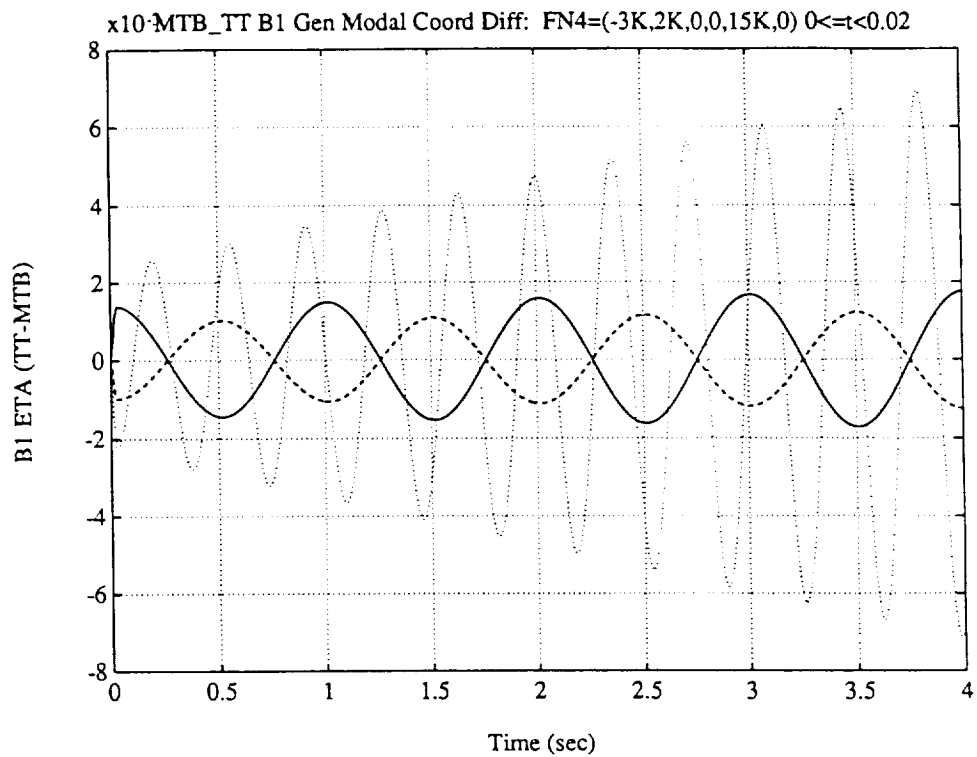


Figure 6.4-13 TREETOPS - MTB Target Body Generalized Modal Coordinates.

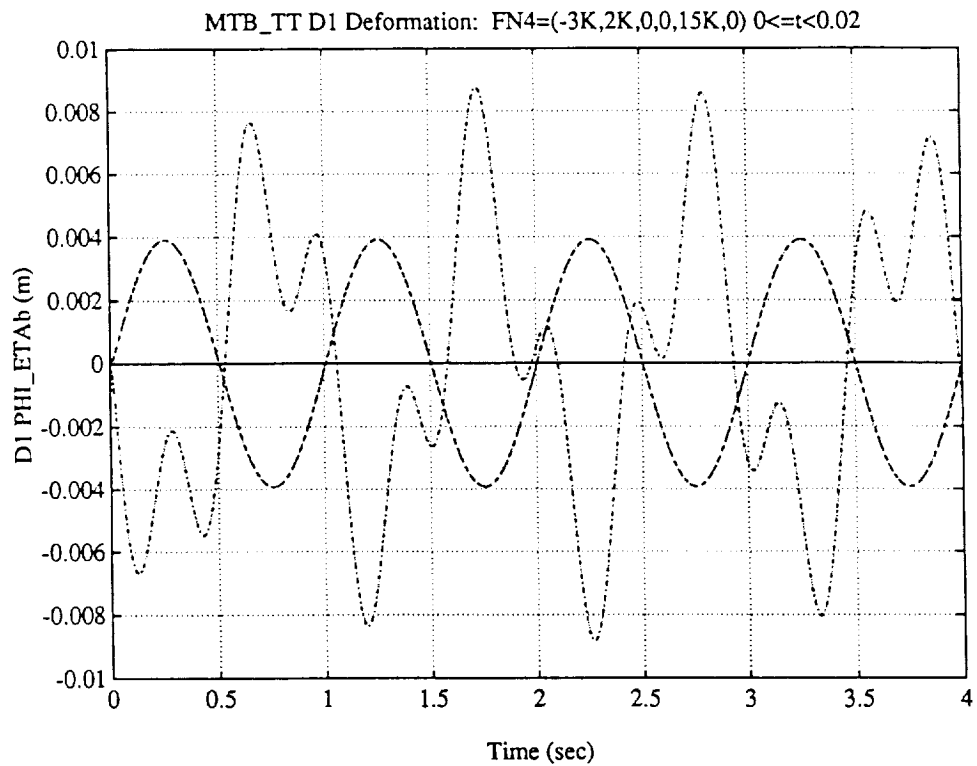


Figure 6.4-14 Target Body Docking Node Translational Deformation.

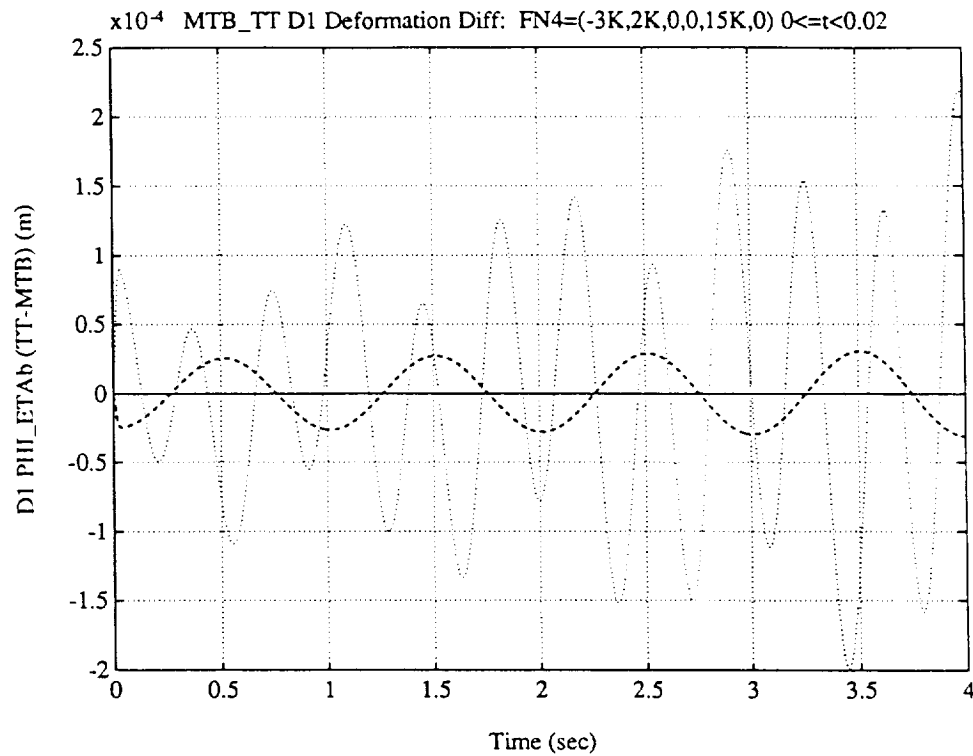


Figure 6.4-15 TREETOPS - MTB Target Body Docking Node Translational Deformation.

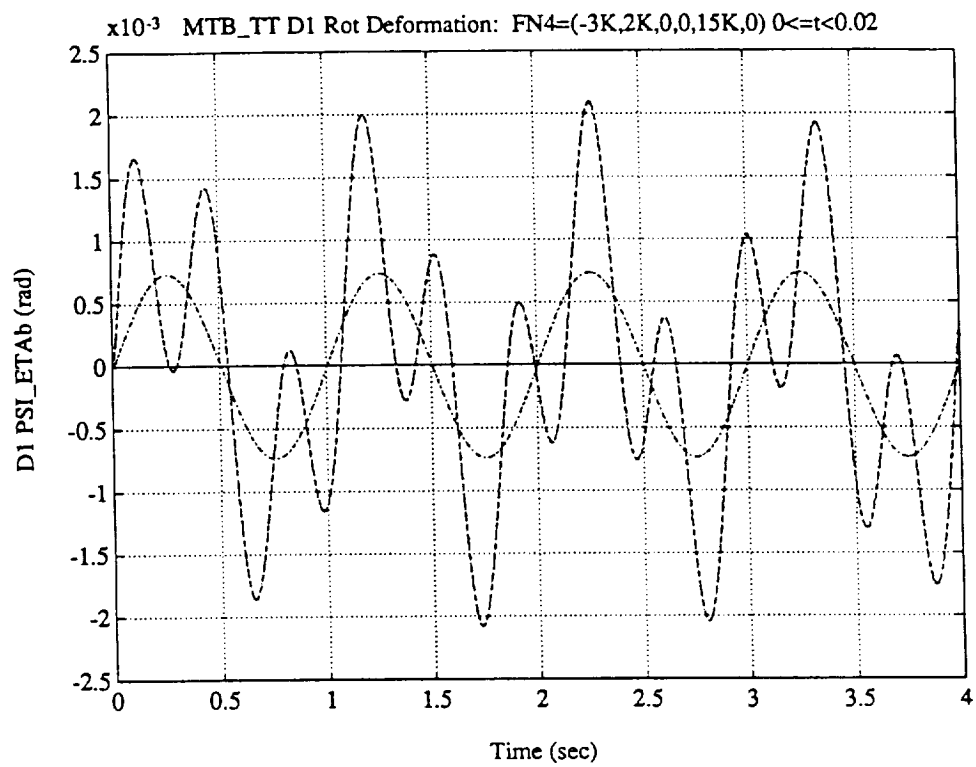


Figure 6.4-16 Target Body Docking Node Rotational Deformation.

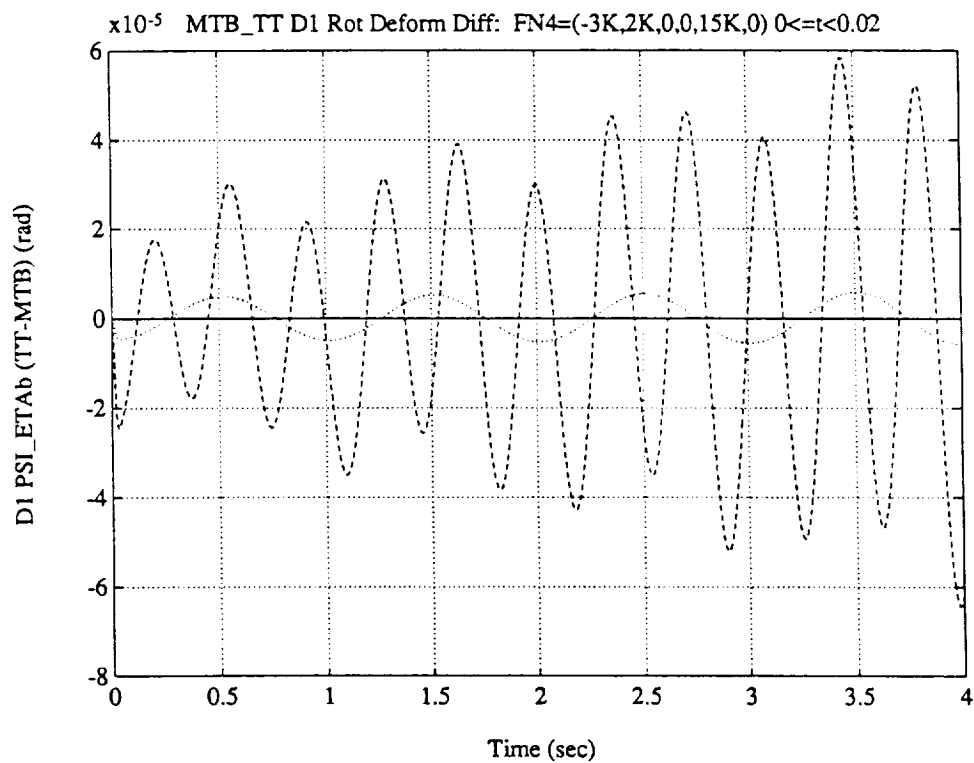


Figure 6.4-17 TREETOPS - MTB Target Body Docking Node Rotational Deformation.

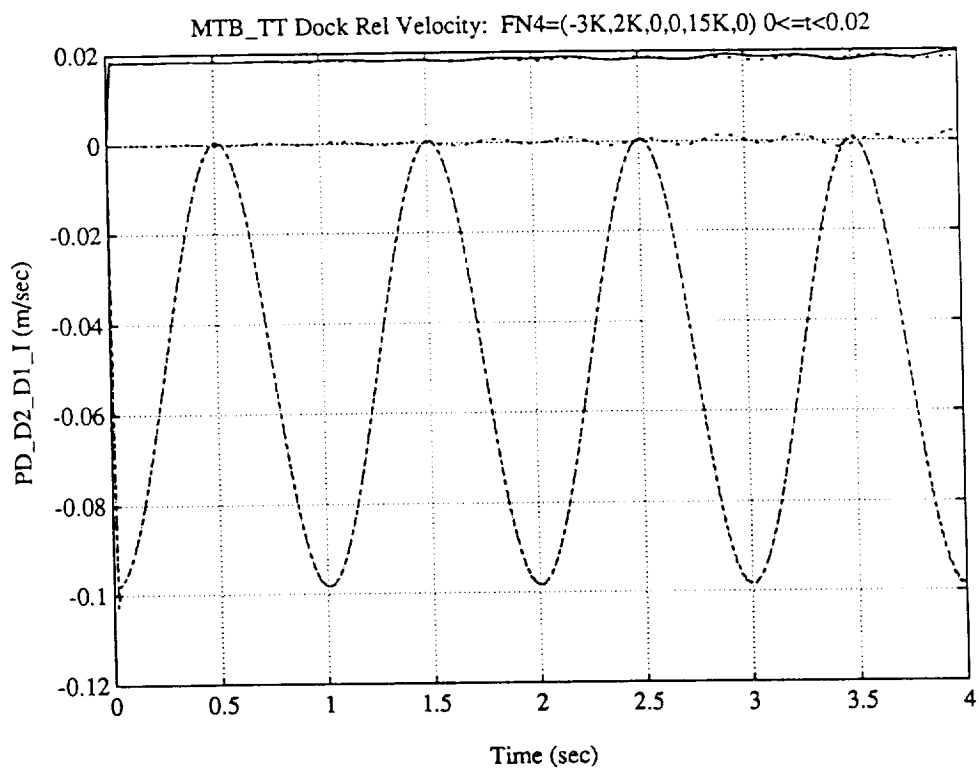


Figure 6.4-18 Docking Node Relative Translational Velocity.

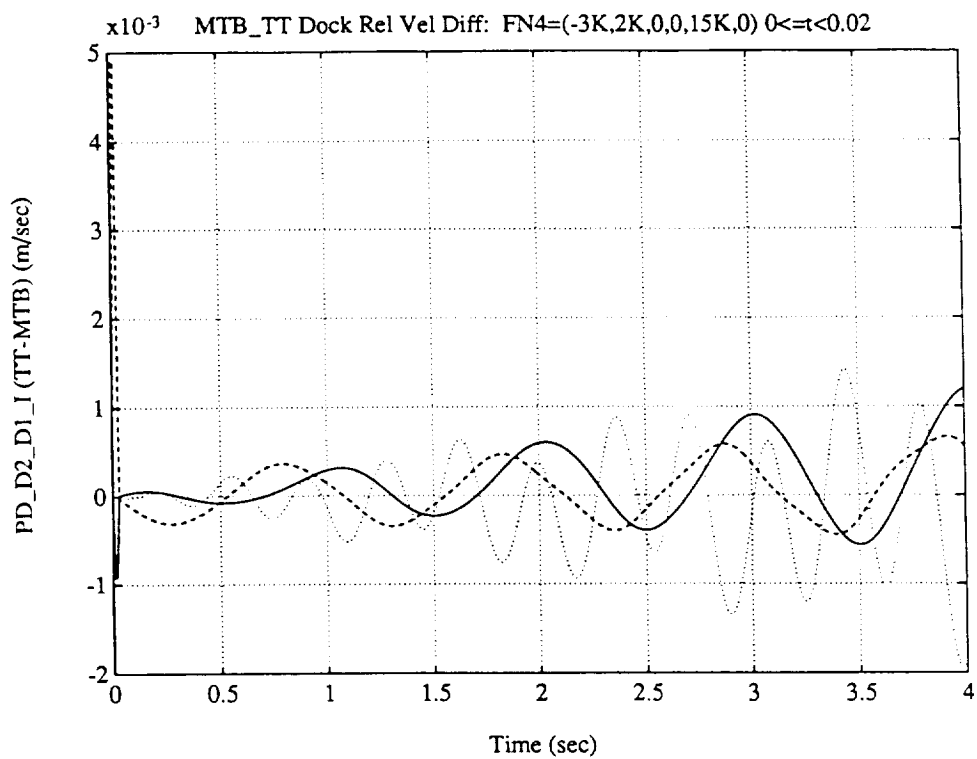


Figure 6.4-19 TREETOPS - MTB Docking Node Relative Translational Velocity.

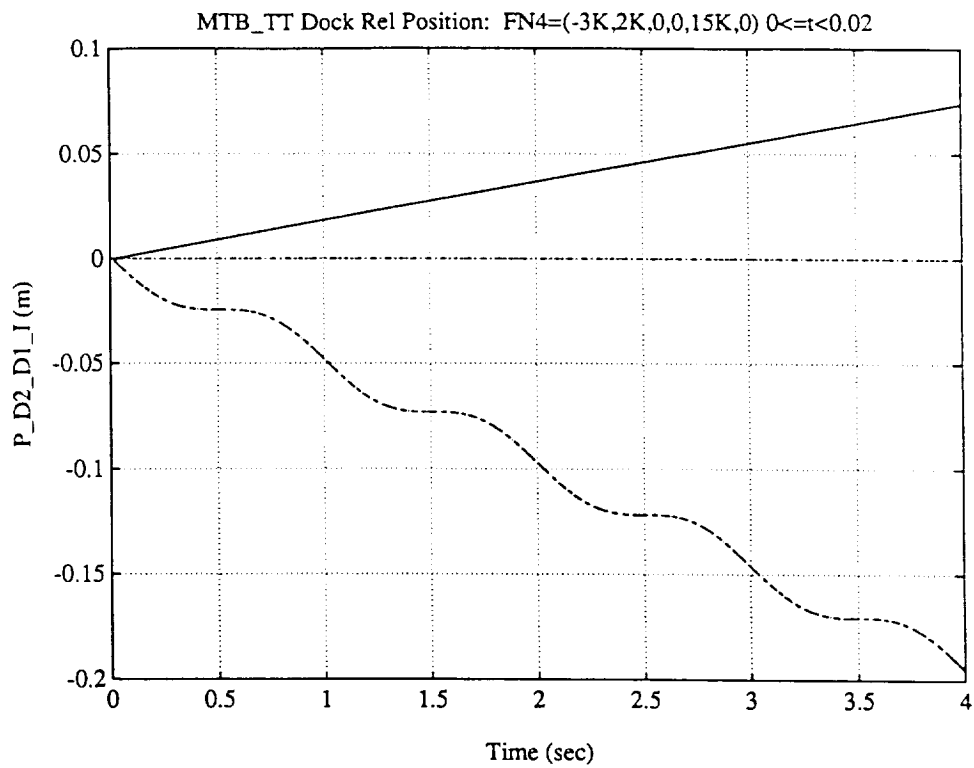


Figure 6.4-20 Docking Node Relative Position.

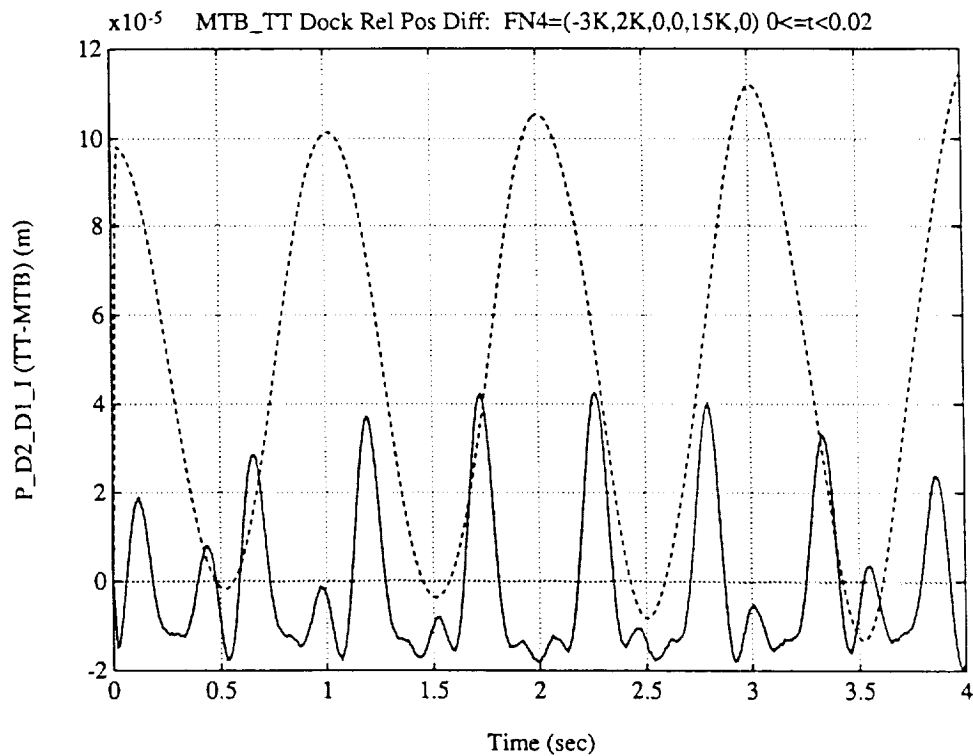


Figure 6.4-21 TREETOPS - MTB Docking Node Relative Position.

7.0 SUMMARY AND CONCLUSIONS

Per the contract objective, a generalized flexible body math model has been developed to further enhance the MTB Facility capabilities. The math model assumes that the flexible modes which characterize body bending will be orthonormal modes. This assumption allows the bending modes to be characterized by the standard second order equation

$$\ddot{\eta}_i + 2 \zeta_i \omega_i \dot{\eta}_i + \omega_i^2 \eta_i = F_i$$

as requested in the statement of work. Only minor modifications to the math model would be required to handle other less common mode types (e.g., non-orthogonal modes which typically result from sub-structure coupling techniques such as Craig-Bampton component mode synthesis).

The model was coded in FORTRAN 77 utilizing double precision floating point numbers. The math model code offers the user several levels of complexity for the equations of motion and three hardware-in-the-loop numerical integration options. Due to the increased volume of data over pure rigid body runs, a flexible body data pre-processing algorithm has been coded to generate simulation input data from FEM output data. The code has been installed and verified on the Alliant computer in the MTB facility.

In the future, the flexible body math model will be incorporated into the host simulation. The cycle times associated with various configurations (i.e., various EOM complexity levels) will need to be determined. Modifications to the flexible body math model may be required in order to reduce the simulation cycle time.

APPENDIX A

MATH MODEL SIMULATION CODE

ASCII Input File: MTBSIM.INP

```
0.002, 500, 5      Integ step size (sec); dt per CRT; dt per write
2.                !Stop time
2                  !INT: 1=Euler (Adams1), 2=Adams3-1, 3=Adams23-16+5
```

ASCII Input File: BODY1.INP

```
6547.              ! mass (c-u); Free-Free Brass Beam Model
32.74, 0., 0.      ! inertia matrix (c-u); symmetric
0., 335000., 0.
0., 0., 335000.
0., 0., 0.          ! P_B_I_I (length c-u); SAME AS GSV(1-3)
0., 0., 0.          ! PD_B_I_B (rate c-u); SAME AS QSV(1-3)
0., 0., 0.          ! W_B_I_B (deg/sec); SAME AS QSV(4-6)
1, 2, 3             ! I <-- BODY: Euler rot seq (1=x, 2=y, 3=z)
0., 0., 0.          ! I <-- BODY: Euler angles (deg)
4, 4                ! # nodes, NODE_DOCK; BODY origin at CM
-12.39, 0., 0.      ! node 1 in BODY
-4.13, 0., 0.       ! node 2 in BODY
4.13, 0., 0.        ! node 3 in BODY
2.39, 0., 0.        ! node 4
0., 0., 0.          ! P_D1_S1_S1
1., 0., 0.          ! T_B1_D1O
0., 1., 0.
0., 0., 1.
1., 0., 0.          ! T_D1_S1
0., 1., 0.
0., 0., 1.
.T.,.F.,.T.        ! EOM: 1)T=EHOMIT; 2)T=(1)+ERNLT 3)T=FREE-FREE
3                  ! # modes
6.2993, 6.2993, 17.4063 ! modal freqs (rad/sec)
0.0, 0.0, 0.0       ! modal damping (zeta_i)
0., 0., 0.          modal coords: eta(NM); GSV(>7)
0., 0., 0.          ! deriv modal coords: etadot(NM); QSV(>6)
1.0000000128348988 0.000000000E+000 -6.931270607381940E-018 !Gen mass
0.000000000E+000 1.0000000128348988 0.000000000000000E+000
-5.143184055728112E-017 0.0000000E+000 1.0000000170771713
.0000000E+00 .0000000E+00 .2485285E-01 ! Translational Modal Gains: Phi
```

.0000000E+00	.2485285E-01	.0000000E+00	
.0000000E+00	.0000000E+00	.2462543E-01	
.0000000E+00	.0000000E+00	-.9216621E-02	
.0000000E+00	-.9216621E-02	.0000000E+00	
.0000000E+00	.0000000E+00	-.1633864E-01	
.0000000E+00	.0000000E+00	-.9216621E-02	
.0000000E+00	-.9216621E-02	.0000000E+00	
.0000000E+00	.0000000E+00	.1633864E-01	
.0000000E+00	.0000000E+00	.2485285E-01	
.0000000E+00	.2485285E-01	.0000000E+00	
.0000000E+00	.0000000E+00	-.2462543E-01	
.0000000E+00	.4663153E-02	.0000000E+00	! Rotational Modal Gains: Psi
.0000000E+00	.0000000E+00	-.4663153E-02	
.0000000E+00	.7881236E-02	.0000000E+00	
.0000000E+00	.2742354E-02	.0000000E+00	
.0000000E+00	.0000000E+00	-.2742354E-02	
.0000000E+00	-.1159555E-02	.0000000E+00	
.0000000E+00	-.2742354E-02	.0000000E+00	
.0000000E+00	.0000000E+00	.2742354E-02	
.0000000E+00	-.1159555E-02	.0000000E+00	
.0000000E+00	-.4663153E-02	.0000000E+00	
.0000000E+00	.0000000E+00	.4663153E-02	
.0000000E+00	.7881236E-02	.0000000E+00	
.0000000E+00	.0000000E+00	-.5443560E-05	! Gamma_0
.0000000E+00	-.5443560E-05	.0000000E+00	
.0000000E+00	.0000000E+00	.3552714E-14	
.0000000E+00	-.3552714E-13	.0000000E+00	! Gamma_1
.0000000E+00	.0000000E+00	.3552714E-13	
.0000000E+00	-.2119259E-04	.0000000E+00	
.0000000E+00	.0000000E+00	.0000000E+00	! Gamma_2jk
-.1000715E+01	.0000000E+00	.0000000E+00	
.0000000E+00	.0000000E+00	.0000000E+00	
.1000715E+01	.0000000E+00	.0000000E+00	
.0000000E+00	.0000000E+00	.0000000E+00	
.5551115E-16	.0000000E+00	.0000000E+00	
.0000000E+00	.0000000E+00	.0000000E+00	
-.5551115E-16	.0000000E+00	.0000000E+00	
.0000000E+00	.0000000E+00	.0000000E+00	
.0000000E+00	.0000000E+00	.0000000E+00	! I_1k
.0000000E+00	.0000000E+00	.0000000E+00	
.0000000E+00	.0000000E+00	.0000000E+00	
.0000000E+00	.0000000E+00	.0000000E+00	
.0000000E+00	.0000000E+00	.0000000E+00	
.0000000E+00	.0000000E+00	.0000000E+00	
.0000000E+00	.0000000E+00	.0000000E+00	
.0000000E+00	.0000000E+00	.0000000E+00	
.0000000E+00	.0000000E+00	.0000000E+00	
.0000000E+00	.0000000E+00	.0000000E+00	
.4752610E+00	.0000000E+00	.0000000E+00	
.1000715E+01	.0000000E+00	.0000000E+00	! I_2jk
.0000000E+00	.1000715E+01	.0000000E+00	

.0000000E+00	.0000000E+00	.0000000E+00
.0000000E+00	.0000000E+00	.0000000E+00
.0000000E+00	.0000000E+00	.0000000E+00
.0000000E+00	-.1000715E+01	.0000000E+00
.5551115E-16	.0000000E+00	.0000000E+00
.0000000E+00	.5551115E-16	.0000000E+00
.0000000E+00	.0000000E+00	.0000000E+00
.0000000E+00	.0000000E+00	.0000000E+00
.0000000E+00	.0000000E+00	-.1000715E+01
.0000000E+00	.0000000E+00	.0000000E+00
.1000715E+01	.0000000E+00	.0000000E+00
.0000000E+00	.0000000E+00	.0000000E+00
.0000000E+00	.0000000E+00	.1000715E+01
.0000000E+00	.0000000E+00	.0000000E+00
.0000000E+00	.0000000E+00	-.5551115E-16
.0000000E+00	.0000000E+00	.0000000E+00
-.2775558E-16	.0000000E+00	.0000000E+00
.0000000E+00	-.2775558E-16	.0000000E+00
.0000000E+00	.0000000E+00	.0000000E+00
.0000000E+00	.0000000E+00	.0000000E+00
.0000000E+00	.0000000E+00	.0000000E+00
.0000000E+00	.0000000E+00	.0000000E+00
.0000000E+00	.2775558E-16	.0000000E+00
.1013732E+01	.0000000E+00	.0000000E+00
.0000000E+00	.1013732E+01	.0000000E+00
.0000000E+00	.0000000E+00	.0000000E+00

ASCII Input File: BODY2.INP

Same structure as BODY1.INP omitting P_D1_S1_S1 and T_D1_S1 data since the F/M sensor is located only on the target body.

```

CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC
cc
cc Filename:      BODY.INC
cc
cc Function:      Include file for variable/common declarations
cc                relating to body data.
cc
cc Source:        JC   Date: 2/91
cc
cc Comments:      Double precision
cc
CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC

```

```

INTEGER MNM                !MAX # OF MODES PER BODY
INTEGER MNN                !MAX # OF NODES PER BODY
INTEGER NBODES             !# OF BODIES
PARAMETER (MNM = 10, MNN = 10, NBODES = 2)
INTEGER NODE_DOCK(NBODES) !BODY DOCKING PORT NODE NUM
INTEGER NMODES(NBODES)    !# FLEX MODES PER BODY
INTEGER NNODES(NBODES)    !# NODES PER BODY
INTEGER INDX(NBODES,6+MNM) !REC OF LUDCMP ROW PERMUTATIONS

```

```

cc Simulation options vars
LOGICAL*2 EOM(NBODES,3)    !OPTIONS FOR EOM COMPLEXITY
INTEGER INT                !INTEGRATION OPTIONS
REAL*8 TIME                !SIM TIME (SEC)
REAL*8 DT                  !INTEGRATION STEP SIZE (SEC)

```

```

cc Mass properties/modal data vars
REAL*8 MASS(NBODES)        !BODY MASS
REAL*8 ICM(NBODES,3,3)     !INERTIA MAT ABOUT BODY CM
REAL*8 GMASS(NBODES,MNM,MNM) !GEN MASS MATRIX PER BODY
REAL*8 NODES(NBODES,MNN,3) !NODE LOCATIONS PER BODY IN
                           !   BODY

REAL*8 MOD_FREQ(NBODES,MNM) !MODAL FREQUENCIES PER BODY
REAL*8 MOD_ZETA(NBODES,MNM) !MODAL DAMPING PER BODY
REAL*8 PHI(NBODES,MNN,MNM,3) !TRANS MODAL GAINS
REAL*8 PSI(NBODES,MNN,MNM,3) !ROT MODAL GAINS

REAL*8 PHIETA(NBODES,MNN,3) !SUM OF PHI*ETA PER NODE
REAL*8 PHIETAD(NBODES,MNN,3) !SUM OF PHI*ETAD PER NODE
REAL*8 PSIETA(NBODES,MNN,3) !SUM OF PSI*ETA PER NODE
REAL*8 PSIETAD(NBODES,MNN,3) !SUM OF PSI*ETAD PER NODE

```

```

cc Mass Integral Terms
REAL*8 GAM_0(NBODES,3,MNM) !ARR OF GAM_0 MASS INTEGS
REAL*8 GAM_1(NBODES,3,MNM) !ARR OF GAM_1 MASS INTEGS
REAL*8 GAM_2JK(NBODES,MNM,MNM,3) !ARR OF GAM_2JK MASS INTS

```

```

REAL*8 I_1K(NBODES,MNM,3,3)      !ARR OF 3X3 I_1K MASS INTS
REAL*8 I_2JK(NBODES,MNM,MNM,3,3) !ARR OF 3X3 I_2JK MASS INTS

cc
REAL*8 LU_MASSM(NBODES,6+MNM,6+MNM) !BODY #NB LUDCMP
!   MASS MAT

cc External body forces/moments
REAL*8 FN_BODY(NBODES,MNM,3)      !FORCES APPLIED AT EACH NODE
REAL*8 MN_BODY(NBODES,MNM,3)      !MOMENTS APPL AT EACH NODE

cc Quasi-coord (with angular body rates) data
REAL*8 QSV(NBODES,6+MNM) ! QUASI COORD VECTOR
REAL*8 QSD(NBODES,6+MNM) !QUASI VECTOR DERIV
REAL*8 QSDO_1(NBODES,6+MNM) !1 OLD QUASI VECTOR DERIV
REAL*8 QSDO_2(NBODES,6+MNM) !2 OLD QUASI VECTOR DERIV

cc Generalized coord (quaternion for orientation) data
REAL*8 GSV(NBODES,7+MNM) !GEN COORD VECTOR
REAL*8 GSD(NBODES,7+MNM) !GEN VECTOR DERIV
REAL*8 GSDO_1(NBODES,7+MNM) !1 OLD GEN VECTOR DERIV
REAL*8 GSDO_2(NBODES,7+MNM) ! 2 OLD GEN VECTOR DERIV

cccccccccccccccccccccccccccc BODY #1 - TARGET ccccccccccccccccccccccccccc
cc Position/rate vectors
REAL*8 P_D1O_B1_B1(3) !UNDEFORM D1 POS VECT WRT B1 IN B1 (L)
REAL*8 P_D1_S1_S1(3) !ASSUMED CONST POS D1 WRT S1 IN S1 (L)

cc Transformations/quaternions
REAL*8 T_B1_D1O(3,3) !UNDEFORMED D1 <-- B1 TRANS
REAL*8 T_D1_S1(3,3) !ASSUMED CONST TRANS D1 <--S1

cccccccccccccccccccccccccccc BODY #2 - TARGET ccccccccccccccccccccccccccc
cc Position/rate vectors
REAL*8 P_D2O_B2_B2(3) !UNDEFORM D2 POS VECT WRT B2 IN B2 (L)

cc Transformations/quaternions
REAL*8 T_B2_D2O(3,3) !UNDEFORMED D2 <-- B2 TRANS

cccccccccccccccccccccccccccc COMMONS ccccccccccccccccccccccccccccccc
cc Simulation options
COMMON / OPTIONS / EOM, INT

cc Simulation control related commons
COMMON / SIMCON / TIME, DT

cc Mass properties related commons
COMMON / RIGID / MASS, ICM, NNODES, NODES
COMMON / FLEX / GMASS
COMMON / MASSINT / GAM_0, GAM_1, GAM_2JK, I_1K, I_2JK

```

COMMON / LHS / LU_MASSM, INDX

cc Modal data related commons

COMMON / MODES / NMODES, NODE_DOCK, PHI, PSI,
& MOD_FREQ, MOD_ZETA, PHIETA, PHJETAD,
& PSIETA, PSIETAD

cc External body forces/moments per body node commons

COMMON / NODEFT / FN_BODY, MN_BODY

cc Coord vector related commons

COMMON / QSTATE / QSV, QSVD, QSVDO_1, QSVDO_2
COMMON / GSTATE / GSV, GSVD, GSVDO_1, GSVDO_2

cc Body #1 related commons

COMMON / VECT1 / P_D1O_B1_B1, P_D1_S1_S1
COMMON / TRANS1 / T_B1_D1O, T_D1_S1

cc Body #2 related commons

COMMON / VECT2 / P_D2O_B2_B2
COMMON / TRANS2 / T_B2_D2O


```

CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC
cc
cc Filename:      INTFAC.INC
cc
cc Function:      Include file for variable/common declarations
cc                  of variables necessary to interface to NASA's
cc                  code.
cc
cc Source:        JC   Date: 3/91
cc
cc Comments:      Double precision
cc
CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC

```

```

REAL*8 T_B1_I(3,3)      !UPDATED IN INTFAC
REAL*8 T_B2_I(3,3)      !UPDATED IN INTFAC

REAL*8 T_D1_D2(3,3)      !UPDATED IN INTFAC
REAL*8 P_D2_D1_D1(3)      !UPDATED IN INTFAC
REAL*8 PD_D2_D1_D1(3)      UPDATED IN INTFAC
REAL*8 W_D2_D1_D2(3)      !UPDATED IN INTFAC
REAL*8 T_B1_B2(3,3)      !UPDATED IN INTFAC
REAL*8 T_B2_B1(3,3)      UPDATED IN INTFAC
REAL*8 P_B2_B1_B1(3)      !UPDATED IN INTFAC
REAL*8 T_D1_B1(3,3)      !UPDATED IN INTFAC
REAL*8 T_D2_B2(3,3)      !UPDATED IN INTFAC
REAL*8 P_D1_B1_B1(3)      !UPDATED IN INTFAC
REAL*8 P_D2_B2_B2(3)      !UPDATED IN INTFAC
REAL*8 FS1_S1(3)          !FROM BODY #1 F/M SENSOR
REAL*8 MS1_S1(3)          !FROM BODY #1 F/M SENSOR

COMMON / BODY_I / T_B1_I, T_B2_I
COMMON / REL_DOCK / T_D1_D2, P_D2_D1_D1, PD_D2_D1_D1,
&      W_D2_D1_D2
COMMON / BODY_BODY / T_B1_B2, T_B2_B1, P_B2_B1_B1
COMMON / BODY_DOCK / T_D1_B1, T_D2_B2, P_D1_B1_B1,
&      P_D2_B2_B2
COMMON / FMSENSOR / FS1_S1, MS1_S1

```

```

CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC
PROGRAM TESTMTBF
CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC
cc
cc Function:      Test driver for checkout of the NAMTB flex
cc                body simulation; formatted MATLAB output.
cc
cc Source:        JC    Date: 3/91
cc
cc Comments:      Double precision
cc
CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC
IMPLICIT NONE

cc Include files
INCLUDE 'BODY.INC'
INCLUDE 'INTFAC.INC'

cc Locals
INTEGER I, NB
INTEGER NCRT, CRT, NWRT, WRIT
INTEGER SEQ(3)
REAL*8 TSTOP
REAL*8 TEMP3_1(3), TEMP3_2(3), TEMP3_3(3)
REAL*8 TEMP3_4(3), TEMP3_5(3), W_B_I_I(3), EA(3)
REAL*8 TEMP4_1(4), TEMP4_2(4)
REAL*8 T_I_B1(3,3), DCM(3,3), T_I_D1(3,3), T_B1_D1(3,3)
REAL*8 T_I_B2(3,3)

CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC
cc Initialization
DATA TIME, NCRT / 0.0, 0 /
DATA SEQ / 1, 2, 3 /

cc First pass sim initialization
CALL DYNAMIC

OPEN(UNIT=13,FILE='MTBSIM.INP',FORM='FORMATTED',STATUS='OLD')

OPEN(UNIT=41,FILE='MTB1.MAT',FORM='FORMATTED')
OPEN(UNIT=42,FILE='MTB2.MAT',FORM='FORMATTED')

WRITE(6,*)
READ(13,*) DT, CRT, WRIT
NWRT = WRIT
NCRT = CRT
READ(13,*) TSTOP
READ(13,*) INT

```



```

        NB=1
        WRITE(41,444) TIME, (QSV(NB,I),I=1,6+NMODES(NB)),
&      (GSV(NB,I),I=1,3), (EA(I),I=1,3), (GSV(NB,I),
&      I=8,7+NMODES(NB)), (PHIETA(NB,4,I),I=1,3),
&      (PHIETAD(NB,4,I),I=1,3), (PSIETA(NB,4,I),I=1,3),
&      (PSIETAD(NB,4,I),I=1,3), (TEMP3_1(I),I=1,3),
&      (TEMP3_2(I),I=1,3), (GSVD(NB,I),I=1,3),
&      (W_B_I_I(I),I=1,3)

cc Body 2
        CALL Q_T_DCM(TEMP4_1, DCM)
        CALL DCM_T_EA(DCM, SEQ, EA)

        CALL D_MTRANS(T_B2_I, T_I_B2, 3, 3)
        CALL D_MMUL(T_I_B2, TEMP3_4, W_B_I_I, 3, 3, 1)

        NB = 2
        WRITE(42,444) TIME, (QSV(NB,I),I=1,6+NMODES(NB)),
&      (GSV(NB,I),I=1,3), (EA(I),I=1,3), (GSV(NB,I),
&      I=8,7+NMODES(NB)), (PHIETA(NB,4,I),I=1,3),
&      (PHIETAD(NB,4,I),I=1,3), (PSIETA(NB,4,I),I=1,3),
&      (PSIETAD(NB,4,I),I=1,3),
&      (GSVD(NB,I),I=1,3),
&      (W_B_I_I(I),I=1,3)

        ENDIF
        NWRT = NWRT + 1

444  FORMAT(2X,50E15.7)
CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC
        CALL DYNAMIC

cc Adjust step size on final step
        IF( (TIME + DT) .GT. TSTOP .AND. (TIME + DT) .LT.
&      (TSTOP + DT/2) ) THEN
            DT = TSTOP - TIME
            NWRT = WRIT
            NCRT = CRT
        ENDIF

cc Increment sim time and check for end
        TIME = TIME + DT
        IF(TIME .LE. TSTOP) GOTO 10

        CLOSE(41)
        CLOSE(42)

cc Create output data label files
        open(unit=41, file='mtb1.lab')

```

```

write(41,500) 1,'Time      '
write(41,500) 2,'PD_B_I_B '
write(41,500) 3,'PD_B_I_B '
write(41,500) 4,'PD_B_I_B '
write(41,500) 5,'W_B_I_B  '
write(41,500) 6,'W_B_I_B  '
write(41,500) 7,'W_B_I_B  '
write(41,500) 8,'ETA_D(1) '
write(41,500) 9,'ETA_D(2) '
write(41,500) 10,'ETA_D(3) '
write(41,500) 11,'P_B_I_I '
write(41,500) 12,'P_B_I_I '
write(41,500) 13,'P_B_I_I '
write(41,500) 14,'B_I EA(1)'
write(41,500) 15,'B_I EA(2)'
write(41,500) 16,'B_I EA(3)'
write(41,500) 17,'ETA(1)  '
write(41,500) 18,'ETA(2)  '
write(41,500) 19,'ETA(3)  '
write(41,500) 20,'PHIETA4x '
write(41,500) 21,'PHIETA4y '
write(41,500) 22,'PHIETA4z '
write(41,500) 23,'PHIETAD4x'
write(41,500) 24,'PHIETAD4y'
write(41,500) 25,'PHIETAD4z'
write(41,500) 26,'PSIETA4x '
write(41,500) 27,'PSIETA4y '
write(41,500) 28,'PSIETA4z '
write(41,500) 29,'PSIETAD4x'
write(41,500) 30,'PSIETAD4y'
write(41,500) 31,'PSIETAD4z'
write(41,500) 32,'P_D2_D1_I'
write(41,500) 33,'P_D2_D1_I'
write(41,500) 34,'P_D2_D1_I'
write(41,500) 35,'PD_D2_D1i'
write(41,500) 36,'PD_D2_D1i'
write(41,500) 37,'PD_D2_D1i'
write(41,500) 38,'PD_B_I_I '
write(41,500) 39,'PD_B_I_I '
write(41,500) 40,'PD_B_I_I '
write(41,500) 41,'W_B_I_I '
write(41,500) 42,'W_B_I_I '
write(41,500) 43,'W_B_I_I '

```

```

open(unit=42, file='mtb2.lab')

```

```

write(42,500) 1,'Time      '
write(42,500) 2,'PD_B_I_B '

```

```

write(42,500) 3,'PD_B_I_B '
write(42,500) 4,'PD_B_I_B '
write(42,500) 5,'W_B_I_B '
write(42,500) 6,'W_B_I_B '
write(42,500) 7,'W_B_I_B '
write(42,500) 8,'ETA_D(1) '
write(42,500) 9,'ETA_D(2) '
write(42,500) 10,'ETA_D(3) '
write(42,500) 11,'P_B_I_I '
write(42,500) 12,'P_B_I_I '
write(42,500) 13,'P_B_I_I '
write(42,500) 14,'B_I EA(1)'
write(42,500) 15,'B_I EA(2)'
write(42,500) 16,'B_I EA(3)'
write(42,500) 17,'ETA(1) '
write(42,500) 18,'ETA(2) '
write(42,500) 19,'ETA(3) '
write(42,500) 20,'PHIETA4x '
write(42,500) 21,'PHIETA4y '
write(42,500) 22,'PHIETA4z '
write(42,500) 23,'PHIETAD4x'
write(42,500) 24,'PHIETAD4y'
write(42,500) 25,'PHIETAD4z'
write(42,500) 26,'PSIETA4x '
write(42,500) 27,'PSIETA4y '
write(42,500) 28,'PSIETA4z '
write(42,500) 29,'PSIETAD4x'
write(42,500) 30,'PSIETAD4y'
write(42,500) 31,'PSIETAD4z'
write(42,500) 32,'PD_B_I_I '
write(42,500) 33,'PD_B_I_I '
write(42,500) 34,'PD_B_I_I '
write(42,500) 35,'W_B_I_I '
write(42,500) 36,'W_B_I_I '
write(42,500) 37,'W_B_I_I '

```

```

500 format(4x,i3,2x,a9)

```

```

CLOSE(42)
CLOSE(42)

```

```

STOP
END

```

```

cccccccccccccccccccccccccccccccccccccccccccccccccccccccccccc
SUBROUTINE DYNAMIC
cccccccccccccccccccccccccccccccccccccccccccccccccccccccccccc
cc
cc Function:      Driver of flexible body contact dynamics model
cc                for the 6-DOF docking/berthing facility.
cc
cc Source:   JC      Date: 2-3/91
cc
cc Comments:   Double precision
cc             Data via common block declarations
cc
cccccccccccccccccccccccccccccccccccccccccccccccccccccccccccc
IMPLICIT NONE

cc Include files
INCLUDE 'BODY.INC'

cc Locals
INTEGER I
LOGICAL PASS          ! SUBROUTINE PASS FLAG
DATA PASS / .FALSE. /

cc Initialization
IF(.NOT. PASS) THEN
  PASS = .TRUE.
  CALL INIT
  CALL INTFAC
  CALL FMTRANS
  CALL CONTROL(1)
c   CALL CONTROL(2)
  CALL LUMASSM(1)
  CALL LUMASSM(2)
  CALL PLANT(1)
  CALL PLANT(2)
  RETURN
ENDIF

cccccccccccccccccccccccc Main Loop ccccccccccccccccccccccccccccc

cc Compute contact forces/torques based on sensor output
CALL FMTRANS

cc Compute control system effects
CALL CONTROL(1)
c   CALL CONTROL(2)

cc Compute time varying mass matrix for Target and Chase if nec
IF(.NOT. EOM(1,1) .OR. .NOT. EOM(1,3)) THEN

```

```

    CALL LUMASSM(1)
ENDIF
IF(.NOT. EOM(2,1) .OR. .NOT. EOM(2,3)) THEN
    CALL LUMASSM(2)
ENDIF

cc Compute right hand side of EOMs for bodies 1 and 2
CALL PLANT(1)
CALL PLANT(2)

cc Integrate EOMs and handle quasi-to-generalized coord.
cc transformations
CALL INTEG(1)
CALL INTEG(2)

cc Compute body relative data
CALL INTFAC

RETURN
END

```



```

CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC
SUBROUTINE INIT
CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC
cc
cc Function:      Initialization of mass properties, geometry,
cc                state vectors, etc.
cc
cc Source:  JC    Date: 2/91
cc
cc Comments:      Double precision
cc                Data input via reads from ASCII input files
cc                Data distributed via common block declarations
cc
CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC
IMPLICIT NONE

cc Include files
INCLUDE 'BODY.INC'
INCLUDE 'INTFAC.INC'

cc Local declarations
INTEGER I, J, K, L
INTEGER NB, ROT_SEQ(3)
REAL*8 PI, RTD
REAL*8 ROT_ANG(3)
REAL*8 TEMP1_33(3,3), TEMP2_33(3,3), TEMP3_33(3,3), DCM(3,3)
REAL*8 Q4(4)
PARAMETER ( PI = 3.141592654d0 )
PARAMETER ( RTD = 180.d0/PI )

cc Open pertinent data files
OPEN(UNIT=11,FILE='BODY1.INP',FORM='FORMATTED',STATUS='OLD')
OPEN(UNIT=12,FILE='BODY2.INP',FORM='FORMATTED',STATUS='OLD')
OPEN(UNIT=20,FILE='INIT.OUT',FORM='FORMATTED',STATUS='UNKNOWN')

cc Read in body #1 (target) data
NB = 1

WRITE(20,*)
WRITE(20,'(10X,"BODY #",I2," INITIALIZATION DATA")') NB
WRITE(20,*)

READ(11,*) MASS(NB)          !MASS
WRITE(20,'(2X,"MASS: ",F12.4)') MASS(NB)

READ(11,*) (ICM(NB,1,J), J=1,3)  !!INERTIA MATRIX CM
READ(11,*) (ICM(NB,2,J), J=1,3)
READ(11,*) (ICM(NB,3,J), J=1,3)
WRITE(20,'(2X,"ICM: ",3F12.4)') (ICM(NB,1,J),J=1,3)

```

```

WRITE(20,'(2X,"ICM: ",3F12.4)') (ICM(NB,2,J),J=1,3)
WRITE(20,'(2X,"ICM: ",3F12.4)') (ICM(NB,3,J),J=1,3)
WRITE(20,*)

```

```

READ(11,*) (GSV(NB,J),J=1,3) !BODY CM POS WRT I IN I
READ(11,*) (QSV(NB,J),J=1,3) !BODY TRANS RATE WRT I IN BODY
READ(11,*) (QSV(NB,J),J=4,6) !BODY ANG RATE WRT I IN BODY

```

```

READ(11,*) (ROT_SEQ(J), J=1,3) !BODY <-- I ROT SEQUENCE
READ(11,*) (ROT_ANG(J), J=1,3) !BODY <-- I ROT ANGLES (DEG)

```

```

cc Compute trans from Inertial (I) to Body 1 (target) from
cc Euler angle sequence (I <-- Body)
DO 3 J=1,3
  QSV(NB,3+J) = QSV(NB,3+J) / RTD
  ROT_ANG(J) = ROT_ANG(J) / RTD
3 CONTINUE

```

```

CALL D_ROT(ROT_SEQ(1), ROT_ANG(1), TEMP1_33)
CALL D_ROT(ROT_SEQ(2), ROT_ANG(2), TEMP2_33)
CALL D_ROT(ROT_SEQ(3), ROT_ANG(3), TEMP3_33)

```

```

CALL D_MMUL(TEMP1_33, TEMP2_33, DCM, 3, 3, 3)
CALL D_MMUL(DCM, TEMP3_33, TEMP1_33, 3, 3, 3)
CALL D_MTRANS(TEMP1_33, T_B1_I, 3, 3)

```

```

cc Compute BODY <-- I quaternion and put in generalized state vect
CALL DCM_T_Q(T_B1_I, Q4)
CALL D_QNORM(Q4)

```

```

WRITE(20,*) ' T_B1_I: ',(T_B1_I(1,I),I=1,3)
WRITE(20,*) ' T_B1_I: ',(T_B1_I(2,I),I=1,3)
WRITE(20,*) ' T_B1_I: ',(T_B1_I(3,I),I=1,3)
WRITE(20,*)

```

```

DO 5 J=1,4
  GSV(NB,3+J) = Q4(J)
5 CONTINUE
WRITE(20,'(3X,"ANG RATES CONVERTED TO (RAD/S) IN SVs")')
WRITE(20,'(2X,"QSV(1->6): ",6F12.5)') (QSV(NB,J),J=1,6)
WRITE(20,*)
WRITE(20,'(2X,"GSV(1->7): ",7F12.5)') (GSV(NB,J),J=1,7)
WRITE(20,*)

```

```

cc Body node data
READ(11,*) NNODES(NB), NODE_DOCK(NB)
WRITE(20,'(2X,"NNODES: ",I3,8X,"NODE_DOCK: ",I3)')
& NNODES(NB), NODE_DOCK(NB)

```

```

DO 10 I=1,NNODES(NB)
  READ(11,*) (NODES(NB,I,J),J=1,3) !BODY NB NODES
  WRITE(20,'(2X,"NODE ",I2," ": ",10F12.5)')
    & I, (NODES(NB,I,J),J=1,3)
10 CONTINUE
  WRITE(20,*)

cc Sensor/docking port location/orientation initialization
DO 15 I=1,3
  P_D1O_B1_B1(I) = NODES(NB,NODE_DOCK(NB),I)
15 CONTINUE

  READ(11,*) (P_D1_S1_S1(I),I=1,3)

  READ(11,*) (T_B1_D1O(1,I),I=1,3)
  READ(11,*) (T_B1_D1O(2,I),I=1,3)
  READ(11,*) (T_B1_D1O(3,I),I=1,3)

  READ(11,*) (T_D1_S1(1,I),I=1,3)
  READ(11,*) (T_D1_S1(2,I),I=1,3)
  READ(11,*) (T_D1_S1(3,I),I=1,3)

cc Equations of motion complexity options
  READ(11,*) (EOM(NB,I),I=1,3) !LOGICAL EOM OPTIONS
  WRITE(20,'(2X,"EOM FLAGS (EHOMIT; ERNLT; FF MODES): ",3L6)')
    & (EOM(NB,I),I=1,3)

cc EOM options: ERNLT implies EHOMIT
  IF( EOM(NB,2) ) EOM(NB,1) = .TRUE.

cc Flex data
  READ(11,*) NMODES(NB)
  WRITE(20,'(2X,"NMODES: ",I3)') NMODES(NB)

  IF(NMODES(NB) .EQ. 0) GOTO 200

cc Modal frequencies (rad/sec)
  READ(11,*) (MOD_FREQ(NB,I),I=1,NMODES(NB))
  WRITE(20,'(2X,"MOD_FREQ: ",10F12.4)') (MOD_FREQ(NB,I),
    & I=1,NMODES(NB))

cc Modal damping
  READ(11,*) (MOD_ZETA(NB,I),I=1,NMODES(NB))
  WRITE(20,'(2X,"MOD_ZETA: ",10F12.4)') (MOD_ZETA(NB,I),
    & I=1,NMODES(NB))

cc Init modal coords (eta) and modal coord derivs (etad)
  READ(11,*) (GSV(NB,7+I),I=1,NMODES(NB))
  READ(11,*) (QSV(NB,6+I),I=1,NMODES(NB))

```

```

WRITE(20,'(2X,"ETA: ",10F12.4)') (GSV(NB,7+I),I=1,NMODES(NB))
WRITE(20,'(2X,"ETAD: ",10F12.4)') (QSV(NB,6+I),I=1,NMODES(NB))
WRITE(20,*)

cc Generalized mass matrix
DO 25 I=1,NMODES(NB)
  READ(11,*) (GMASS(NB,I,J),J=1,NMODES(NB)) !GEN MASS MAT
25 CONTINUE

cc Translation modal gains
DO 35 J=1,NNODES(NB)
  DO 30 K=1,NMODES(NB)
    READ(11,*) (PHI(NB,J,K,I),I=1,3)
30 CONTINUE
35 CONTINUE

cc Rotation modal gains
DO 45 J=1,NNODES(NB)
  DO 40 K=1,NMODES(NB)
    READ(11,*) (PSI(NB,J,K,I),I=1,3)
40 CONTINUE
45 CONTINUE

cc Standard mass integral terms
DO 50 J=1,NMODES(NB)
  READ(11,*) (GAM_0(NB,I,J), I=1,3) !GAM_0 MASS INTEGS
50 CONTINUE

  DO 60 J=1,NMODES(NB)
    READ(11,*) (GAM_1(NB,I,J), I=1,3) !GAM_1 MASS INTEGS
60 CONTINUE

cc Zero standard MITs for free-free mode option
IF( EOM(NB,3) ) THEN
  DO 65 J=1,NMODES(NB)
    DO 66 I=1,3
      GAM_0(NB,I,J) = 0.d0
      GAM_1(NB,I,J) = 0.d0
66 CONTINUE
65 CONTINUE
ENDIF

cc Higher order mass integral terms
DO 70 I=1,NMODES(NB)
  DO 80 J=1,NMODES(NB)
    READ(11,*) (GAM_2JK(NB,I,J,K),K=1,3) !GAM_2JK MASS INTEGS
80 CONTINUE
70 CONTINUE

```

```

DO 90 J=1,NMODES(NB)
  DO 100 K=1,3
    READ(11,*) (I_1K(NB,J,K,L),L=1,3)  !!_1K MASS INTEGERS
100  CONTINUE
90  CONTINUE

DO 110 I=1,NMODES(NB)
  DO 120 J=1,NMODES(NB)
    DO 130 K=1,3
      READ(11,*) (I_2JK(NB,I,J,K,L),L=1,3)  !!_1K MASS INTEGERS
130  CONTINUE
120  CONTINUE
110  CONTINUE

200  CONTINUE

cc Read in body #2 (CHASE) data
NB = 2

WRITE(20,*)
WRITE(20,*)
WRITE(20,*(10X,"BODY #",I2," INITIALIZATION DATA")) NB
WRITE(20,*)

READ(12,*) MASS(NB)          !MASS
WRITE(20,*(2X,"MASS: ",F12.4)) MASS(NB)

READ(12,*) (ICM(NB,1,J), J=1,3)  !!INERTIA MATRIX CM
READ(12,*) (ICM(NB,2,J), J=1,3)
READ(12,*) (ICM(NB,3,J), J=1,3)
WRITE(20,*(2X,"ICM: ",3F12.4)) (ICM(NB,1,J),J=1,3)
WRITE(20,*(2X,"ICM: ",3F12.4)) (ICM(NB,2,J),J=1,3)
WRITE(20,*(2X,"ICM: ",3F12.4)) (ICM(NB,3,J),J=1,3)
WRITE(20,*)

READ(12,*) (GSV(NB,J),J=1,3)  !BODY CM POS WRT I IN I
READ(12,*) (QSV(NB,J),J=1,3)  !BODY TRANS RATE WRT I IN BODY
READ(12,*) (QSV(NB,J),J=4,6)  !BODY ANG RATE WRT I IN BODY
READ(12,*) (ROT_SEQ(J), J=1,3) !BODY <-- I ROT SEQUENCE
READ(12,*) (ROT_ANG(J), J=1,3) !BODY <-- I ROT ANGLES (DEG)

cc Compute trans from Inertial (I) to Body 2 (CHASE) from
cc Euler angle sequence (I <-- Body)
DO 203 J=1,3
  QSV(NB,3+J) = QSV(NB,3+J) / RTD
  ROT_ANG(J) = ROT_ANG(J) / RTD
203 CONTINUE
CALL D_ROT(ROT_SEQ(1), ROT_ANG(1), TEMP1_33)
CALL D_ROT(ROT_SEQ(2), ROT_ANG(2), TEMP2_33)

```

```

CALL D_ROT(ROT_SEQ(3), ROT_ANG(3), TEMP3_33)

CALL D_MMUL(TEMP1_33, TEMP2_33, DCM, 3, 3, 3)
CALL D_MMUL(DCM, TEMP3_33, TEMP1_33, 3, 3, 3)
CALL D_MTRANS(TEMP1_33, T_B2_I, 3, 3)

cc Compute BODY <-- I quaternion and put in generalized state vect
CALL DCM_T_Q(T_B2_I, Q4)
CALL D_QNORM(Q4)

WRITE(20,*) ' T_B2_I: ',(T_B2_I(1,I),I=1,3)
WRITE(20,*) ' T_B2_I: ',(T_B2_I(2,I),I=1,3)
WRITE(20,*) ' T_B2_I: ',(T_B2_I(3,I),I=1,3)
WRITE(20,*)

DO 205 J=1,4
  GSV(NB,3+J) = Q4(J)
205 CONTINUE
WRITE(20, '(3X,"ANG RATES CONVERTED TO (RAD/S) IN SVs")')
WRITE(20, '(2X,"QSV(1->6): ",6F12.5)') (QSV(NB,J),J=1,6)
WRITE(20,*)
WRITE(20, '(2X,"GSV(1->7): ",7F12.5)') (GSV(NB,J),J=1,7)
WRITE(20,*)

cc Body node data
READ(12,*) NNODES(NB), NODE_DOCK(NB)
WRITE(20, '(2X,"NNODES: ",I3,8X,"NODE_DOCK: ",I3)')
&      NNODES(NB), NODE_DOCK(NB)

DO 210 I=1,NNODES(NB)
  READ(12,*) (NODES(NB,I,J),J=1,3) !BODY NB NODES
  WRITE(20, '(2X,"NODE ",I2," ": ",10F12.5)')
&      I, (NODES(NB,I,J),J=1,3)
210 CONTINUE
WRITE(20,*)

cc Sensor/docking port location/orientation initialization
DO 215 I=1,3
  P_D2O_B2_B2(I) = NODES(NB,NODE_DOCK(NB),I)
215 CONTINUE

READ(12,*) (T_B2_D2O(1,I),I=1,3)
READ(12,*) (T_B2_D2O(2,I),I=1,3)
READ(12,*) (T_B2_D2O(3,I),I=1,3)

cc Equations of motion complexity options
READ(12,*) (EOM(NB,I),I=1,3) !LOGICAL EOM OPTIONS
WRITE(20, '(2X,"EOM FLAGS (EHOMIT; ERNLT; FF MODES): ",3L6)')
&      (EOM(NB,I),I=1,3)

```

```

cc EOM options: ERNLT implies EHOMIT
  IF( EOM(NB,2) ) EOM(NB,1) = .TRUE.

cc Flex data
  READ(12,*) NMODES(NB)
  WRITE(20,'(2X,"NMODES: ",I3)') NMODES(NB)

  IF(NMODES(NB) .EQ. 0) GOTO 400

cc Modal frequencies (rad/sec)
  READ(12,*) (MOD_FREQ(NB,I),I=1,NMODES(NB))
  WRITE(20,'(2X,"MOD_FREQ: ",10F12.4)') (MOD_FREQ(NB,I),
&      I=1,NMODES(NB))

cc Modal damping
  READ(12,*) (MOD_ZETA(NB,I),I=1,NMODES(NB))
  WRITE(20,'(2X,"MOD_ZETA: ",10F12.4)') (MOD_ZETA(NB,I),
&      I=1,NMODES(NB))

cc Init modal coords (eta) and modal coord derivs (etad)
  READ(12,*) (GSV(NB,7+I),I=1,NMODES(NB))
  READ(12,*) (QSV(NB,6+I),I=1,NMODES(NB))
  WRITE(20,'(2X,"ETA: ",10F12.4)') (GSV(NB,7+I),I=1,NMODES(NB))
  WRITE(20,'(2X,"ETAD: ",10F12.4)') (QSV(NB,6+I),I=1,NMODES(NB))
  WRITE(20,*)

cc Generalized mass matrix
  DO 245 I=1,NMODES(NB)
    READ(12,*) (GMASS(NB,I,J),J=1,NMODES(NB)) !GEN MASS MAT
245  CONTINUE

cc Translation modal gains
  DO 225 J=1,NNODES(NB)
    DO 230 K=1,NMODES(NB)
      READ(12,*) (PHI(NB,J,K,I),I=1,3)
230  CONTINUE
225  CONTINUE

cc Rotation modal gains
  DO 235 J=1,NNODES(NB)
    DO 240 K=1,NMODES(NB)
      READ(12,*) (PSI(NB,J,K,I),I=1,3)
240  CONTINUE
235  CONTINUE

cc Standard mass integral terms
  DO 250 J=1,NMODES(NB)
    READ(12,*) (GAM_0(NB,I,J), I=1,3) !GAM_0 MASS INTEG

```

```

250  CONTINUE

      DO 260 J=1,NMODES(NB)
        READ(12,*) (GAM_1(NB,I,J), I=1,3) !GAM_1 MASS INTEG
260  CONTINUE

cc Zero standard MITs for free-free mode option
      IF( EOM(NB,3) ) THEN
        DO 265 J=1,NMODES(NB)
          DO 266 I=1,3
            GAM_0(NB,I,J) = 0.d0
            GAM_1(NB,I,J) = 0.d0
266    CONTINUE
265    CONTINUE
        ENDIF

cc Higher order mass integral terms
      DO 270 I=1,NMODES(NB)
        DO 280 J=1,NMODES(NB)
          READ(12,*) (GAM_2JK(NB,I,J,K),K=1,3) !GAM_2JK MASS INTEG
280    CONTINUE
270    CONTINUE

        DO 290 J=1,NMODES(NB)
          DO 300 K=1,3
            READ(12,*) (I_1K(NB,J,K,L),L=1,3)  !!_1K MASS INTEG
300    CONTINUE
290    CONTINUE

        DO 310 I=1,NMODES(NB)
          DO 320 J=1,NMODES(NB)
            DO 330 K=1,3
              READ(12,*) (I_2JK(NB,I,J,K,L),L=1,3) !!_1K MASS INTEG
330    CONTINUE
320    CONTINUE
310    CONTINUE

400  CONTINUE

      CLOSE(11)
      CLOSE(12)
      CLOSE(20)

      RETURN
      END

```



```

cc Compute contact moments (at D1 and D2) in respect body coords
  CALL D_VCROSS(P_D1_S1_S1, FS1_S1, TEMP1_3)
  CALL D_VCROSS(P_D2_S1_S1, FS1_S1, TEMP2_3)

  DO 40 I=1,3
    MD1_S1(I) = MS1_S1(I) - TEMP1_3(I)
    MD2_S1(I) = -MS1_S1(I) + TEMP2_3(I)
40  CONTINUE

  CALL D_MXV(T_B1_S1, MD1_S1, MD1_B1, 3, 3)
  CALL D_MXV(T_B2_S1, MD2_S1, MD2_B2, 3, 3)

cc Add contact moments to node moment vector
  DO 50 I=1,3
    MN_BODY(1,NODE_DOCK(1),I) = MD1_B1(I)
    MN_BODY(2,NODE_DOCK(2),I) = MD2_B2(I)
50  CONTINUE

  RETURN
  END

```

```

CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC
SUBROUTINE CONTROL(NB)
CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC
cc
cc Function:      Computes control forces and torques applied to
cc                bodies in the respective body fixed frames.
cc
cc Source:        JC   Date: 3/91
cc
cc Comments:      Double precision
cc                Data via common block declarations
cc                File serves as a place holder when no control
cc                is implemented
cc
CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC

cc Include files
    INCLUDE 'BODY.INC'
    INCLUDE 'INTFAC.INC'

cc Locals

RETURN
END

```

```

CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC
SUBROUTINE LUMASSM(NB)
CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC
cc
cc Function:      Computes mass matrix of body #NB in lower-upper
cc                decomposed format (i.e., LU decomposition);
cc                the resulting system can be solved using
cc                forwards/backwards substitution.
cc
cc Source:        JC    Date: 2/91
cc
cc Comments:      Double precision
cc                Data via common block declarations
cc
CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC
IMPLICIT NONE

cc Include files
INCLUDE 'BODY.INC'

cc Local declarations
INTEGER I, J, K, L
INTEGER NB, NM, INDXX(6+MNM)
REAL*8 MASSM(6+MNM,6+MNM)
REAL*8 TEMP_3MNM(3,MNM), TEMP_MNM(MNM)
REAL*8 GAM0_ETA(3), G0_ETA_T(3,3), GAM_2(3,MNM)
REAL*8 II_1(3,3), II_2(3,3), D

cc Start building lower triangular mass matrix (symmetric)
cc Add diagonal mass submatrix (3x3) assoc with rigid body trans
MASSM(1,1) = MASS(NB)
MASSM(2,2) = MASS(NB)
MASSM(3,3) = MASS(NB)

cc Add rigid body moment of inertia about body CM submatrix (3x3)
cc to lower triangular mass matrix
DO 10 I=1,3
  DO 20 J=1,I
    MASSM(3+I,3+J) = ICM(NB,I,J)
20  CONTINUE
10  CONTINUE

cc Check for no flex body modes
NM = NMODES(NB)
IF(NM .EQ. 0) GOTO 500

cc Add generalized mass submatrix (NMxNM) assoc with modal
cc coords to lower triangular mass matrix
DO 30 I=1,NM

```

```

        DO 40 J=1,I
            MASSM(6+I,6+J) = GMASS(NB,I,J)
40    CONTINUE
30    CONTINUE

cc Coupling between rigid body translation and rotation
cc via mass integral terms
cc Pull gamma0 (3xNM) matrix and eta vector (NMx1) out
    IF( .NOT. EOM(NB,3) ) THEN
        CALL D_ZERO(TEMP_3MNM,3,MNM)

        DO 50 J=1,NM
            DO 60 I=1,3
                TEMP_3MNM(I,J) = GAM_0(NB,I,J)
60        CONTINUE
            TEMP_MNM(J) = GSV(NB,7+J)
50    CONTINUE

cc Create gamma_0 (3 x MNM) * eta (MNM x 1) tilde matrix
        CALL D_MMUL(TEMP_3MNM,TEMP_MNM,GAM0_ETA,3,MNM,1)
        CALL D_TILDE(GAM0_ETA, G0_ETA_T)
    ELSE
        CALL D_ZERO(G0_ETA_T, 3, 3)
    ENDIF

cc Add gam0_eta tilde (3x3), gam0 (3xNM), and gam1 (3xNM)
cc submatrices to lower triangular mass matrix
    DO 70 I=1,3
        DO 80 J=1,3
            MASSM(3+I,J) = G0_ETA_T(I,J)
80    CONTINUE
        DO 90 J=1,NM
            MASSM(6+J,I) = GAM_0(NB,I,J)
            MASSM(6+J,3+I) = GAM_1(NB,I,J)
90    CONTINUE
70    CONTINUE

cc Check for elimination of higher order mass integral terms
    IF(EOM(NB,1)) GOTO 500

cc Compute gamma2 (3 x NM) and add to lower triangular mass mat
    DO 100 I=1,3
        DO 110 J=1,NM
            GAM_2(I,J) = 0.0D0
            DO 120 K=1,NM
                GAM_2(I,J) = GAM_2(I,J) + GSV(NB,7+K)*
&                GAM_2JK(NB,K,J,I)
120    CONTINUE
            MASSM(6+J,3+I) = MASSM(6+J,3+I) + GAM_2(I,J)

```

```

110  CONTINUE
100  CONTINUE

```

```

cc Compute II_1 (3 x 3) and II_2 (3 x 3); HOM-ITs affecting
cc moment of inertia matrix
DO 170 I=1,3
  DO 180 J=1,3
    II_1(I,J) = 0.0D0
    II_2(I,J) = 0.0D0
    DO 190 K=1,NM
      II_1(I,J) = II_1(I,J) + GSV(NB,7+K)*
&      ( I_1K(NB,K,I,J) + I_1K(NB,K,J,I) )
      DO 200 L=1,NM
        II_2(I,J) = II_2(I,J) + 0.5*GSV(NB,7+K)*GSV(NB,7+L)*
&      ( I_2JK(NB,K,L,I,J) + I_2JK(NB,L,K,I,J) )
200    CONTINUE
190    CONTINUE
180    CONTINUE
170  CONTINUE

```

```

cc Add II_1 and II_2 to lower triangular mass matrix
DO 220 I=1,3
  DO 230 J=1,I
    MASSM(3+I,3+J) = MASSM(3+I,3+J) + II_1(I,J) + II_2(I,J)
230  CONTINUE
220  CONTINUE

```

```

cc Now have FUN-WHOM-IT mass matrix

```

```

500  CONTINUE

```

```

cc Fill in upper triangle of mass matrix using symmetry prop
DO 400 I=1,6+NM
  DO 420 J=I,6+NM
    MASSM(I,J) = MASSM(J,I)
420  CONTINUE
400  CONTINUE

```

```

CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC
c  WRITE(6,*)
c  WRITE(6,*) ' MASSM:'
c  DO 425 I=1,6+NM
c    WRITE(6,'(1X,10E12.5)') (MASSM(I,J),J=1,6+NM)
c425 CONTINUE
c  WRITE(6,*)
CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC

```

```

cc Compute lower-upper decomposition of mass matrix
CALL LUDCMP(MASSM,6+NM,6+NM,INDXX,D)

```

```
cc Record in memory
  DO 510 I=1,6+NM
    INDX(NB,I) = INDXX(I)
    DO 520 J=1,6+NM
      LU_MASSM(NB,I,J) = MASSM(I,J)
520  CONTINUE
510  CONTINUE

  RETURN
  END
```



```

CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC
SUBROUTINE PLANT(NB)
CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC
cc
cc Function:           Computes right hand side of equations of motion.
cc                     These terms represent the effects of damping,
cc                     stiffness, and forcing functions.
cc
cc Source:             JC    Date: 2/91
cc
cc Comments:          Double precision
cc                     Data via common block declarations
cc
CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC
IMPLICIT NONE

cc Include files
INCLUDE 'BODY.INC'

cc Local declarations
INTEGER I, J, K, L, M, N
INTEGER NB, NM, INDXX(6+MNM)
REAL*8 MASSM(6+MNM,6+MNM)
REAL*8 TEMP1_MNM3(MNM,3), TEMP2_MNM3(MNM,3)
REAL*8 TEMP1_3MNM(3,MNM), TEMP2_3MNM(3,MNM)
REAL*8 TEMP1_3(3), TEMP2_3(3), TEMP3_3(3), TEMP4_3(3)
REAL*8 TEMP5_3(3), TEMP6_3(3), TEMP_33(3,3)
REAL*8 TEMP1_MNM(MNM), TEMP2_MNM(MNM)
REAL*8 GAM0_ETA(3), GAM0_ETAD(3)
REAL*8 GAM_2D(3,MNM), GAM_2(3,MNM)
REAL*8 RHS(6+MNM), FNM(MNM), TNM(MNM)
REAL*8 F_BODY(3), T_BODY(3)
REAL*8 WB(3)
REAL*8 II_1(3,3), II_2(3,3), II_1D(3,3), II_2D(3,3)
REAL*8 II(3,3), IID(3,3)
REAL*8 D_DOT

cc Rigid body translational/rotational forcing functions
cc Based on nodes locations defined over each body
cc Mult forces, pure torques, and POS x FORCE torques due to
cc node forces are summed into a single force/torque at CM.
CALL D_ZERO(F_BODY, 3, 1)
CALL D_ZERO(T_BODY, 3, 1)
CALL D_ZERO(RHS, 6+MNM, 1)
NM = NMODES(NB)

cc Compute torque at CM due to POS x FORCE of node forces
cc Include flexibility effects on node position vectors
DO 10 N=1,NNODES(NB)

```

```

DO 15 J=1,3
  TEMP1_3(J) = NODES(NB,N,J) + PHIETA(NB,N,J)
  TEMP2_3(J) = FN_BODY(NB,N,J)
15  CONTINUE
  CALL D_VCROSS(TEMP1_3, TEMP2_3, TEMP3_3)
  DO 20 J=1,3
    F_BODY(J) = F_BODY(J) + FN_BODY(NB,N,J)
    T_BODY(J) = T_BODY(J) + MN_BODY(NB,N,J) + TEMP3_3(J)
20  CONTINUE
10  CONTINUE

cc Build right hand side terms; pull out body ang vel vector
DO 30 J=1,3
  RHS(J) = F_BODY(J)
  RHS(3+J) = T_BODY(J)
  WB(J) = QSV(NB,3+J)
30  CONTINUE

cc Check for no flex body modes
IF(NM .EQ. 0) GOTO 500

cc Compute EHOMIT + ERNLT flex RHS terms
cc Generalized flex eqns forcing functions
DO 40 N=1,NNODES(NB)
  DO 50 M=1,NM
    DO 60 I=1,3
      TEMP1_MNM3(M,I) = PHI(NB,N,M,I)
      TEMP2_MNM3(M,I) = PSI(NB,N,M,I)
60  CONTINUE
50  CONTINUE
  DO 70 I=1,3
    TEMP1_3(I) = FN_BODY(NB,N,I)
    TEMP2_3(I) = MN_BODY(NB,N,I)
70  CONTINUE
  CALL D_MMUL(TEMP1_MNM3, TEMP1_3, FNM, MNM, 3, 1)
  CALL D_MMUL(TEMP2_MNM3, TEMP2_3, TNM, MNM, 3, 1)
  DO 80 M=1,NM
    RHS(6+M) = RHS(6+M) + FNM(M) + TNM(M)
80  CONTINUE
40  CONTINUE

cc Strain energy and damping terms
DO 90 M=1,NM
  RHS(6+M) = RHS(6+M) - GSV(NB,7+M) * MOD_FREQ(NB,M)**2 -
& 2.D0 * MOD_ZETA(NB,M) * MOD_FREQ(NB,M) * GSVD(NB,7+M)
90  CONTINUE

cc Check for EHOMIT + ERNLT option; ERNLT encompasses EHOMIT
IF(EOM(NB,2)) GOTO 500

```

```

cc Nonlinear translational RHS terms
cc Pull out gamma0 (3xNM) matrix, gamm1 (3xNM),
cc eta vector (NMx1), and etad vector (NMx1)
CALL D_ZERO(TEMP1_3MNM,3,MNM)
CALL D_ZERO(TEMP2_3MNM,3,MNM)

DO 100 J=1,NM
DO 110 I=1,3
TEMP1_3MNM(I,J) = GAM_0(NB,I,J)
TEMP2_3MNM(I,J) = GAM_1(NB,I,J)
110 CONTINUE
TEMP1_MNM(J) = GSV(NB,7+J)
TEMP2_MNM(J) = QSV(NB,6+J)
100 CONTINUE

cc Create gamma_0 (3 x MNM) * eta (MNM x 1) vector
cc Compute wb x (wb x gam0*eta)
IF( .NOT. EOM(NB,3) ) THEN
CALL D_MMUL(TEMP1_3MNM, TEMP1_MNM, GAM0_ETA, 3, MNM, 1)
CALL D_VCROSS(WB, GAM0_ETA, TEMP1_3)
CALL D_VCROSS(WB, TEMP1_3, TEMP2_3)
ELSE
TEMP2_3(1) = 0.D0
TEMP2_3(2) = 0.D0
TEMP2_3(3) = 0.D0
ENDIF

cc Compute wb x (mass*Rdcmb + 2*gam0*etad)
IF( .NOT. EOM(NB,3) ) THEN
CALL D_MMUL(TEMP1_3MNM, TEMP2_MNM, GAM0_ETAD, 3, MNM, 1)
ELSE
GAM0_ETAD(1) = 0.D0
GAM0_ETAD(2) = 0.D0
GAM0_ETAD(3) = 0.D0
ENDIF

DO 120 I=1,3
TEMP1_3(I) = MASS(NB) * QSV(NB,I) +
& 2.D0 * GAM0_ETAD(I)
120 CONTINUE

CALL D_VCROSS(WB, TEMP1_3, TEMP3_3)

cc Add translational terms to RHS
DO 130 I=1,3
RHS(I) = RHS(I) - TEMP2_3(I) - TEMP3_3(I)
TEMP1_3(I) = QSV(NB,I)
130 CONTINUE

```

```

cc Nonlinear rotational RHS terms
cc wb x (gam1 * etad)
  IF( .NOT. EOM(NB,3) ) THEN
    CALL D_MMUL(TEMP2_3MNM, TEMP2_MNM, TEMP2_3, 3, MNM, 1)
    CALL D_VCROSS(WB, TEMP2_3, TEMP3_3)

cc gam0*eta x ( wb x Rdcmb)
    CALL D_VCROSS(WB, TEMP1_3, TEMP4_3)
    CALL D_VCROSS(GAM0_ETA, TEMP4_3, TEMP5_3)
  ELSE
    DO 135 I=1,3
      TEMP3_3(I) = 0.D0
      TEMP5_3(I) = 0.D0
135  CONTINUE
  ENDIF

cc wb x (ICM*wb)
  IF(EOM(NB,1)) THEN
    DO 140 I=1,3
      DO 150 J=1,3
        TEMP_33(I,J) = ICM(NB,I,J)
150  CONTINUE
140  CONTINUE

    CALL D_MMUL(TEMP_33, WB, TEMP6_3, 3, 3, 1)
    CALL D_VCROSS(WB, TEMP6_3, TEMP4_3)
  ELSE
    TEMP4_3(1) = 0.D0
    TEMP4_3(2) = 0.D0
    TEMP4_3(3) = 0.D0
  ENDIF

cc Add nonlinear rotational terms to RHS
  DO 160 I=1,3
    RHS(3+I) = RHS(3+I) - TEMP3_3(I) - TEMP4_3(I) -
    &      TEMP5_3(I)
160  CONTINUE

cc Nonlinear RHS terms in flex eqns
cc Rdcmb . (wb x gam0) ; performed on sub vector basis in arrays
  IF( .NOT. EOM(NB,3) ) THEN
    DO 170 M=1,NM
      DO 180 I=1,3
        TEMP2_3(I) = GAM_0(NB,I,M)
180  CONTINUE
      CALL D_VCROSS(WB, TEMP2_3, TEMP3_3)

cc Add to RHS

```

```

      RHS(6+M) = RHS(6+M) + D_DOT(TEMP1_3, TEMP3_3, 3)
170  CONTINUE
      ENDIF

```

```

cc Check for elimination of higher order mass integral terms
   IF(EOM(NB,1)) GOTO 500

```

```

cc HOMITs on RHS of rotational eqns
cc Compute II (3 x 3); moment of inertia matrix including
cc effects of HOMITs
cc Compute IID (derivative of II) (3 x 3)
   DO 200 I=1,3
     DO 210 J=1,3
       II_1(I,J) = 0.D0
       II_2(I,J) = 0.D0
       II_1D(I,J) = 0.D0
       II_2D(I,J) = 0.D0
       DO 220 K=1,NM
         II_1(I,J) = II_1(I,J) + GSV(NB,7+K)*
&           ( I_1K(NB,K,I,J) + I_1K(NB,K,J,I) )
         II_1D(I,J) = II_1D(I,J) + QSV(NB,6+K)*
&           ( I_1K(NB,K,I,J) + I_1K(NB,K,J,I) )
         DO 230 L=1,NM
           II_2(I,J) = II_2(I,J) + 0.5*GSV(NB,7+K)*GSV(NB,7+L)*
&           ( I_2JK(NB,K,L,I,J) + I_2JK(NB,L,K,I,J) )
           II_2D(I,J) = II_2D(I,J) + 0.5*QSV(NB,6+K)*GSV(NB,7+L)*
&           ( I_2JK(NB,K,L,I,J) + I_2JK(NB,L,K,I,J) ) +
&           0.5*GSV(NB,7+K)*QSV(NB,6+L)*
&           ( I_2JK(NB,K,L,I,J) + I_2JK(NB,L,K,I,J) )
230      CONTINUE
220      CONTINUE
       II(I,J) = ICM(NB,I,J) + II_1(I,J) + II_2(I,J)
       IID(I,J) = II_1D(I,J) + II_2D(I,J)
210    CONTINUE
200    CONTINUE

```

```

cc Compute wb x (II * wb)
   CALL D_MMUL(II, WB, TEMP2_3, 3, 3, 1)
   CALL D_VCROSS(WB, TEMP2_3, TEMP1_3)

```

```

cc Compute IID * wb
   CALL D_MMUL(IID, WB, TEMP2_3, 3, 3, 1)

```

```

cc Compute wb x (gam2 * etad)
   DO 240 I=1,3
     DO 250 J=1,NM
       GAM_2(I,J) = 0.D0
       DO 260 K=1,NM
         GAM_2(I,J) = GAM_2(I,J) + GSV(NB,7+K)*

```

```

&          GAM_2JK(NB,K,J,I)
260  CONTINUE
250  CONTINUE
240  CONTINUE

      CALL D_MMUL(GAM_2, TEMP2_MNM, TEMP3_3, 3, MNM, 1)
      CALL D_VCROSS(WB, TEMP3_3, TEMP4_3)

cc  Add rotational HOMITs to RHS
      DO 270 I=1,3
        RHS(3+I) = RHS(3+I) - TEMP1_3(I) - TEMP2_3(I) -
&          TEMP4_3(I)

cc  Compute gam_2 dot (3 x NM) for flex eqns
      DO 280 J=1,NM
        GAM_2D(I,J) = 0.D0
        DO 290 K=1,NM
          GAM_2D(I,J) = GAM_2D(I,J) + QSV(NB,6+K)*
&          GAM_2JK(NB,K,J,I)
290  CONTINUE
280  CONTINUE
270  CONTINUE

cc  Compute -2*wb . gam_2dot on array sub vector basis
      DO 300 M=1,NM
        DO 310 I=1,3
          TEMP1_3(I) = GAM_2D(I,M)
310  CONTINUE

      RHS(6+M) = RHS(6+M) - 2.D0 * D_DOT(WB, TEMP1_3, 3)

cc  Pull out each I_1k 3x3 submatrix, compute wb'*I1m*wb,
cc  and add to RHS
      DO 320 I=1,3
        DO 330 J=1,3
          TEMP_33(I,J) = I_1K(NB, M, I, J)
330  CONTINUE
320  CONTINUE

      CALL D_MMUL(TEMP_33, WB, TEMP2_3, 3, 3, 1)
      RHS(6+M) = RHS(6+M) + D_DOT(WB, TEMP2_3, 3)

cc  Compute col vect over I: 0.5 * wb' * sumj(I_2jl + I_2lj) * wb
      CALL D_ZERO(TEMP_33, 3, 3)

      DO 340 K=1,NM
        DO 350 I=1,3
          DO 360 J=1,3
            TEMP_33(I,J) = TEMP_33(I,J) + GSV(NB,7+K) *

```

```

&      ( I_2JK(NB,K,M,I,J) + I_2JK(NB,M,K,I,J) )
360    CONTINUE
350    CONTINUE
340    CONTINUE

      CALL D_MMUL(TEMP_33, WB, TEMP3_3, 3, 3, 1)
      RHS(6+M) = RHS(6+M) + 0.5D0 * D_DOT(WB, TEMP3_3, 3)
300    CONTINUE

cc Solve system for deriv of quasi coord vector

500    CONTINUE

cc Extract appropriate upper triangular mass matrix
      DO 510 I=1,6+NM
        DO 520 J=1,6+NM
          MASSM(I,J) = LU_MASSM(NB,I,J)
520    CONTINUE
510    CONTINUE

CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC
c  WRITE(6,*)
c  WRITE(6,*) ' RHS: '
c  WRITE(6, '(1X,10E12.5)') (RHS(I),I=1,6+NM)
c  WRITE(6,*)
CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC

cc Solve upper triangular system (A*x = B) using back sub
      DO 530 J=1,6+NM
        INDXX(J) = INDX(NB,J)
530    CONTINUE

      CALL LUBKSB(MASSM, 6+NM, 6+NM, INDXX, RHS)

cc Store quasi-coord state vector derivative
      DO 550 I=1,6+NM
        QSVD(NB,I) = RHS(I)
550    CONTINUE

      RETURN
      END

```

```

CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC
SUBROUTINE INTEG(NB)
CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC
cc
cc Function:      Integrates derivatives of quasi-coord state
cc                vector (W); transforms (W) to generalized
cc                coord state vector derivative (qdot) and
cc                integrates qdot to produce q. Normalizes
cc                quaternion in generalized coord vector.
cc
cc                Four integration options based on integer INT:
cc                INT = 1 --> Euler or Adams 1
cc                INT = 2 --> Adams (3 - 1) / 2
cc                INT = 3 --> Adams (23 - 16 + 5) / 12
cc
cc Source:        JC    Date: 2/91
cc
cc Comments:      Double precision
cc                Data via common block declarations
cc
CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC
IMPLICIT NONE

cc Include files
cc INCLUDE 'BODY.INC'

cc Local declarations
cc INTEGER I
cc INTEGER NB, NM
cc REAL*8 TEMP1_3(3), TEMP2_3(3), TEMP_4(4), Q4(4)
cc REAL*8 WB(3), ET(4,3), TEMP1_33(3,3), TEMP2_33(3,3)

cc Initialization
cc NM = NMODES(NB)

cc Transform quasi-coord state vector (QSV) to deriv of gen coord
cc state vector (GSVD) in sub blocks; {GSVD} = [beta] * {QSV}

cc Rigid body translation rates
cc Extract body vel vector PD_B_I_B, ang vel vector W_B_I_B,
cc and BODY <-- I quaternion
cc DO 20 I=1,3
cc     TEMP1_3(I) = QSV(NB,I)
cc     WB(I) = QSV(NB,3+I)
cc     Q4(I) = GSV(NB,3+I)
20 CONTINUE
cc     Q4(4) = GSV(NB,7)

cc Transform PD_B_I_B to PD_B_I_I

```



```

CALL Q_T_DCM(Q4, TEMP1_33)
CALL D_MTRANS(TEMP1_33, TEMP2_33, 3, 3)
CALL D_MMUL(TEMP2_33, TEMP1_3, TEMP2_3, 3, 3, 1)

cc Build (Q_B_I)dot <-- angular body rate vector trans
ET(1,1) = Q4(4)
ET(1,2) = -Q4(3)
ET(1,3) = Q4(2)
ET(2,1) = Q4(3)
ET(2,2) = Q4(4)
ET(2,3) = -Q4(1)
ET(3,1) = -Q4(2)
ET(3,2) = Q4(1)
ET(3,3) = Q4(4)
ET(4,1) = -Q4(1)
ET(4,2) = -Q4(2)
ET(4,3) = -Q4(3)

cc Update Q_B_I derivative
CALL D_MMUL(ET, WB, TEMP_4, 4, 3, 1)

cc Add to GSVD
DO 30 I=1,3
  GSVD(NB,I) = TEMP2_3(I)
  GSVD(NB,3+I) = 0.5D0 * TEMP_4(I)
30  CONTINUE
  GSVD(NB,7) = 0.5D0 * TEMP_4(4)

cc Modal coord trans is an identity matrix
DO 50 I=1,NM
  GSVD(NB,7+I) = QSV(NB,6+I)
50  CONTINUE

cc Integrate gen coord state deriv vector based on INT option
cc Integrate quasi-coord state deriv vector based on INT option
IF(INT.EQ. 1) THEN
  DO 60 I=1,7+NM
    GSV(NB,I) = GSV(NB,I) + DT*GSVD(NB,I)
60  CONTINUE

  DO 70 I=1,6+NM
    QSV(NB,I) = QSV(NB,I) + DT*QSV(NB,I)
70  CONTINUE

  ELSEIF(INT.EQ. 2) THEN
    DO 80 I=1,7+NM
      GSV(NB,I) = GSV(NB,I) + (DT/2.D0) * ( 3.D0 * GSVD(NB,I) -
&      GSVD_1(NB,I) )
80  CONTINUE

```

```

DO 90 I=1,6+NM
  QSV(NB,I) = QSV(NB,I) + (DT/2.D0) * ( 3.D0 * QSVDO_1(NB,I) -
&      QSVDO_1(NB,I) )
90  CONTINUE

  ELSEIF(INT .EQ. 3) THEN
    DO 100 I=1,7+NM
      GSV(NB,I) = GSV(NB,I) + (DT/12.D0) * ( 23.D0 * GSVDO_1(NB,I) -
&      16.D0 * GSVDO_1(NB,I) + 5.D0 * GSVDO_2(NB,I) )
100  CONTINUE

    DO 110 I=1,6+NM
      QSV(NB,I) = QSV(NB,I) + (DT/12.D0) * ( 23.D0 * QSVDO_1(NB,I) -
&      16.D0 * QSVDO_1(NB,I) + 5.D0 * QSVDO_2(NB,I) )
110  CONTINUE

    ELSE
      WRITE(6,*) '*** INTEG: Integration Option Invalid ***'
      STOP
    ENDIF

cc Normalize Q_B_I and replace in generalized-coord vector
DO 120 I=1,4
  Q4(I) = GSV(NB,3+I)
120  CONTINUE

  CALL D_QNORM(Q4)

DO 130 I=1,4
  GSV(NB,3+I) = Q4(I)
130  CONTINUE

cc Save old state derivative data
DO 150 I=1,7+NM
  GSVDO_2(NB,I) = GSVDO_1(NB,I)
  GSVDO_1(NB,I) = GSVDO_2(NB,I)
150  CONTINUE

DO 160 I=1,6+NM
  QSVDO_2(NB,I) = QSVDO_1(NB,I)
  QSVDO_1(NB,I) = QSVDO_2(NB,I)
160  CONTINUE

RETURN
END

```

```

CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC
SUBROUTINE INTFAC
CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC
cc
cc Function:      Updates transformations and relative position
cc                and rate vectors; includes flexibility effects.
cc
cc Source:        JC    Date: 3/91
cc
cc Comments:      Double precision
cc                Data via common block declarations
cc
CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC
IMPLICIT NONE

cc Include files
INCLUDE 'BODY.INC'
INCLUDE 'INTFAC.INC'

cc Locals
INTEGER I, J, N, NB, M
REAL*8 T_I_B1(3,3)
REAL*8 T_D2_D1(3,3), T_D1O_B1(3,3)
REAL*8 T_D1_D1O(3,3), T_D2O_D2(3,3)
REAL*8 TEMP1_33(3,3), TEMP2_33(3,3)
REAL*8 TEMP1_3(3), TEMP2_3(3), TEMP3_3(3), TEMP4_3(3)
REAL*8 IDENT_33(3,3), Q4(4)
REAL*8 W_B1_I_B1(3), W_B2_I_B2(3)

DATA IDENT_33 / 1.D0, 0.D0, 0.D0, 0.D0, 1.D0, 0.D0,
&              0.D0, 0.D0, 1.D0/

cc Build transformations from undeformed docking port systems
cc to deformed docking point systems (due to infinitesimal
cc rotations resulting from flexibility).
DO 40 NB=1,NBODES

cc Compute BODY <-- I transformations from generalized coord
cc vector quaternions
DO 10 I= 1,4
    Q4(I) = GSV(NB,3+I)
10 CONTINUE

IF(NB.EQ. 1) THEN
    CALL Q_T_DCM(Q4, T_B1_I)
ELSEIF(NB.EQ. 2) THEN
    CALL Q_T_DCM(Q4, T_B2_I)
ENDIF

```

```

cc Compute flexibility effects at each node
  DO 20 N=1,NNODES(NB)
    DO 30 J=1,3

      PHIETA(NB,N,J) = 0.D0
      PHJETAD(NB,N,J) = 0.D0
      PSIETA(NB,N,J) = 0.D0
      PSJETAD(NB,N,J) = 0.D0

      DO 35 M=1,NMODES(NB)
        PHIETA(NB,N,J) = PHIETA(NB,N,J) + PHI(NB,N,M,J)*
&          GSV(NB,7+M)
        PHJETAD(NB,N,J) = PHJETAD(NB,N,J) + PHI(NB,N,M,J)*
&          QSV(NB,6+M)
        PSIETA(NB,N,J) = PSIETA(NB,N,J) + PSI(NB,N,M,J)*
&          GSV(NB,7+M)
        PSJETAD(NB,N,J) = PSJETAD(NB,N,J) + PSI(NB,N,M,J)*
&          QSV(NB,6+M)
35      CONTINUE
30      CONTINUE
20      CONTINUE

cc Compute and add in infinitesimal rotations at DOCKING node
cc due to flexibility
  TEMP1_3(1) = PSIETA(NB,NODE_DOCK(NB),1)
  TEMP1_3(2) = PSIETA(NB,NODE_DOCK(NB),2)
  TEMP1_3(3) = PSIETA(NB,NODE_DOCK(NB),3)

  CALL D_TILDE(TEMP1_3, TEMP1_33)

  IF(NB.EQ. 1) THEN
    CALL D_ADDM(IDENT_33, TEMP1_33, TEMP2_33, 3, 3)
    CALL D_MTRANS(TEMP2_33, T_D1_D1O, 3, 3)
  ELSEIF(NB.EQ. 2) THEN
    CALL D_ADDM(IDENT_33, TEMP1_33, T_D2O_D2, 3, 3)
  ENDIF

40 CONTINUE

cc Compute transformation D1 <-- D2
  CALL D_MTRANS(T_B1_D1O, T_D1O_B1, 3, 3)
  CALL D_MMUL(T_D1_D1O, T_D1O_B1, T_D1_B1, 3, 3, 3)
  CALL D_MTRANS(T_B1_I, T_I_B1, 3, 3)
  CALL D_MMUL(T_B2_I, T_I_B1, T_B2_B1, 3, 3, 3)
  CALL D_MTRANS(T_B2_B1, T_B1_B2, 3, 3)
  CALL D_MMUL(T_D1_B1, T_B1_B2, TEMP1_33, 3, 3, 3)
  CALL D_MMUL(T_B2_D2O, T_D2O_D2, TEMP2_33, 3, 3, 3)
  CALL D_MTRANS(TEMP2_33, T_D2_B2, 3, 3)
  CALL D_MMUL(TEMP1_33, TEMP2_33, T_D1_D2, 3, 3, 3)

```

```

CALL D_MTRANS(T_D1_D2, T_D2_D1, 3, 3)

cc Compute position vector D2 <-- D1 expressed in D1 frame
DO 45 I=1,3
  TEMP1_3(I) = GSV(2,I) - GSV(1,I)
45 CONTINUE

CALL D_MXV(T_B1_I, TEMP1_3, P_B2_B1_B1, 3, 3)

DO 50 I=1,3
  P_D1_B1_B1(I) = P_D1O_B1_B1(I) + PHIETA(1,NODE_DOCK(1),I)
  TEMP1_3(I) = -P_D1_B1_B1(I) + P_B2_B1_B1(I)
50 CONTINUE

CALL D_MXV(T_D1_B1, TEMP1_3, TEMP2_3, 3, 3)

DO 60 I=1,3
  P_D2_B2_B2(I) = P_D2O_B2_B2(I) + PHIETA(2,NODE_DOCK(2),I)
60 CONTINUE

CALL D_MXV(T_D2_B2, P_D2_B2_B2, TEMP3_3, 3, 3)
CALL D_MXV(T_D1_D2, TEMP3_3, TEMP1_3, 3, 3)

DO 70 I=1,3
  P_D2_D1_D1(I) = TEMP2_3(I) + TEMP1_3(I)
  W_B1_I_B1(I) = QSV(1,3+I)
  W_B2_I_B2(I) = QSV(2,3+I)
  TEMP1_3(I) = PSIETA(1,NODE_DOCK(1),I)
  TEMP2_3(I) = PSIETA(2,NODE_DOCK(2),I)
70 CONTINUE

cc Compute angular velocity of D2 wrt D1 expressed in D2
CALL D_VCROSS(W_B1_I_B1, TEMP1_3, TEMP3_3)
CALL D_VCROSS(W_B2_I_B2, TEMP2_3, TEMP4_3)

DO 80 I=1,3
  TEMP1_3(I) = -TEMP3_3(I) - W_B1_I_B1(I) -
&      PSIETAD(1,NODE_DOCK(1),I)
  TEMP2_3(I) = TEMP4_3(I) + W_B2_I_B2(I) +
&      PSIETAD(2,NODE_DOCK(2),I)
80 CONTINUE

CALL D_MMUL(T_D2_B2, TEMP2_3, TEMP4_3, 3, 3, 3)
CALL D_MMUL(T_D1_B1, TEMP1_3, TEMP3_3, 3, 3, 3)
CALL D_MMUL(T_D2_D1, TEMP3_3, TEMP1_3, 3, 3, 3)

DO 90 I=1,3
  W_D2_D1_D2(I) = TEMP1_3(I) + TEMP4_3(I)
90 CONTINUE

```

```

cc  Compute derivative of D2 <-- D1 position vector wrt D1
cc  expressed in D1
    CALL D_VCROSS(W_B1_I_B1, P_D1_B1_B1, TEMP3_3)
    CALL D_VCROSS(W_B2_I_B2, P_D2_B2_B2, TEMP4_3)

    DO 110 I=1,3
        TEMP1_3(I) = -TEMP3_3(I) - PHJETAD(1,NODE_DOCK(1),I)
    &    - QSV(1,I)
        TEMP2_3(I) = TEMP4_3(I) + PHJETAD(2,NODE_DOCK(2),I)
    &    + QSV(2,I)
        TEMP4_3(I) = PSJETAD(1,NODE_DOCK(1),I)
110  CONTINUE

    CALL D_MXV(T_D1_B1, TEMP1_3, TEMP3_3, 3, 3)
    CALL D_MXV(T_D2_B2, TEMP2_3, TEMP1_3, 3, 3)
    CALL D_MXV(T_D1_D2, TEMP1_3, TEMP2_3, 3, 3)
    CALL D_VCROSS(W_B1_I_B1, TEMP4_3, TEMP1_3)

    DO 120 I=1,3
        TEMP4_3(I) = -W_B1_I_B1(I) - PSJETAD(1,NODE_DOCK(1),I)
    &    - TEMP1_3(I)
120  CONTINUE

    CALL D_MXV(T_D1_B1, TEMP4_3, TEMP1_3, 3, 3)
    CALL D_VCROSS(TEMP1_3, P_D2_D1_D1, TEMP4_3)

    DO 130 I=1,3
        PD_D2_D1_D1(I) = TEMP3_3(I) + TEMP2_3(I) + TEMP4_3(I)
130  CONTINUE

    RETURN
    END

```

APPENDIX B

MATH MODEL GLOBAL VARIABLE DEFINITIONS

DT	-	Integration step size (sec)
EOM(NBODES,3)	-	Equations of motion complexity flags
FN_BODY(NBODES,MNN,3)	-	External forces applied at each body node
FS1_S1(3)	-	Sensed contact forces
GAM_0(NBODES,3,MNM)	-	Array of GAM_0 modal integrals
GAM_1(NBODES,3,MNM)	-	Array of GAM_1 modal integrals
GAM_2JK(NBODES,MNM,MNM,3)	-	Array of GAM_2JK modal integrals
GMASS(NBODES,MNM,MNM)	-	Generalized mass matrices
GSV(NBODES,7+MNM)	-	Generalized coordinate vectors
GSVD(NBODES,7+MNM)	-	Generalized coordinate vector derivatives
GSVDO_1(NBODES,7+MNM)	-	Previous generalized coordinate vector derivatives
GSVDO_2(NBODES,7+MNM)	-	Second previous generalized coordinate vector derivatives
ICM(NBODES,3,3)	-	Body inertia matrices about respective body center of mass
INDX(NBODES,6+MNM)	-	Record of row permutations in LU decomposition
INT	-	Integration option flag
I_1K(NBODES,MNM,3,3)	-	Array of 3X3 I_1K modal integrals
I_2JK(NBODES,MNM,MNM,3,3)	-	Array of 3X3 I_2JK modal integrals
LU_MASSM(NBODES,6+MNM,6+MNM)	-	LU decomposed mass matrices
MASS(NBODES)	-	Body mass
MN_BODY(NBODES,MNN,3)	-	External moments applied at each body node

MNM	-	Maximum number of modes per body
MNN	-	Maximum number of nodes per body
MOD_FREQ(NBODES,MNM)	-	Flexible mode frequencies
MOD_ZETA(NBODES,MNM)	-	Damping per flexible mode
MS1_S1(3)	-	Sensed contact moments
NBODES	-	Number of bodies in simulation
NMODES(NBODES)	-	Number of flexible modes per body
NNODES(NBODES)	-	Number of nodes per body
NODE_DOCK(NBODES)	-	Body docking port node number
NODES(NBODES,MNN,3)	-	Node locations in respective body coordinates
P_B2_B1_B1(3)	-	Position vector of B2 center of mass with respect to B1 center of mass in B1 coordinates
PD_D2_D1_D1(3)	-	Velocity vector of chase docking node with respect to target docking node in D1 coordinates
P_D1O_B1_B1(3)	-	Position vector of undeformed target docking node with respect to target center of mass in B1 coordinates
P_D1_B1_B1(3)	-	Position vector of target docking node with respect to target center of mass in B1 coordinates
P_D1_S1_S1(3)	-	Position vector of target docking node with respect to force/moment sensor in S1 coordinates
P_D2O_B2_B2(3)	-	Position vector of undeformed chase docking node with respect to chase center of mass in B2 coordinates
P_D2_B2_B2(3)	-	Position vector of chase docking node with respect to chase center of mass in B2 coordinates

P_D2_D1_D1(3)	-	Position vector of chase docking node with respect to target docking node in D1 coordinates
PHI(NBODES,MNN,MNM,3)	-	Translational modal gains
PHIETA(NBODES,MNN,3)	-	Node translational deformation in respective body coordinates
PHIETAD(NBODES,MNN,3)	-	Node translational deformation velocity in respective body coordinates
PSI(NBODES,MNN,MNM,3)	-	Rotational modal gains
PSIETA(NBODES,MNN,3)	-	Node rotational deformation in respective body coordinates
PSIETAD(NBODES,MNN,3)	-	Node rotational deformation velocity in respective body coordinates
QSV(NBODES,6+MNM)	-	Quasi-coordinate vectors
QSVD(NBODES,6+MNM)	-	Quasi-coordinate vector derivatives
QSVDO_1(NBODES,6+MNM)	-	Previous quasi-coordinate vector derivatives
QSVDO_2(NBODES,6+MNM)	-	Second previous quasi-coordinate vector derivatives
T_B1_B2(3,3)	-	Transformation from B2 to B1 coordinates
T_B1_D1O(3,3)	-	Transformation from undeformed target docking frame to B1 coordinates
T_B1_I(3,3)	-	Transformation from inertial to B1 coordinates
T_B2_B1(3,3)	-	Transformation from B1 to B2 coordinates
T_B2_D2O(3,3)	-	Transformation from undeformed chase docking frame to B2 coordinates
T_B2_I(3,3)	-	Transformation from inertial to B2 coordinates
T_D1_B1(3,3)	-	Transformation from B1 to D1 coordinates
T_D1_D2(3,3)	-	Transformation from D2 to D1 coordinates

T_D1_S1(3,3)	- Transformation from force/moment sensor to D1 coordinates
T_D2_B2(3,3)	- Transformation from B2 to D2 coordinates
TIME	- Simulation time (sec)
W_D2_D1_D2(3)	- Angular velocity of chase docking frame with respect to target docking frame in D2 coordinates

APPENDIX C

FLEXIBLE BODY DATA PRE-PROCESSING CODE

```

cccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccc
program massint
cccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccc
cc
cc Purpose:      Compute mass integral terms for a consistent mass
cc                matrix. Output modal data for use in flex body sims.
cc
cc Source:       MC and JC      4/91
cc
cc Explicit Inputs:
cc   ndof                - number of degrees of freedom (DOF) of
cc                        the system
cc   nnode               - number of nodes of the system
cc   nmode               - number of modes of the system
cc   idof_type(ndof)     - DOF array from DOF map
cc   inod(ndof)          - node array from DOF map
cc   mass_mat(ndof,ndof) - consistent mass matrix
cc   gen_mass(nmode,nmode) - generalized mass mat ( $\phi^T \text{mass\_mat} \phi$ )
cc   phi(ndof,nmode)     - matrix of eigenvectors (mode shapes) of the
cc                        system
cc   node_mat(nnode,3)   - nodal geometry matrix
cc   r_offset(3)         - offset vector for nodal geometry matrix
cc
cc Explicit Outputs:
cc   gen_mass            - generalized mass matrix ( $\phi^T M \phi$ )
cc   gamma_0(3,nmode)    - zeroth mass integral vector
cc   gamma_1(3,nmode)    - first mass integral vector
cc   gamma_2(nmode,nmode,3) - second mass integral vector
cc   I_1(nmode,3,3)      - first mass integral matrix
cc   I_2(nmode,nmode,3,3) - second mass integral matrix
cc
cc Others:
cc   rho(nnode,3)        - offset nodal geometry matrix
cc   a(ndof,nmode)       - matrix of the "a_j" vectors, stored column-wise
cc   a_t(nnode,nmode,3)  - subset of the "a" matrix describing the
cc                        translation of the nodes (x,y,z)
cc   a_a(nnode,nmode,3)  - subset of the "a" matrix describing the
cc                        rotation of the nodes (theta_x, theta_y, theta_z)
cc   phi_t(nnode,nmode,3) - subset of the modal matrix describing the
cc                        translation of the nodes (x,y,z)
cc   psi(nnode,nmode,3)  - subset of the modal matrix describing the
cc                        rotation of the nodes (theta_x, theta_y, theta_z)
cc

```



```

cc "A" matrix variables
  real*8 a(ndofmax,nmodemax)
  real*8 a_t(nnodemax, nmodemax, 3), a_a(nnodemax, nmodemax, 3)
  real*8 a_t_i(3), a_x_phi(3), a_dot_phi
  real*8 a_phi_mat(3,3)

cc Mass integral vectors
  real*8 gamma_0(3,nmodemax), gamma_1(3,nmodemax)
  real*8 gamma_2(nmodemax, nmodemax, 3)

cc First mass integral matrix and associated temporary variables
  real*8 l_1_1(3,3), l_1_2(3,3), l_1_3(3,3), l_1_4(3,3)
  real*8 l_1(nmodemax,3,3), l_1_n_old(3,3)

cc Second mass integral matrix and associated temporary variables
  real*8 l_2_1(3,3), l_2_2(3,3), l_2_3(3,3), l_2_4(3,3)
  real*8 l_2(nmodemax,nmodemax,3,3), l_2_n_old(3,3)

cc 3x3 zero and identity matrices
  real*8 zeros(3,3), ident(3,3)
  character*15 file1, file2

  data zeros / 9*0.d0 /
  data ident / 1.d0, 3*0.d0, 1.d0, 3*0.d0, 1.d0/

cc Open data files
  write(6,*)
  write(6,'(2x,"Enter raw data file name: ",$)')
  read(5,'(a)') file1
  write(6,'(2x,"Enter output data file name: ",$)')
  read(5,'(a)') file2
  write(6,*)

  open(unit=1, file = file1, status = 'old',
    & form = 'formatted')
  open(unit=2, file = file2, status = 'unknown',
    & form = 'formatted')

cc Read DOF map
  read(1,*) nnode,ndof,nmode
  write(6,'(4x,"nnode; ndof; nmode: ",3i5)') nnode,ndof,nmode
  read(1,*)
  do 10 i=1,ndof
    read(1,*) inod(i), idof_type(i)
10  continue

cc Read nodal geometry matrix
  read(1,*)

```

```

    read(1,*) (r_offset(i),i=1,3)
    do 35 i=1,nnode
        read(1,*) (node_mat(i,j),j=1,3)
35    continue

cc Read mass matrix
    read(1,*)
    do 20 i=1,ndof
        read(1,*) (mass_mat(i,j),j=1,ndof)
20    continue

    write(6,*) '    Mass matrix read complete'

cc Read modal (eigenvector) matrix
    read(1,*)
    do 30 i=1,ndof
        read(1,*) (phi(i,j),j=1,nmode)
30    continue

    write(6,*) '    Modal matrix read complete'

cc Calculate the nodal position vector based on the nodal geometry
cc matrix and the offset vector
    do 33 i=1,nnode
        do 34 j=1,3
            rho(i,j) = node_mat(i,j) - r_offset(j)
34        continue
33    continue

cc Calculate generalized mass matrix
    call d_mmul(mass_mat, phi, temp, ndof, ndof, nmode,
&             ndofmax, ndofmax, nmodemax)

    do 36 i=1,nmode
        do 37 j=1,nmode
            gen_mass(i,j) = 0.d0
            do 38 k=1,ndof
                gen_mass(i,j) = gen_mass(i,j) + phi(k,i) * temp(k,j)
38            continue
37        continue
36    continue

    write(6,*) '    Generalized mass matrix formed'

cc Calculate the "a" matrix
    do 60 j=1,nmode
        do 40 i=1,ndof
            a(i,j) = 0.d0
            do 50 k=1,ndof

```

```

        a(i,j) = a(i,j) + mass_mat(i,k)*phi(k,j)
50    continue
40    continue
60    continue

    write(6,*) '  A matrix formed'

cc Divide the "a" matrix into translational and rotational subsets
do 80 j=1,nmode
    do 70 i=1,ndof
        if(idof_type(i).le.3) then
            a_t(inod(i), j, idof_type(i)) = a(i,j)
        elseif(idof_type(i).ge.4.and.idof_type(i).le.6) then
            a_a(inod(i), j, (idof_type(i)-3)) = a(i,j)
        endif
70    continue
80    continue

cc Calculate the zeroth mass integral vector
do 110 j=1,nmode
    do 100 k=1,3
        gamma_0(k,j) = 0.0d0
        do 90 i=1,nnode
            gamma_0(k,j) = gamma_0(k,j) + a_t(i,j,k)
90    continue
100    continue
110    continue

    write(6,*) '  Gamma_0J calculated'

cc Calculate the first mass integral vector
do 150 j=1,nmode
    do 130 k=1,3
        gamma_1(k,j) = 0.0d0
        do 140 i=1,nnode
            do 141 n=1,3
                a_t_i(n) = a_t(i,j,n)
                rho_i(n) = rho(i,n)
141            continue
                call d_vcross( rho_i, a_t_i, rho_x_a )
                gamma_1(k,j) = gamma_1(k,j) + a_a(i,j,k) + rho_x_a(k)
140            continue
130        continue
150    continue

    write(6,*) '  Gammo_1J calculated'

cc Divide the modal matrix into translational and rotational subsets
do 170 j=1,nmode

```

```

do 160 i=1,ndof
  if(idof_type(i).le.3) then
    phi_t( inod(i), j, idof_type(i) ) = phi(i,j)
  elseif(idof_type(i).ge.4.and.idof_type(i).le.6) then
    psi( inod(i), j, (idof_type(i)-3) ) = phi(i,j)
  endif
160  continue
170  continue

  write(6,*) ' Separation of trans and rot modes'

cc Initialize counter
  icount = 1

cc Calculate the second mass integral vector
  do 230 j=icount,nmode
    do 220 k=icount,nmode
      do 180 i=1,3
        gamma_2(j,k,i) = 0.0d0
180      continue
      do 210 i1 = 1,nnode
        do 181 i=1,3
          a_t_i(i) = a_t(i1,j,i)
          phi_t_i(i) = phi_t(i1,k,i)
181        continue
          call d_vcros( a_t_i, phi_t_i, a_x_phi )
          do 200 i=1,3
            gamma_2(j,k,i) = gamma_2(j,k,i) + a_x_phi(i)
            gamma_2(k,j,i) = -gamma_2(j,k,i)
200          continue
210        continue
220      continue
      icount = icount+1
230    continue

    write(6,*) ' Gamma_2JK calculated'

cc Calculate the first mass integral matrix
  do 280 k=1,nmode
    call d_fmim( l_1, zeros, k, nmodemax )
    do 270 j=1,nnode
      call d_mf3dm( l_1, l_1_n_old, k, nmodemax )
      call d_vf3dm( a_t, a_t_i, j, k, nnodemax, nmodemax )
      call d_vf3dm( phi_t, phi_t_i, j, k, nnodemax, nmodemax )
      call d_vf2dm( rho, rho_i, j, nnodemax )
      call d_dot( rho_i, a_t_i, rho_dot_a )
      call d_cxm( ident, rho_dot_a, l_1_1, 3, 3 )
      call d_mmul( phi_t_i, rho_i, rho_phi_mat, 3,1,3,3,1,3)
      call d_chsn( rho_phi_mat, l_1_2, 3, 3 )

```



```

        call d_addm( l_1_1, l_1_2, l_1_3, 3, 3 )
        call d_addm( l_1_3, l_1_n_old, l_1_4, 3, 3 )
        call d_fmim( l_1, l_1_4, k, nmodemax )
270    continue
280    continue

    write(6,*) '    l_1K calculated'

cc Calculate the second mass integral matrix
do 310 j=1,nmode
    do 300 k=1,nmode
        call d_smim( l_2, zeros, j, k, nmodemax )
        do 290 i=1,nnode
            call d_mf4dm( l_2, l_2_n_old, j, k, nmodemax )
            call d_vf3dm( a_t, a_t_i, i, j, nmodemax, nmodemax )
            call d_vf3dm( phi_t, phi_t_i, i, k, nmodemax,
&                nmodemax )
            call d_dot( a_t_i, phi_t_i, a_dot_phi )
            call d_cxm( ident, a_dot_phi, l_2_1, 3, 3 )
            call d_mmul( a_t_i, phi_t_i, a_phi_mat, 3, 1, 3, 3, 1, 3 )
            call d_chsn( a_phi_mat, l_2_2, 3, 3 )
            call d_addm( l_2_1, l_2_2, l_2_3, 3, 3 )
            call d_addm( l_2_3, l_2_n_old, l_2_4, 3, 3 )
            call d_smim( l_2, l_2_4, j, k, nmodemax )
290        continue
300    continue
310    continue

```

```

    write(6,*) '    l_2JK calculated'
    write(6,*)

```

```

cc Write output in format consistent with NAMTB flex body
cc simulation input

```

```

    write(6, '(2x, "Writing output data to: ", a15)') file2

```

```

    write(6,*)

```

```

    write(6,*) '    Generalized mass matrix'
    do i=1,nmode
        write(2,*) (gen_mass(i,j),j=1,nmode)
    enddo

```

```

    write(6,*) '    Trans and rot modal gains'
    do i=1,nnode
        do j=1,nmode
            write(2,500) (phi_t(i,j,k),k=1,3)
        enddo
    enddo

```